

**PHY 712 Electrodynamics
9-9:50 AM MWF Olin 103**

Plan for Lecture 32:

Special Topics in Electrodynamics:

Cherenkov radiation

References: Jackson Chapter 13.4

Zangwill Chapter 23.7

Smith Chapter 6.4

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20	Mon: 03/16/2015	Chap. 8	Review Exam; Wave guides	#19	03/16/2015
21	Wed: 03/18/2015	Chap. 8	Wave guides	#20	03/20/2015
22	Fri: 03/20/2015	Chap. 9	Radiation sources	#21	03/23/2015
23	Mon: 03/23/2015	Chap. 9 & 10	Radiation and scattering	#22	03/25/2015
24	Wed: 03/25/2015	Chap. 9 & 10	Radiation and scattering		
25	Fri: 03/27/2015	Chap. 11	Special relativity	#23	03/30/2015
26	Mon: 03/30/2015	Chap. 11	Special relativity	#24	04/01/2015
27	Wed: 04/01/2015	Chap. 11	Special relativity	#25	04/06/2015
	Fri: 04/03/2015	Good Friday	No class		
28	Mon: 04/06/2015	Chap. 14	Radiation from moving charges	#26	04/08/2015
29	Wed: 04/08/2015	Chap. 14	Radiation from moving charges	#27	04/10/2015
30	Fri: 04/10/2015	Chap. 14	Radiation from moving charges	#28	04/13/2015
31	Mon: 04/13/2015	Chap. 15	Radiation due to scattering	#29	04/15/2015
32	Wed: 04/15/2015	Chap. 13	Cherenkov radiation	#30	04/17/2015
33	Fri: 04/17/2015		Special topics -- superconductivity		04/20/2015
34	Mon: 04/20/2015		Special topics -- superconductivity		
35	Wed: 04/22/2015		Review		
36	Fri: 04/24/2015		Review		
	Mon: 04/27/2015		Presentations I		
	Wed: 04/29/2015		Presentations II		
	Fri: 05/01/2015		Presentations III & Take home exam		

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Department of Physics

News



Senior Abdul Obaid awarded Gates Cambridge Scholarship



Senior Derek Fogel wins Best Presentation Award at APS March Meeting



Prof. Jurchescu receives 2015 Excellence in Research Award

Events

Wed. Apr. 15, 2015
Physics Colloquium:
Biomolecular Recognition
Prof. Wang, SUNY Stony Brook
Olin 101 4:00 PM
Refreshments at 3:30 PM
Olin Lobby

Profiles in Physics

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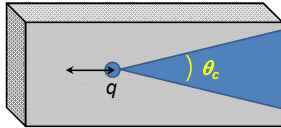
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References for notes: Glenn S. Smith, *An Introduction to Electromagnetic Radiation* (Cambridge UP, 1997), Andrew Zangwill, *Modern Electrodynamics* (Cambridge UP, 2013)

Cherenkov radiation

Discovered ~1930; bluish light emitted by energetic charged particles traveling within dielectric materials



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Maxwell's potential equations within a material having permittivity and permeability (Lorentz gauge; cgs Gaussian units)

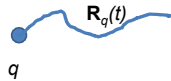
$$\nabla^2 \Phi - \mu\epsilon \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = -\frac{4\pi}{\epsilon} \rho$$

$$\nabla^2 \mathbf{A} - \mu\epsilon \frac{1}{c^2} \frac{\partial^2 \mathbf{A}}{\partial t^2} = -\frac{4\pi\mu}{c} \mathbf{J}$$

Source: charged particle moving on trajectory $\mathbf{R}_q(t)$:

$$\rho(\mathbf{r}, t) = q \delta(\mathbf{r} - \mathbf{R}_q(t))$$

$$\mathbf{J}(\mathbf{r}, t) = q \dot{\mathbf{R}}_q(t) \delta(\mathbf{r} - \mathbf{R}_q(t))$$



