

**PHY 770 -- Statistical Mechanics**  
**10-10:50 AM MWF Olin 107**

Instructor: Natalie Holzwarth (Olin 300)  
 Course Webpage: <http://www.wfu.edu/~natalie/s14phy770>

**Lecture 2 -- Chapter 3**  
**Review of Thermodynamics – continued**

1. Some empirically obtained equations of state
2. Some properties of entropy
3. Thermodynamic potentials

1/16/2014

PHY 770 Spring 2014 -- Lecture 2

1

---

---

---

---

---

---

---

---

---

---

**Equations of state**

| Variable        | Units              |
|-----------------|--------------------|
| T (temperature) | °K                 |
| P (pressure)    | Pa                 |
| V (volume)      | m <sup>3</sup>     |
| n (moles)       | n=N/N <sub>A</sub> |
| N (particles)   | N=nN <sub>A</sub>  |

Avogadro's number: N<sub>A</sub>=6.022 141 29 x 10<sup>23</sup> mol<sup>-1</sup>

Boltzmann constant: k=1.380 6488 x 10<sup>-23</sup> J K<sup>-1</sup>

Molar gas constant: R=N<sub>A</sub>k=8.314 4621 J mol<sup>-1</sup> K<sup>-1</sup>

<http://physics.nist.gov/cuu/Constants/index.html>

1/16/2014

PHY 770 Spring 2014 -- Lecture 2

2

---

---

---

---

---

---

---

---

---

---

**Equations of State -- examples**

**Range of validity**

Ideal Gas Law

$$PV = nRT \equiv NkT$$

**dilute limit; ignore  
particle interactions**

Virial expansion

$$P = \left( \frac{nRT}{V} \right) \left( 1 + \frac{n}{V} B_2(T) + \frac{n^2}{V^2} B_3(T) + \dots \right)$$

**includes effects of  
higher density in terms  
of virial coefficients B<sub>i</sub>**

Van der Waals equation of state

$$\left( P + \frac{an^2}{V^2} \right) (V - nb) = nRT$$

**approximates gases and  
liquids in terms of an  
excluded volume nb and  
a cohesion parameter a**

1/16/2014

PHY 770 Spring 2014 -- Lecture 2

3

---

---

---

---

---

---

---

---

---

---

**Special properties of entropy**

For a reversible process:  $dS = \frac{dQ}{T}$

Thermo "laws" involving entropy

2. Heat flows spontaneously from high temperature to low temperature
3. It is not possible to reach the coldest temperature using a finite set of reversible steps

These relationships, together with the notion that entropy is an extensive and additive property leads to the **Fundamental equation of thermodynamics:**

$$TS = U - YX - \sum_i \mu_i N_i$$

1/16/2014 PHY 770 Spring 2014 – Lecture 2 4

---

---

---

---

---

---

---

---

---

---

---

---

**Fundamental equation of thermodynamics**

internal energy
generalized displacement (Y)
chemical potential

$$TS = U - YX - \sum_i \mu_i N_i$$

generalized force (-P)
number of particles

Derivation of fundamental equation of thermodynamics:

First law of thermodynamics for a reversible process :

$$dU = TdS + YdX + \sum_i \mu_i dN_i$$

$$dS = \frac{1}{T}dU - \frac{Y}{T}dX - \sum_i \frac{\mu_i}{T}dN_i$$

1/16/2014 PHY 770 Spring 2014 – Lecture 2 5

---

---

---

---

---

---

---

---

---

---

---

---

Derivation of fundamental equation of thermodynamics – continued:

$$dS = \frac{1}{T}dU - \frac{Y}{T}dX - \sum_i \frac{\mu_i}{T}dN_i$$

also :

$$dS = \left(\frac{\partial S}{\partial U}\right)_{X, N_i} dU + \left(\frac{\partial S}{\partial X}\right)_{U, N_i} dX + \sum_i \left(\frac{\partial S}{\partial N_i}\right)_{U, X, N_j} dN_i$$

$$\Rightarrow \left(\frac{\partial S}{\partial U}\right)_{X, N_i} = \frac{1}{T} \quad \left(\frac{\partial S}{\partial X}\right)_{U, N_i} = -\frac{Y}{T} \quad \left(\frac{\partial S}{\partial N_i}\right)_{U, X, N_j} = -\frac{\mu_i}{T}$$

