

PHY 770 -- Statistical Mechanics
12:30-1:45 PM TR Olin 107

Instructor: Natalie Holzwarth (Olin 300)
 Course Webpage: <http://www.wfu.edu/~natalie/s14phy770>

Lecture 5 -- Chapter 3
Review of Thermodynamics – continued

1. Continue the analysis of thermodynamic stability
2. Examples

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PHY 770 Statistical Mechanics

TR 12:30-1:45 PM OPL 107 <http://www.wfu.edu/~natalie/s14phy770/>

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Course schedule for Spring 2014

(Preliminary schedule – subject to frequent adjustment.) Please note that makeup lectures (indicated in red) are scheduled for Tuesdays or Thursdays at 11 AM - 12:15 PM in Olin 107.

Lecture date	Text Reading	Topic	Assign.	Due date
1 Tue: 01/14/2014	Chap. 3	Review of macroscopic thermodynamics	#1	02/04/2014
2 Thu: 01/16/2014	Chap. 3	Review of macroscopic thermodynamics	#2	02/04/2014
3 Tue: 01/21/2014	Chap. 3	Thermodynamic potentials	#3	02/04/2014
4 Thu: 01/23/2014	Chap. 3	Thermodynamic stability	#4	02/04/2014
5 Tue: 01/28/2014		IAWH out of town - no class		
6 Thu: 01/30/2014		IAWH out of town - no class		
7 Tue: 02/04/2014	Chap. 3			02/06/2014

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Summary of thermodynamic potentials

Potential	Variables	Total Diff	Fund. Eq.
U	S, X, N_i	$dU = TdS + YdX + \sum_i \mu_i dN_i$	$U = TS + YX + \sum_i \mu_i N_i$
H	S, Y, N_i	$dH = TdS - XdY + \sum_i \mu_i dN_i$	$H = U - YX$
A	T, X, N_i	$dA = -SdT + YdX + \sum_i \mu_i dN_i$	$A = U - TS$
G	T, Y, N_i	$dG = -SdT - XdY + \sum_i \mu_i dN_i$	$G = U - TS - YX$
Ω	T, X, μ_i	$dΩ = -SdT + YdX - \sum_i \mu_i dN_i$	$Ω = U - TS - \sum_i \mu_i N_i$

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Derivative relationships of thermodynamic potentials

Internal energy $U(S, X, \{N_i\})$:

$$T = \left(\frac{\partial U}{\partial S}\right)_{X, \{N_i\}} \quad Y = \left(\frac{\partial U}{\partial X}\right)_{S, \{N_i\}} \quad \mu_i = \left(\frac{\partial U}{\partial N_i}\right)_{S, X, \{N_j\}}$$

Enthalpy $H(S, Y, \{N_i\})$:

$$T = \left(\frac{\partial H}{\partial S}\right)_{Y, \{N_i\}} \quad X = \left(\frac{\partial H}{\partial Y}\right)_{S, \{N_i\}} \quad \mu_i = \left(\frac{\partial H}{\partial N_i}\right)_{S, Y, \{N_j\}}$$

Helmholz free energy $A(T, X, \{N_i\})$:

$$S = -\left(\frac{\partial A}{\partial T}\right)_{X, \{N_i\}} \quad Y = \left(\frac{\partial A}{\partial X}\right)_{T, \{N_i\}} \quad \mu_i = \left(\frac{\partial A}{\partial N_i}\right)_{T, X, \{N_j\}}$$

Gibb's free energy $G(T, Y, \{N_i\})$:

$$S = -\left(\frac{\partial G}{\partial T}\right)_{Y, \{N_i\}} \quad X = -\left(\frac{\partial G}{\partial Y}\right)_{T, \{N_i\}} \quad \mu_i = \left(\frac{\partial G}{\partial N_i}\right)_{T, Y, \{N_j\}}$$

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Equilibrium properties of thermodynamic potentials

At equilibrium $dS \rightarrow 0$ and $S \rightarrow$ maximum

At equilibrium $dU \rightarrow 0$ and $U(S, X, \{N_i\}) \rightarrow$ minimum

At equilibrium $dH \rightarrow 0$ and $H(S, Y, \{N_i\}) \rightarrow$ minimum

At equilibrium $dA \rightarrow 0$ and $A(T, X, \{N_i\}) \rightarrow$ minimum

At equilibrium $dG \rightarrow 0$ and $G(T, Y, \{N_i\}) \rightarrow$ minimum

At equilibrium $d\Omega \rightarrow 0$ and $\Omega(T, X, \{\mu_i\}) \rightarrow$ minimum