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Liénard-Wiechert potential solutions:

$$\Phi(\mathbf{r}, t) = \frac{q}{\epsilon} \frac{1}{R(t_r) - \boldsymbol{\beta}_n \cdot \mathbf{R}(t_r)}$$

$$\mathbf{A}(\mathbf{r}, t) = q\mu \frac{\boldsymbol{\beta}_n}{R(t_r) - \boldsymbol{\beta}_n \cdot \mathbf{R}(t_r)}$$

$$\mathbf{R}(t_r) \equiv \mathbf{r} - \mathbf{R}_q(t_r)$$

$$\boldsymbol{\beta}_n(t_r) \equiv \frac{\dot{\mathbf{R}}_q(t_r)}{c_n} \quad c_n \equiv \sqrt{\mu\epsilon} c \equiv \frac{c}{n}$$

$$t_r = t - \frac{R(t_r)}{c_n}$$

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Some algebra

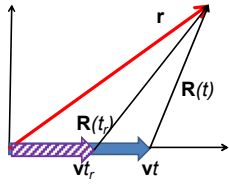
$$\mathbf{R}(t_r) = \mathbf{r} - \mathbf{v}t_r$$

$$\mathbf{R}(t_r) = \mathbf{r} - \mathbf{v}t_r = \mathbf{R}(t) + \mathbf{v}(t - t_r)$$

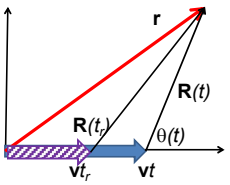
$$(t - t_r)c_n = R(t_r) = |\mathbf{R}(t) + \mathbf{v}(t - t_r)|$$

Quadratic equation for $(t - t_r)c_n$:

$$((t - t_r)c_n)^2 = R^2(t) + 2\mathbf{R}(t) \cdot \boldsymbol{\beta}_n (t - t_r)c_n + \beta_n^2 ((t - t_r)c_n)^2$$

$$(t - t_r)c_n = \frac{-\mathbf{R}(t) \cdot \boldsymbol{\beta}_n \pm \sqrt{(\mathbf{R}(t) \cdot \boldsymbol{\beta}_n)^2 - (\beta_n^2 - 1)R^2(t)}}{\beta_n^2 - 1}$$


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$$\mathbf{R}(t_r) = \mathbf{r} - \mathbf{v}t_r = \mathbf{R}(t) + \mathbf{v}(t - t_r)$$

$$(t - t_r)c_n = R(t_r)$$

$$R(t_r) - \mathbf{R}(t_r) \cdot \boldsymbol{\beta}_n = (t - t_r)c_n(1 - \beta_n^2) - \mathbf{R}(t) \cdot \boldsymbol{\beta}_n$$

$$R(t_r) = \frac{-\mathbf{R}(t) \cdot \boldsymbol{\beta}_n \pm \sqrt{(\mathbf{R}(t) \cdot \boldsymbol{\beta}_n)^2 - (\beta_n^2 - 1)R^2(t)}}{\beta_n^2 - 1}$$

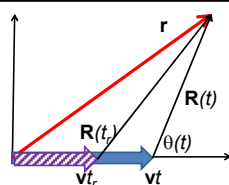
$$R(t_r) = \frac{R(t)}{\beta_n^2 - 1} \left(-\beta_n \cos \theta \pm \sqrt{1 - \beta_n^2 \sin^2 \theta} \right)$$

$$R(t_r) - \mathbf{R}(t_r) \cdot \boldsymbol{\beta}_n = \mp R(t) \sqrt{1 - \beta_n^2 \sin^2 \theta}$$

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Liénard-Wiechert potentials:

$$\Phi(\mathbf{r}, t) = \pm \frac{q}{\epsilon} \frac{1}{R(t) \sqrt{1 - \beta_n^2 \sin^2 \theta}}$$

$$\mathbf{A}(\mathbf{r}, t) = \pm q\boldsymbol{\mu} \frac{\boldsymbol{\beta}_n}{R(t) \sqrt{1 - \beta_n^2 \sin^2 \theta}}$$