$$\Rightarrow S = S(U, X, \{N_i\})$$

Also note because all of these quantities are extensive we can infer the following scaling relationship :

$$S(\lambda U, \lambda X, \{\lambda N_i\}) = \lambda S(U, X, \{N_i\})$$

1/16/2014 PHY 770 Spring 2014 – Lecture 2 6

---

---

---

---

---

---

---

---

---

---

---

---

Derivation of fundamental equation of thermodynamics – continued:

$$S(\lambda U, \lambda X, \{\lambda N_i\}) = \lambda S(U, X, \{N_i\})$$

$$\frac{d(\lambda S)}{d\lambda} = \left(\frac{\partial S}{\partial(\lambda U)}\right)_{X, N_i} \frac{d(\lambda U)}{d\lambda} + \left(\frac{\partial S}{\partial(\lambda X)}\right)_{U, N_i} \frac{d(\lambda X)}{d\lambda} + \sum_i \left(\frac{\partial S}{\partial(\lambda N_i)}\right)_{U, X, N_j} \frac{d(\lambda N_i)}{d\lambda}$$

$$\Rightarrow S = \left(\frac{\partial S}{\partial(\lambda U)}\right)_{X, N_i} U + \left(\frac{\partial S}{\partial(\lambda X)}\right)_{U, N_i} X + \sum_i \left(\frac{\partial S}{\partial(\lambda N_i)}\right)_{U, X, N_j} N_i$$

$$\Rightarrow S = \frac{1}{T} U - \frac{Y}{T} X - \sum_i \frac{\mu_i}{T} N_i$$

$$\Rightarrow TS = U - YX - \sum_i \mu_i N_i$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

7

---

---

---

---

---

---

---

---

---

---

**Fundamental equation of thermodynamics**

$$TS = U - YX - \sum_i \mu_i N_i$$

Some consequences:

From first law of thermodynamics :

$$dU = TdS + YdX + \sum_i \mu_i dN_i$$

$$d(TS) = d\left(U - YX - \sum_i \mu_i N_i\right)$$

$\Rightarrow$  Gibbs - Duhem equation :

$$SdT + XdY + \sum_i N_i d\mu_i = 0$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

8

---

---

---

---

---

---

---

---

---

---

**Thermodynamic potentials**

Internal energy :  $U = U(S, X, \{N_i\})$

Differential :  $dU = TdS + YdX + \sum_i \mu_i dN_i$

Fundamental equation :  $U = TS + YX + \sum_i \mu_i N_i$

Further relationships :

$$dU = \left(\frac{\partial U}{\partial S}\right)_{X, N_i} dS + \left(\frac{\partial U}{\partial X}\right)_{S, N_i} dX + \sum_i \left(\frac{\partial U}{\partial N_i}\right)_{S, X, N_j} dN_i$$

Equations of state :

$$\left(\frac{\partial U}{\partial S}\right)_{X, N_i} = T \quad \left(\frac{\partial U}{\partial X}\right)_{S, N_i} = Y \quad \left(\frac{\partial U}{\partial N_i}\right)_{S, X, N_j} = \mu_i$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

9

---

---

---

---

---

---

---

---

---

---

Analysis of internal energy continued:

Further relationships:

Equations of state:

$$\left(\frac{\partial U}{\partial S}\right)_{X,N_i} = T \quad \left(\frac{\partial U}{\partial X}\right)_{S,N_i} = Y \quad \left(\frac{\partial U}{\partial N_i}\right)_{S,X,N_j} = \mu_i$$

Maxwell's relations:

$$\left(\frac{\partial}{\partial X} \left(\frac{\partial U}{\partial S}\right)_{X,N_i}\right)_{S,N_i} = \left(\frac{\partial T}{\partial X}\right)_{S,N_i} = \left(\frac{\partial Y}{\partial S}\right)_{X,N_i}$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