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Stability analysis of the equilibrium state

$P_A, V_A, T_A,$	$P_B, V_B, T_B,$
$\{N_{iA}\}, \{\mu_{iA}\}$	$\{N_{iB}\}, \{\mu_{iB}\}$

←

→

$$\Delta U_A = -\Delta U_B$$

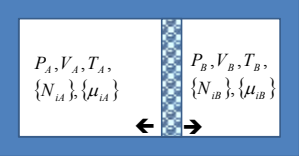
$$\Delta V_A = -\Delta V_B$$

$$\Delta N_{iA} = -\Delta N_{iB}$$

Assume that the total system is isolated, but that there can be exchange of variables A and B .

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Stability analysis of the equilibrium state



If we can assume :

$$\Delta U_A = -\Delta U_B$$

$$\Delta V_A = -\Delta V_B$$

$$\Delta N_{iA} = -\Delta N_{iB}$$

Analysis of entropy fluctuations :

$$\Delta S_{total} = \sum_{\alpha=A,B} \left[\left(\frac{\partial S_{\alpha}}{\partial U_{\alpha}} \right) \Delta U_{\alpha} + \left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right) \Delta V_{\alpha} + \sum_i \left(\frac{\partial S_{\alpha}}{\partial N_{i\alpha}} \right) \Delta N_{i\alpha} \right]$$

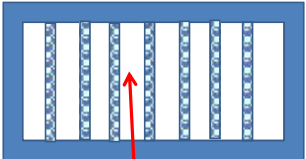
$$= \left(\frac{1}{T_A} - \frac{1}{T_B} \right) \Delta U_A + \left(\frac{P_A}{T_A} - \frac{P_B}{T_B} \right) \Delta V_A + \sum_i \left(\frac{\mu_{iA}}{T_A} - \frac{\mu_{iB}}{T_B} \right) \Delta N_{iA}$$

At equilibrium m : $T_A = T_B, P_A = P_B, \mu_{iA} = \mu_{iB}$

Note that this result does not hold for non-porous partition.

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Stability analysis of the equilibrium state – multi-partioned box



$P_{\alpha}, V_{\alpha}, T_{\alpha}, \{N_{i\alpha}\}, \{\mu_{i\alpha}\}$

Entropy fluctuations in partition α near equilibrium :

$$S_{\alpha}(U_{\alpha}, V_{\alpha}, \{N_{i\alpha}\}) = S_{\alpha}^0(U_{\alpha}^0, V_{\alpha}^0, \{N_{i\alpha}^0\}) +$$

+ (first order fluctuations)
+ (second order fluctuations) +

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Stability analysis of the equilibrium state – multi-partioned box

Entropy fluctuations in partition α near equilibrium :

$$S_{\alpha}(U_{\alpha}, V_{\alpha}, \{N_{i\alpha}\}) = S_{\alpha}^0(U_{\alpha}^0, V_{\alpha}^0, \{N_{i\alpha}^0\}) +$$

+ (first order fluctuations)
+ (second order fluctuations) +

(first order fluctuations) =

$$\left(\frac{\partial S_{\alpha}}{\partial U_{\alpha}} \right)^0 \Delta U_{\alpha} + \left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right)^0 \Delta V_{\alpha} + \sum_i \left(\frac{\partial S_{\alpha}}{\partial N_{i\alpha}} \right)^0 \Delta N_{i\alpha}$$

(second order fluctuations) =

$$\frac{1}{2} \left\{ \Delta \left(\frac{\partial S_{\alpha}}{\partial U_{\alpha}} \right) \Delta U_{\alpha} + \Delta \left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right) \Delta V_{\alpha} + \sum_i \Delta \left(\frac{\partial S_{\alpha}}{\partial N_{i\alpha}} \right) \Delta N_{i\alpha} \right\}$$

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Stability analysis of the equilibrium state – multi-partioned box
(second order fluctuations) =

$$\frac{1}{2} \left\{ \Delta \left(\frac{\partial S_\alpha}{\partial U_\alpha} \right) \Delta U_\alpha + \Delta \left(\frac{\partial S_\alpha}{\partial V_\alpha} \right) \Delta V_\alpha + \sum_i \Delta \left(\frac{\partial S_\alpha}{\partial N_{i\alpha}} \right) \Delta N_{i\alpha} \right\}$$

notation :

