

PHY 712 Electrodynamics
9-9:50 AM MWF Olin 105

Plan for Lecture 12:

Continue reading Chapter 5

A. Examples of magnetostatic fields

B. Magnetic dipoles

C. Hyperfine interaction

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Course schedule for Spring 2019
(Preliminary schedule -- subject to frequent adjustment.)


Lecture date	JDJ Reading	Topic	HW	Due date
1 Mon: 01/14/2019	Chap. 1 & Appen.	Introduction, units and Poisson equation	#1	01/23/2019
2 Wed: 01/16/2019	Chap. 1	Electrostatic energy calculations	#2	01/23/2019
3 Fri: 01/18/2019	Chap. 1	Electrostatic potentials and fields	#3	01/23/2019
Mon: 01/21/2019	No class	Martin Luther King Holiday		
4 Wed: 01/23/2019	Chap. 1 - 3	Poisson's equation in 2 and 3 dimensions		
5 Fri: 01/25/2019	Chap. 1 - 3	Brief introduction to numerical methods	#4	01/28/2019
6 Mon: 01/28/2019	Chap. 2 & 3	Image charge constructions	#5	01/30/2019
7 Wed: 01/30/2019	Chap. 2 & 3	Cylindrical and spherical geometries		
8 Fri: 02/01/2019	Chap. 3 & 4	Spherical geometry and multipole moments	#6	02/04/2019
9 Mon: 02/04/2019	Chap. 4	Dipoles and Dielectrics	#7	02/06/2019
10 Wed: 02/06/2019	Chap. 4	Polarization and Dielectrics		
11 Fri: 02/08/2019	Chap. 5	Magnetostatics	#8	02/11/2019
12 Mon: 02/11/2019	Chap. 5	Magnetic dipoles and hyperfine interaction	#9	02/13/2019
13 Wed: 02/13/2019				
14 Fri: 02/15/2019				
15 Mon: 02/18/2019				

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Events

Colloquium: "Tailored Light-Matter Interactions in Scalable and Artificial Nanomaterials" - Wednesday, February 13, 2019, at 4:00 PM
Peijun Guo, PhD, Named Fellowship-Enrico Fermi, Argonne National Laboratory George P. Williams, Jr. Lecture Hall, (Olin 101)
Wednesday, February 13, 2019, at 4:00 PM
There will be a reception with ...

News

Randall Ledford receives Physics Department Outstanding Alumni Award

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Various forms of Ampere's law :

$$\nabla \times \mathbf{B}(\mathbf{r}) = \mu_0 \mathbf{J}(\mathbf{r})$$

Vector potential: $\mathbf{B}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r})$

For Coulomb gauge: $\nabla \cdot \mathbf{A}(\mathbf{r}) = 0$

$$\Rightarrow \nabla^2 \mathbf{A}(\mathbf{r}) = -\mu_0 \mathbf{J}(\mathbf{r})$$

For confined current density :

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^3 r' \frac{\mathbf{J}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

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Other examples of current density sources:

Quantum mechanical expression for current density

for a particle of mass M and charge e and of probability amplitude $\Psi(\mathbf{r})$:

$$\mathbf{J}(\mathbf{r}) = -\frac{e\hbar}{2Mi} (\Psi^*(\mathbf{r}) \nabla \Psi(\mathbf{r}) - \Psi(\mathbf{r}) \nabla \Psi^*(\mathbf{r}))$$

For an electron in a spherical potential (such as in an atom):

$$\Psi(\mathbf{r}) \equiv \Psi_{nlm}(\mathbf{r}) = R_{nl}(r) Y_{lm}(\hat{\mathbf{r}})$$

$$\begin{aligned} \mathbf{J}(\mathbf{r}) &= \frac{e\hbar}{2Mi} |R_{nl}(r)|^2 \frac{1}{r \sin \theta} \left(Y_{lm}^*(\hat{\mathbf{r}}) \frac{\partial Y_{lm}(\hat{\mathbf{r}})}{\partial \varphi} - Y_{lm}(\hat{\mathbf{r}}) \frac{\partial Y_{lm}^*(\hat{\mathbf{r}})}{\partial \varphi} \right) \hat{\boldsymbol{\phi}} \\ &= \frac{e\hbar}{M} \frac{m_l}{r \sin \theta} |\Psi_{nlm}(\mathbf{r})|^2 \hat{\boldsymbol{\phi}} \end{aligned}$$