10

---

---

---

---

---

---

---

---

---

---

Mathematical transformations for continuous functions of several variables & Legendre transforms:

$$z(x, y) \Leftrightarrow x(y, z)???$$

$$z(x, y) \Rightarrow dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy$$

$$x(y, z) \Rightarrow dx = \left(\frac{\partial x}{\partial y}\right)_z dy + \left(\frac{\partial x}{\partial z}\right)_y dz$$

$$\text{But: } \left(\frac{\partial x}{\partial y}\right)_z = -\frac{(\partial z / \partial y)_x}{(\partial z / \partial x)_y}$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

11

---

---

---

---

---

---

---

---

---

---

Mathematical transformations for continuous functions of several variables & Legendre transforms continued:

$$z(x, y) \Rightarrow dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy$$

$$\text{Let } u \equiv \left(\frac{\partial z}{\partial x}\right)_y \quad \text{and} \quad v \equiv \left(\frac{\partial z}{\partial y}\right)_x$$

Define new function

$$w(u, y) \Rightarrow dw = \left(\frac{\partial w}{\partial u}\right)_y du + \left(\frac{\partial w}{\partial y}\right)_u dy$$

$$\text{For } w = z - ux, \quad dw = dz - udx - xdu = udx + vdy - udx - xdu$$

$$dw = -xdu + vdy \quad \Rightarrow \left(\frac{\partial w}{\partial u}\right)_y = -x \quad \left(\frac{\partial w}{\partial y}\right)_u = \left(\frac{\partial z}{\partial y}\right)_x = v$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

12

---

---

---

---

---

---

---

---

---

---

For thermodynamic functions:

Internal energy:  $U = U(S, V)$

$$dU = TdS - PdV$$

$$dU = \left(\frac{\partial U}{\partial S}\right)_V dS + \left(\frac{\partial U}{\partial V}\right)_S dV$$

$$\Rightarrow T = \left(\frac{\partial U}{\partial S}\right)_V \quad P = -\left(\frac{\partial U}{\partial V}\right)_S$$

Enthalpy:  $H = H(S, P) = U + PV$

$$dH = dU + PdV + VdP = TdS + VdP = \left(\frac{\partial H}{\partial S}\right)_P dS + \left(\frac{\partial H}{\partial P}\right)_S dP$$

$$\Rightarrow T = \left(\frac{\partial H}{\partial S}\right)_P \quad V = \left(\frac{\partial H}{\partial P}\right)_S$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

13

---

---

---

---

---

---

---

---

---

---

Thermodynamic potentials – Internal Energy

| Internal energy      | $U(S, X, \{N_j\})$  |
|----------------------|---|
| Total Differential   | $dU = T dS + Y dX + \sum_{j=1}^r \mu_j dN_j$  |
| Fundamental Equation | $U = TS + YX + \sum_{j=1}^r \mu_j N_j$  |
| Equations of State   | $T = \left(\frac{\partial U}{\partial S}\right)_{X, \{N_j\}}$   |
|                      | $Y = \left(\frac{\partial U}{\partial X}\right)_{S, \{N_j\}}$   |
|                      | $\mu_j = \left(\frac{\partial U}{\partial N_j}\right)_{S, X, \{N_{i \neq j}\}}$   |
| Maxwell Relations    |   |
|                      | $\left(\frac{\partial T}{\partial X}\right)_{S, \{N_j\}} = \left(\frac{\partial Y}{\partial S}\right)_{X, \{N_j\}}$                                     |
|                      | $\left(\frac{\partial T}{\partial N_j}\right)_{S, X, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial S}\right)_{X, \{N_j\}}$                   |
|                      | $\left(\frac{\partial Y}{\partial N_j}\right)_{S, X, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial X}\right)_{S, \{N_j\}}$                   |
|                      | $\left(\frac{\partial \mu_j}{\partial N_i}\right)_{S, X, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_i}{\partial N_j}\right)_{S, X, \{N_{i \neq j}\}}$ |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