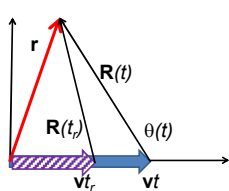
For $\beta_n > 1$, range of θ is limited:

$$R(t_r) = \frac{R(t)}{\beta_n^2 - 1} \left(-\beta_n \cos \theta \pm \sqrt{1 - \beta_n^2 \sin^2 \theta} \right) \geq 0$$

$$\Rightarrow \frac{\pi}{2} \leq \theta \leq \sin^{-1} \left(\frac{1}{\beta_n} \right)$$

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Physical fields for $\beta_n > 1$



$$\frac{\pi}{2} \leq \theta \leq \sin^{-1}\left(\frac{1}{\beta_n}\right)$$

Define $\cos \theta_c \equiv \sqrt{1 - \frac{1}{\beta_n^2}}$

$$\Rightarrow \cos \theta \leq \cos \theta_c$$

$$\Phi(\mathbf{r}, t) = \frac{2q}{\epsilon} \frac{1}{R(t)\sqrt{1 - \beta_n^2 \sin^2 \theta}} \Theta(\cos \theta_c - \cos \theta(t))$$

$$\mathbf{A}(\mathbf{r}, t) = 2q\mu \frac{\beta_n}{R(t)\sqrt{1 - \beta_n^2 \sin^2 \theta}} \Theta(\cos \theta_c - \cos \theta(t))$$

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Physical fields for $\beta > 1$

$$\Phi(\mathbf{r}, t) = \frac{2q}{\epsilon} \frac{1}{R(t)\sqrt{1 - \beta_n^2 \sin^2 \theta}} \Theta(\cos \theta_c - \cos \theta(t))$$

$$\mathbf{A}(\mathbf{r}, t) = 2q\mu \frac{\beta_n}{R(t)\sqrt{1 - \beta_n^2 \sin^2 \theta}} \Theta(\cos \theta_c - \cos \theta(t))$$

$$\mathbf{E}(\mathbf{r}, t) = -\nabla\Phi - \frac{1}{c_n} \frac{\partial \mathbf{A}}{\partial t} \quad \mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}$$

$$\mathbf{E}(\mathbf{r}, t) = \frac{2q}{\epsilon} \frac{\hat{\mathbf{R}}}{(R(t))^2 \sqrt{1 - \beta_n^2 \sin^2 \theta}} \times \left(\frac{\beta_n^2 - 1}{1 - \beta_n^2 \sin^2 \theta} \Theta(\cos \theta_c - \cos \theta(t)) + \sqrt{\beta_n^2 - 1} \delta(\cos \theta_c - \cos \theta(t)) \right)$$

$$\mathbf{B}(\mathbf{r}, t) = -\beta_n \sin \theta (\hat{\theta} \times \mathbf{E}(\mathbf{r}, t))$$

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Intermediate steps:

$$\frac{d\theta}{dt} = \frac{v \sin \theta}{R} \quad \frac{dR}{dt} = -v \cos \theta$$

Using instantaneous polar coordinates: $\nabla \equiv \hat{\mathbf{R}} \frac{\partial}{\partial R} + \hat{\boldsymbol{\theta}} \frac{1}{R} \frac{\partial}{\partial \theta}$

$$\nabla \Theta(\cos \theta_c - \cos \theta(t)) = \delta(\cos \theta_c - \cos \theta(t)) \frac{\sin \theta(t)}{R(t)} \hat{\boldsymbol{\theta}}$$

$$\frac{\partial \Theta(\cos \theta_c - \cos \theta(t))}{\partial t} = \delta(\cos \theta_c - \cos \theta(t)) \frac{v \sin^2 \theta(t)}{R(t)}$$

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Some details: Cherenkov radiation observed near the angle θ_c at time $t=t_c+\Delta t$

$$\cos \theta_c - \cos \theta(t) \approx \frac{c_n \Delta t}{\beta_n R_C}$$

$$1 - \beta_n^2 \sin^2 \theta(t) \approx \frac{2c_n \Delta t \sqrt{\beta_n^2 - 1}}{R_C}$$