Note that: $\hat{\boldsymbol{\phi}} = -\sin \varphi \hat{\mathbf{x}} + \cos \varphi \hat{\mathbf{y}} = \frac{\hat{\mathbf{z}} \times \mathbf{r}}{r \sin \theta}$

$$\mathbf{J}(\mathbf{r}) = \frac{e\hbar}{M} \frac{m_l}{r^2 \sin^2 \theta} |\Psi_{nlm}(\mathbf{r})|^2 (\hat{\mathbf{z}} \times \mathbf{r})$$

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Details of the electron orbital magnetic dipole moment

$$\mathbf{J}(\mathbf{r}) = \frac{e\hbar}{m_e} \frac{m_l}{r \sin \theta} |\Psi_{nlm}(\mathbf{r})|^2 \hat{\boldsymbol{\phi}}$$

Note that: $\hat{\boldsymbol{\phi}} = -\sin \varphi \hat{\mathbf{x}} + \cos \varphi \hat{\mathbf{y}}$

Magnetic dipole moment:

$$\begin{aligned} \mathbf{m} &= \frac{1}{2} \int d^3 r' \mathbf{r}' \times \mathbf{J}(\mathbf{r}') = -\frac{e\hbar m_l}{2m_e} \int d^3 r' \frac{\mathbf{r}' \times \hat{\boldsymbol{\phi}}}{r' \sin \theta'} |\Psi_{nlm}(\mathbf{r}')|^2 \\ &= -\frac{e\hbar m_l}{2m_e} \int d^3 r' \frac{-r' \hat{\boldsymbol{\theta}}}{r' \sin \theta'} |\Psi_{nlm}(\mathbf{r}')|^2 \end{aligned}$$

Note that: $\hat{\boldsymbol{\theta}} = \cos \theta \cos \varphi \hat{\mathbf{x}} + \cos \theta \sin \varphi \hat{\mathbf{y}} - \sin \theta \hat{\mathbf{z}}$

$$\begin{aligned} \mathbf{m} &= -\frac{e\hbar m_l \hat{\mathbf{z}}}{2m_e} \int d^3 r' |\Psi_{nlm}(\mathbf{r}')|^2 \\ &= -\frac{e\hbar m_l}{2m_e} \hat{\mathbf{z}} \end{aligned}$$

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Summary of magnetic field generated by point magnetic dipole moment:

$$\mathbf{B}_{\mu_e}(\mathbf{r}) = \frac{\mu_0}{4\pi} \left(\frac{3\hat{\mathbf{r}}(\boldsymbol{\mu}_e \cdot \hat{\mathbf{r}}) - \boldsymbol{\mu}_e}{r^3} + \frac{8\pi}{3} \boldsymbol{\mu}_e \delta(\mathbf{r}) \right)$$

Magnetic field near nucleus due to orbiting electron:

$$\mathbf{B}_O(\mathbf{r}) = \frac{\mu_0}{4\pi} \frac{e}{m_e} L_z \hat{\mathbf{z}} \left(\frac{1}{r^3} \right)$$

"Hyperfine" interaction energy:

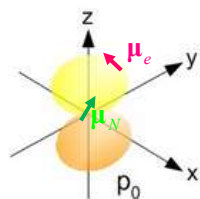
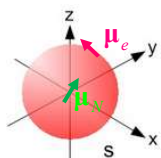
$$\begin{aligned} \mathcal{H}_{HF} &= -\boldsymbol{\mu}_N \cdot (\mathbf{B}_{\mu_e}(\mathbf{r}) + \mathbf{B}_O(\mathbf{r})) \\ &= \frac{\mu_0}{4\pi} \left(\frac{3(\boldsymbol{\mu}_N \cdot \hat{\mathbf{r}})(\boldsymbol{\mu}_e \cdot \hat{\mathbf{r}}) - \boldsymbol{\mu}_N \cdot \boldsymbol{\mu}_e}{r^3} + \frac{8\pi}{3} \boldsymbol{\mu}_N \cdot \boldsymbol{\mu}_e \delta(\mathbf{r}) + \frac{e}{m_e} \left(\frac{\mathbf{L} \cdot \boldsymbol{\mu}_N}{r^3} \right) \right) \end{aligned}$$

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$$\mathcal{H}_{HF} = \frac{\mu_0}{4\pi} \left(\frac{3(\boldsymbol{\mu}_N \cdot \hat{\mathbf{r}})(\boldsymbol{\mu}_e \cdot \hat{\mathbf{r}}) - \boldsymbol{\mu}_N \cdot \boldsymbol{\mu}_e}{r^3} + \frac{8\pi}{3} \boldsymbol{\mu}_N \cdot \boldsymbol{\mu}_e \delta(\mathbf{r}) + \frac{e}{m_e} \left(\frac{\mathbf{L} \cdot \boldsymbol{\mu}_N}{r^3} \right) \right)$$



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