

PHY 712 Electrodynamics
9-9:50 AM MWF Olin 105

Plan for Lecture 2:

Reading: Chapter 1 (especially 1.11) in JDJ;
Ewald summation methods

1. Motivation
2. Expression to evaluate the electrostatic energy of an extended periodic system
3. Examples

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PHY 712 Electrodynamics

MWF 9-9:50 AM | OPL 105 | <http://www.wfu.edu/~natalie/s18phy712/>

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Course schedule for Spring 2018
 (Preliminary schedule -- subject to frequent adjustment.)

Lecture date	JDJ Reading	Topic	HW	Due date
Wed: 01/17/2018	No class	Snow		
1 Fri: 01/19/2018	Chap. 1 & I	Introduction, units and Poisson equation	#1	01/26/2018
2 Mon: 01/22/2018	Chap. 1	Electrostatic energy calculations	#2	01/26/2018
3 Wed: 01/24/2018				
4 Thu: 01/25/2018				
5 Fri: 01/26/2018				
6 Mon: 01/29/2018				
7 Wed: 01/31/2018				

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Ewald summation methods -- motivation

Consider a collection of point charges $\{q_i\}$ located at points $\{\mathbf{r}_i\}$.
 The energy to separate these charges to infinity ($\mathbf{r}_i \rightarrow \infty$) is

$$W = \frac{1}{4\pi\epsilon_0} \sum_{(i,j); i>j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$$

Here the summation is over all pairs of (i, j) , excluding $i = j$.
 It is convenient to sum over all particles and divide by 2 in order to compensate for the double counting:

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i,j; i \neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$$

Now the summation is over all i and j , excluding $i = j$.
 The energy W scales as the number of particles N . As $N \rightarrow \infty$, the ratio W / N remains well-defined in principle, but difficult to calculate in practice.

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Ewald summation methods – slight digression

When the discrete charge distribution becomes a continuous charge density: $q_i \rightarrow \rho(\mathbf{r})$, the electrostatic energy becomes

$$W = \frac{1}{8\pi\epsilon_0} \int d^3r \int d^3r' \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Notice, in this case, it is not possible to exclude the "self-interaction".

Electrostatic energy in terms of $\Phi(\mathbf{r})$ and field $\mathbf{E}(\mathbf{r})$:

Previous expression can be rewritten in terms of the electrostatic potential or field:

$$W = \frac{1}{2} \int d^3r \rho(\mathbf{r})\Phi(\mathbf{r}) = -\frac{\epsilon_0}{2} \int d^3r (\nabla^2\Phi(\mathbf{r}))\Phi(\mathbf{r})$$

$$W = \frac{\epsilon_0}{2} \int d^3r |\nabla\Phi(\mathbf{r})|^2 = \frac{\epsilon_0}{2} \int d^3r |\mathbf{E}(\mathbf{r})|^2$$

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Evaluation of the electrostatic energy for N point charges:

$$\frac{W}{N} = \frac{1}{8\pi\epsilon_0} \frac{1}{N} \sum_{i,j;i \neq j} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$$

Ewald summation methods – exact results for periodic systems

$$\frac{W}{N} = \sum_{\alpha\beta} \frac{q_\alpha q_\beta}{8\pi\epsilon_0} \left(\frac{4\pi}{\Omega} \sum_{\mathbf{G} \neq 0} \frac{e^{-\alpha|\mathbf{G} \cdot \mathbf{r}_{\alpha\beta}} e^{-\beta G^2/\eta}}{G^2} - \sqrt{\frac{\eta}{\pi}} \delta_{\alpha\beta} + \sum_{\mathbf{T}} \frac{\text{erfc}(\frac{1}{2}\sqrt{\eta}|\mathbf{r}_{\alpha\beta} + \mathbf{T}|)}{|\mathbf{r}_{\alpha\beta} + \mathbf{T}|} \right) - \frac{4\pi Q^2}{8\pi\epsilon_0 \Omega \eta}$$

Note that the results should not depend upon η (assuming that all summations are carried to convergence). In the example of CsCl having a lattice constant a , we show two calculations produce the result:

$$\frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070722970}{a} \quad \text{or} \quad \frac{W}{N} = -\frac{e^2}{8\pi\epsilon_0} \frac{4.070723039}{a}$$

See lecture notes for details.

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