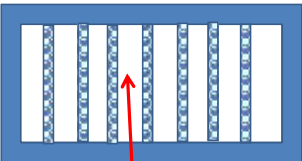
$$\Delta \left(\frac{\partial S_\alpha}{\partial U_\alpha} \right) \equiv \left(\frac{\partial^2 S_\alpha}{\partial U_\alpha^2} \right)^0 \Delta U_\alpha + \left(\frac{\partial^2 S_\alpha}{\partial U_\alpha \partial V_\alpha} \right)^0 \Delta V_\alpha + \sum_i \left(\frac{\partial^2 S_\alpha}{\partial N_{i\alpha} \partial U_\alpha} \right)^0 \Delta N_{i\alpha}$$

$$\Delta \left(\frac{\partial S_\alpha}{\partial V_\alpha} \right) \equiv \left(\frac{\partial^2 S_\alpha}{\partial V_\alpha^2} \right)^0 \Delta V_\alpha + \left(\frac{\partial^2 S_\alpha}{\partial U_\alpha \partial V_\alpha} \right)^0 \Delta U_\alpha + \sum_i \left(\frac{\partial^2 S_\alpha}{\partial N_{i\alpha} \partial V_\alpha} \right)^0 \Delta N_{i\alpha}$$

$$\Delta \left(\frac{\partial S_\alpha}{\partial N_{i\alpha}} \right) \equiv \left(\frac{\partial^2 S_\alpha}{\partial V_\alpha \partial N_{i\alpha}} \right)^0 \Delta V_\alpha + \left(\frac{\partial^2 S_\alpha}{\partial U_\alpha \partial N_{i\alpha}} \right)^0 \Delta U_\alpha + \sum_j \left(\frac{\partial^2 S_\alpha}{\partial N_{j\alpha} \partial N_{i\alpha}} \right)^0 \Delta N_{j\alpha}$$

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Stability analysis of the equilibrium state – multi-partioned box



$P_\alpha, V_\alpha, T_\alpha, \{N_{i\alpha}\}, \{\mu_{i\alpha}\}$

Entropy fluctuations total box near equilibrium m :

$\Delta S_{total} =$ (first order fluctuations) + (second order fluctuations) +

(second order fluctuations) =

$$\sum_\alpha \frac{1}{2} \left\{ \Delta \left(\frac{\partial S_\alpha}{\partial U_\alpha} \right) \Delta U_\alpha + \Delta \left(\frac{\partial S_\alpha}{\partial V_\alpha} \right) \Delta V_\alpha + \sum_i \Delta \left(\frac{\partial S_\alpha}{\partial N_{i\alpha}} \right) \Delta N_{i\alpha} \right\}$$

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Stability analysis of the equilibrium state – multi-partioned box

Entropy fluctuations total box near equilibrium m :

$\Delta S_{total} =$ (second order fluctuations) +

(second order fluctuations)

$$= \sum_\alpha \frac{1}{2} \left\{ \Delta \left(\frac{\partial S_\alpha}{\partial U_\alpha} \right) \Delta U_\alpha + \Delta \left(\frac{\partial S_\alpha}{\partial V_\alpha} \right) \Delta V_\alpha + \sum_i \Delta \left(\frac{\partial S_\alpha}{\partial N_{i\alpha}} \right) \Delta N_{i\alpha} \right\}$$

$$= \sum_\alpha \frac{1}{2} \left\{ \Delta \left(\frac{1}{T} \right)_\alpha \Delta U_\alpha + \Delta \left(\frac{P}{T} \right)_\alpha \Delta V_\alpha + \sum_i \Delta \left(\frac{\mu_i}{T} \right)_\alpha \Delta N_{i\alpha} \right\}$$

Noting that : $T\Delta S = \Delta U + P\Delta V - \sum_i \mu_i \Delta N_i$

$$\Rightarrow \Delta S_{total} = -\frac{1}{2T} \sum_\alpha \left\{ \Delta T_\alpha \Delta S_\alpha - \Delta P_\alpha \Delta V_\alpha + \sum_i \Delta \mu_{i\alpha} \Delta N_{i\alpha} \right\}$$