When the dust clears

$$\frac{d^2 I}{d\omega d\ell} \propto \left(1 - \frac{c_n^2}{v^2}\right) \omega$$

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Some details: Cherenkov radiation observed near the angle θ_c

$$\mathbf{R}(t) = \mathbf{r} - \mathbf{vt}$$

$$\sin \theta_c = \frac{c_n}{v}$$

$$\pi \geq \theta(t) \geq \theta_c$$

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Cherenkov radiation observed near the angle θ_c -- continued

$$\cos \theta_c - \cos \theta(\Delta t) \approx \sin \theta_c \theta(\Delta t)$$

$$\approx \frac{c_n \Delta t}{R_C}$$

$$1 - \beta_n^2 \sin^2 \theta(t) \approx 2\sqrt{\beta_n^2 - 1} \frac{c_n \Delta t}{R_C}$$

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Cherenkov radiation observed near the angle θ_c -- continued

$$\mathbf{E}(\mathbf{r}, t) = \frac{2q}{\epsilon} \frac{\hat{\mathbf{R}}}{(R(t))^2 \sqrt{1 - \beta_n^2 \sin^2 \theta}} \times \left(\frac{\beta_n^2 - 1}{1 - \beta_n^2 \sin^2 \theta} \Theta(\cos \theta_c - \cos \theta(t)) + \sqrt{\beta_n^2 - 1} \delta(\cos \theta_c - \cos \theta(t)) \right)$$

$$\mathbf{B}(\mathbf{r}, t) = -\beta_n \sin \theta (\hat{\boldsymbol{\theta}} \times \mathbf{E}(\mathbf{r}, t))$$

Estimates at $t = t_c + \Delta t$

$$\mathbf{E}(\mathbf{r}, t) \approx -\frac{2q}{\epsilon} \hat{\mathbf{R}}_c \frac{(\beta_n^2 - 1)^{1/4}}{(2c_n^3 R_c)^{1/2}} \left[(\Delta t)^{-1/2} \delta(\Delta t) - \frac{1}{2} (\Delta t)^{-3/2} \Theta(\Delta t) \right]$$

$$\mathbf{B}(\mathbf{r}, t) = -(\hat{\boldsymbol{\theta}}(0) \times \mathbf{E}(\mathbf{r}, t))$$

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Cherenkov radiation observed near the angle θ_c -- continued

Spectral analysis:


$$\tilde{\mathbf{E}}(\omega) = -\frac{2q}{\epsilon} \hat{\mathbf{R}}_c \frac{(\beta_n^2 - 1)^{1/4}}{(2c_n^3 R_c)^{1/2}} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dt \left[t^{-1/2} \delta(t) - \frac{1}{2} t^{-3/2} \Theta(t) \right] e^{i\omega t}$$

$$= -i\omega \frac{2q}{\epsilon} \hat{\mathbf{R}}_c \frac{(\beta_n^2 - 1)^{1/4}}{(2c_n^3 R_c)^{1/2}} \frac{1}{\sqrt{2\pi}} \int_0^{\infty} dt t^{-1/2} e^{i\omega t}$$

$$= \frac{q}{\epsilon} \hat{\mathbf{R}}_c \frac{(\beta_n^2 - 1)^{1/4}}{(2c_n^3 R_c)^{1/2}} (1 - i) \sqrt{\omega}$$

Spectral intensity: $\frac{d^2 I}{d\Omega d\omega} \propto |\tilde{\mathbf{E}}(\omega)|^2 = \frac{q^2 (\beta_n^2 - 1)^{1/2}}{\epsilon^2 c_n^3 R_c} \omega$

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Cherenkov radiation emitted by the core of the Reed Research Reactor located at Reed College in Portland, Oregon, U.S.
Cherenkov radiation. Photograph. *Encyclopædia Britannica Online*. Web. 12 Apr. 2013.
<http://www.britannica.com/EBchecked/media/174732>

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