

PHY 752 Solid State Physics
11-11:50 AM MWF Olin 107

Plan for Lecture 7:

Reading: Chapter 7 in MPM; Electronic Structure

- 1. Bloch's Theorem**
- 2. Eigenstates of a simple model potential**

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Course schedule for Spring 2015
(Preliminary schedule -- subject to frequent adjustment.)

Lecture date	MPM Reading	Topic	Assign.	Due date
1 Mon: 01/12/2015	Chap. 1 & 2	Crystal structures	#1	01/23/2015
2 Wed: 01/14/2015	Chap. 1 & 2	Some group theory	#2	01/23/2015
Fri: 01/16/2015	No class	NAWH out of town		
Mon: 01/19/2015	No class	MLK Holiday		
3 Wed: 01/21/2015	Chap. 1 & 2	Some group theory	#3	01/23/2015
4 Fri: 01/23/2015	Chap. 1 & 2	Some more group theory	#4	01/26/2015
5 Mon: 01/26/2015	Chap. 7.3	Some more group theory	#5	01/28/2015
6 Wed: 01/28/2015	Chap. 6	Electronic structure; Free electron gas	#6	01/30/2015
7 Fri: 01/30/2015	Chap. 7	Electronic structure; Model potentials	#7	02/02/2015

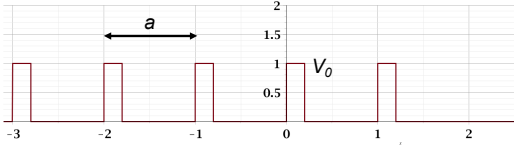
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Consider an electron moving in a one-dimensional model potential (Kronig and Penney, *Proc. Roy. Soc. (London)* **130**, 499 (1931))

Schrodinger equation for electron:

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V(x) \right) \Psi(x) = E \Psi(x)$$

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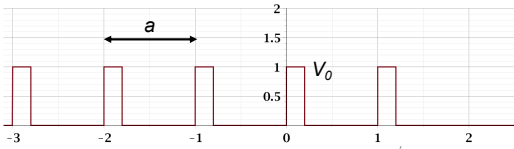
Effects of periodicity: $V(x) = V(x + na)$

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V(x)\right) \Psi(x) = E \Psi(x)$$

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V(x + na)\right) \Psi(x + na) = E \Psi(x + na)$$

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V(x)\right) \Psi(x + na) = E \Psi(x + na)$$

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Since $\Psi(x + na)$ and $\Psi(x)$ are solutions of the same eigenvalue problem: $\Psi(x + na) = K \Psi(x)$

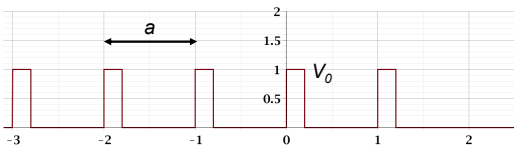
Assume $K = e^{i\theta} = e^{ikna}$

Bloch theorem: $\Psi_k(x + na) = e^{ikna} \Psi_k(x)$

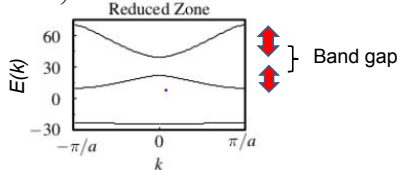
$$\Psi_{kk}(x) = e^{ikx} u(x)$$

where $u_k(x + na) = u_k(x)$

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Eigenstates:

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V(x)\right) \Psi_k(x) = E(k) \Psi_k(x)$$


Reduced Zone

Band gap

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What causes band gaps in the electronic structure?

Consider a single potential well

Spectrum

E

x

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Define: $\alpha = \sqrt{\frac{2mE}{\hbar^2}}$ $\beta = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$

For $0 \leq x \leq (a-b)$ $\Psi_1(x) = Ae^{\beta x} + Be^{-\beta x}$

For $(a-b) \leq x \leq a$ $\Psi_2(x) = Ce^{i\alpha x} + De^{-i\alpha x}$

Continuity conditions: $\Psi_1(0) = \Psi_2(0)$ $\frac{d\Psi_1(0)}{dx} = \frac{d\Psi_2(0)}{dx}$

$\Psi_1(a-b) = \Psi_2(a-b)$ $\frac{d\Psi_1(a-b)}{dx} = \frac{d\Psi_2(a-b)}{dx}$

Also note: $\Psi(x+a) = e^{ika}\Psi(x)$

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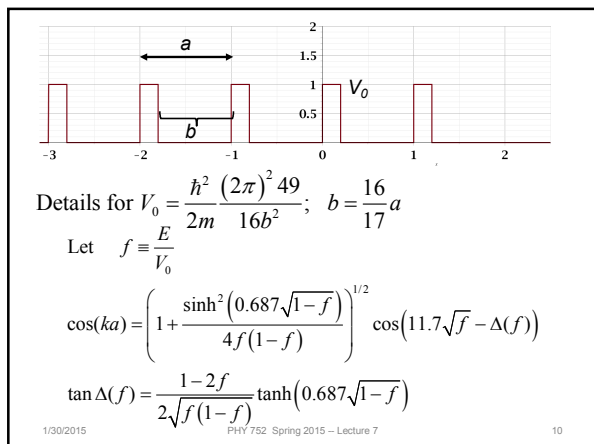
Matching conditions reduce to:

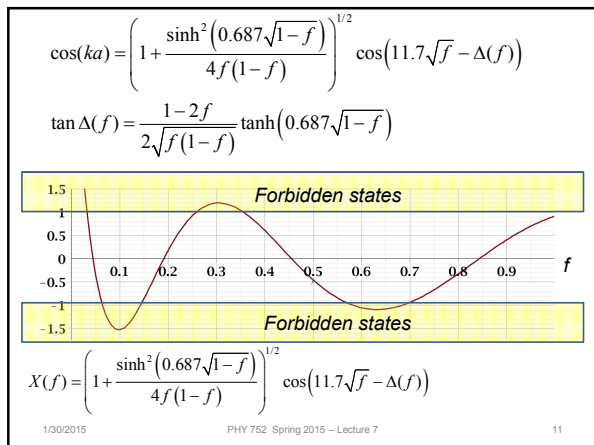
$\cos(ka) = F(E)\cos(ab - \Delta(E))$

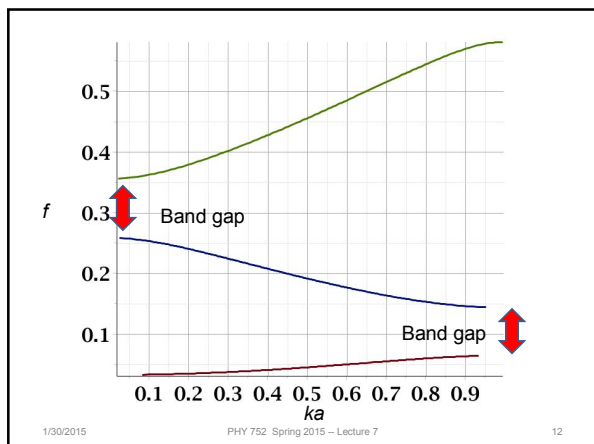
$F(E) \equiv \left(1 + \frac{V_0^2}{4E(E - V_0)} \sinh^2(\beta(a-b))\right)^{1/2}$

$\tan \Delta(E) = \frac{V_0 - 2E}{2\sqrt{E(V_0 - E)}} \tanh(\beta(a-b))$

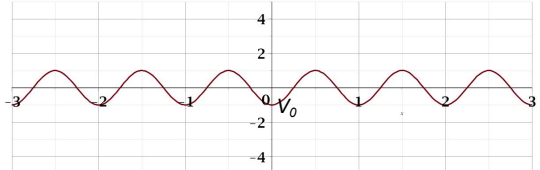
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Electrons in the presence of a weak periodic potential



$$\mathcal{H} = \mathcal{H}^0 + \mathcal{H}^1$$

$$\mathcal{H}^0 = -\frac{\hbar^2 d^2}{2m dx^2} \quad \mathcal{H}^1 = V_0 \cos\left(\frac{2\pi x}{a}\right)$$

$$\Psi(x) = \sum_n C_n e^{i\left(k+n\frac{2\pi}{a}\right)x} \quad -\frac{\pi}{a} \leq k \leq \frac{\pi}{a}$$

$$E_{nk}^0 = \frac{\hbar^2 \left|k + n\frac{2\pi}{a}\right|^2}{2m}$$

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$$\mathcal{H}^0 = -\frac{\hbar^2 d^2}{2m dx^2} \quad \mathcal{H}^1 = V_0 \cos\left(\frac{2\pi x}{a}\right)$$

$$\Psi(x) = \sum_n C_n e^{i\left(k+n\frac{2\pi}{a}\right)x} \quad -\frac{\pi}{a} \leq k \leq \frac{\pi}{a}$$

$$E_{nk}^0 = \frac{\hbar^2 \left|k + n\frac{2\pi}{a}\right|^2}{2m}$$

Note that $E_{0\pm\frac{\pi}{a}}^0 = E_{-1\pm\frac{\pi}{a}}^0 = \frac{\hbar^2 \pi^2}{2ma^2}$

Degenerate perturbation theory

$$\begin{pmatrix} \frac{\hbar^2 \pi^2}{2ma^2} & V_0 \\ V_0 & \frac{\hbar^2 \pi^2}{2ma^2} \end{pmatrix} \begin{pmatrix} C_0 \\ C_1 \end{pmatrix} = E^1 \begin{pmatrix} C_0 \\ C_1 \end{pmatrix}$$

$$E^1 = \pm V_0$$

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