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Stability analysis of the equilibrium state
 Note that at equilibrium, $S_{total} = \text{maximum}$
 $\Rightarrow \Delta S_{total} \leq 0$

$$\Delta S_{total} = -\frac{1}{2T} \sum_{\alpha} \left\{ \underbrace{\Delta T_{\alpha} \Delta S_{\alpha} - \Delta P_{\alpha} \Delta V_{\alpha} + \sum_i \Delta \mu_{i\alpha} \Delta N_{i\alpha}}_{\geq 0} \right\}$$

After a few more steps --

$$\Delta S_{total} = -\frac{1}{2T} \sum_{\alpha} \left\{ \left(\frac{\partial S}{\partial T} \right)_{T,P,\{N_i\}}^0 \Delta T_{\alpha}^2 - \left(\frac{\partial P}{\partial V} \right)_{T,\{N_i\}}^0 \Delta V_{\alpha}^2 + \sum_{i,j} \left(\frac{\partial \mu_i}{\partial N_i} \right)_{T,P,\{N_i\}}^0 \Delta N_{i\alpha} \Delta N_{j\alpha} \right\}$$

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Stability analysis of the equilibrium state -- continued

Le Chatelier's principle: If a system is in stable equilibrium, then any spontaneous change in its parameters must bring about processes which tend to restore the system to equilibrium.

$$\Delta S_{total} = -\frac{1}{2T} \sum_{\alpha} \left\{ \left(\frac{\partial S}{\partial T} \right)_{T,P,\{N_i\}}^0 \Delta T_{\alpha}^2 - \left(\frac{\partial P}{\partial V} \right)_{T,\{N_i\}}^0 \Delta V_{\alpha}^2 + \sum_{i,j} \left(\frac{\partial \mu_i}{\partial N_i} \right)_{T,P,\{N_i\}}^0 \Delta N_{i\alpha} \Delta N_{j\alpha} \right\}$$

By assumption, the fluctuations of each variable is independent so that we can examine each term separately.

$$\left(\frac{\partial S}{\partial T} \right)_{T,P,\{N_i\}}^0 \geq 0 \Rightarrow C_{V,\{N_i\}} = T \left(\frac{\partial S}{\partial T} \right)_{T,P,\{N_i\}} \geq 0$$

$$\left(\frac{\partial P}{\partial V} \right)_{T,\{N_i\}}^0 \leq 0 \Rightarrow \kappa_{T,\{N_i\}} = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_{T,\{N_i\}} \geq 0$$

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Stability analysis of the equilibrium state -- continued

$$\Delta S_{total} = -\frac{1}{2T} \sum_{\alpha} \left\{ \left(\frac{\partial S}{\partial T} \right)_{T,P,\{N_i\}}^0 \Delta T_{\alpha}^2 - \left(\frac{\partial P}{\partial V} \right)_{T,\{N_i\}}^0 \Delta V_{\alpha}^2 + \sum_{i,j} \left(\frac{\partial \mu_i}{\partial N_i} \right)_{T,P,\{N_i\}}^0 \Delta N_{i\alpha} \Delta N_{j\alpha} \right\}$$

Contributions from fluctuations in the chemical potentials

$$\sum_{i,j} \left(\frac{\partial \mu_i}{\partial N_i} \right)_{T,P,\{N_i\}}^0 \Delta N_{i\alpha} \Delta N_{j\alpha} \geq 0$$

Define: $\mathbf{M} \equiv \begin{pmatrix} m_{11} & m_{12} & \dots & m_{1v} \\ m_{21} & m_{22} & \dots & m_{2v} \\ \vdots & \vdots & \dots & \vdots \\ m_{v1} & m_{v2} & \dots & m_{vv} \end{pmatrix} \Rightarrow \mathbf{M}$ must be positive definite

where: $m_{ij} \equiv \left(\frac{\partial \mu_i}{\partial N_j} \right)_{T,P,\{N_i\}} = \left(\frac{\partial \mu_j}{\partial N_i} \right)_{T,P,\{N_i\}}$

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Example 3.10 from Reichl:

Consider a system at constant P, T containing two kinds of particles A and B with the following form for the Gibbs free energy:

$$G(T, P, N_A, N_B) = N_A \mu_A^0 + N_B \mu_B^0 + kT N_A \ln \left(\frac{N_A}{N_A + N_B} \right) + kT N_B \ln \left(\frac{N_B}{N_A + N_B} \right) + \lambda \frac{N_A N_B}{N_A + N_B}$$

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