14

---

---

---

---

---

---

---

---

---

---

Thermodynamic potentials – Enthalpy

| Enthalpy             | $H(S, Y, \{N_j\})$  |
|----------------------|---|
| Total Differential   | $dH = T dS - X dY + \sum_{j=1}^r \mu_j dN_j$  |
| Fundamental Equation | $H = U - XY = TS + \sum_{j=1}^r \mu_j N_j$  |
| Equations of State   | $T = \left(\frac{\partial H}{\partial S}\right)_{Y, \{N_j\}}$   |
|                      | $X = -\left(\frac{\partial H}{\partial Y}\right)_{S, \{N_j\}}$  |
|                      | $\mu_j = \left(\frac{\partial H}{\partial N_j}\right)_{S, Y, \{N_{i \neq j}\}}$   |
| Maxwell Relations    |   |
|                      | $\left(\frac{\partial T}{\partial Y}\right)_{S, \{N_j\}} = -\left(\frac{\partial X}{\partial S}\right)_{Y, \{N_j\}}$                                    |
|                      | $\left(\frac{\partial T}{\partial N_j}\right)_{S, Y, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial S}\right)_{Y, \{N_j\}}$                   |
|                      | $\left(\frac{\partial X}{\partial N_j}\right)_{S, Y, \{N_{i \neq j}\}} = -\left(\frac{\partial \mu_j}{\partial Y}\right)_{S, \{N_j\}}$                  |
|                      | $\left(\frac{\partial \mu_j}{\partial N_i}\right)_{S, Y, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_i}{\partial N_j}\right)_{S, Y, \{N_{i \neq j}\}}$ |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

15

---

---

---

---

---

---

---

---

---

---

Thermodynamic potentials – Helmholtz Free Energy

| Helmholtz free energy | $A(T, X, \{N_j\})$  |
|-----------------------|---|
| Total Differential    | $dA = -S dT + Y dX + \sum_{j=1}^r \mu_j dN_j$   |
| Fundamental Equation  | $A = U - ST = XY + \sum_{j=1}^r \mu_j N_j$  |
| Equations of State    | $S = -\left(\frac{\partial A}{\partial T}\right)_{X, \{N_j\}}$  |
|                       | $Y = \left(\frac{\partial A}{\partial X}\right)_{T, \{N_j\}}$   |
|                       | $\mu_j = \left(\frac{\partial A}{\partial N_j}\right)_{T, X, \{N_{i \neq j}\}}$   |
| Maxwell Relations     |   |
|                       | $\left(\frac{\partial S}{\partial X}\right)_{T, \{N_j\}} = -\left(\frac{\partial Y}{\partial T}\right)_{X, \{N_j\}}$                                    |
|                       | $\left(\frac{\partial S}{\partial N_j}\right)_{T, X, \{N_{i \neq j}\}} = -\left(\frac{\partial \mu_j}{\partial T}\right)_{X, \{N_j\}}$                  |
|                       | $\left(\frac{\partial Y}{\partial N_j}\right)_{T, X, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial X}\right)_{T, X, \{N_{i \neq j}\}}$       |
|                       | $\left(\frac{\partial \mu_i}{\partial N_j}\right)_{T, X, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial N_i}\right)_{T, X, \{N_{i \neq j}\}}$ |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

16

---

---

---

---

---

---

---

---

---

---

---

---

Thermodynamic potentials – Gibbs Free Energy

| Gibbs free energy    | $G(T, Y, \{N_j\})$  |
|----------------------|---|
| Total Differential   | $dG = -S dT - X dY + \sum_{j=1}^r \mu_j dN_j$   |
| Fundamental Equation | $G = U - TS - YX = \sum_{j=1}^r \mu_j N_j$  |
| Equations of State   | $S = -\left(\frac{\partial G}{\partial T}\right)_{Y, \{N_j\}}$  |
|                      | $X = -\left(\frac{\partial G}{\partial Y}\right)_{T, \{N_j\}}$  |
|                      | $\mu_j = \left(\frac{\partial G}{\partial N_j}\right)_{T, Y, \{N_{i \neq j}\}}$   |
| Maxwell Relations    |   |
|                      | $\left(\frac{\partial S}{\partial Y}\right)_{T, \{N_j\}} = \left(\frac{\partial X}{\partial T}\right)_{Y, \{N_j\}}$                                     |
|                      | $\left(\frac{\partial S}{\partial N_j}\right)_{T, Y, \{N_{i \neq j}\}} = -\left(\frac{\partial \mu_j}{\partial T}\right)_{Y, \{N_j\}}$                  |
|                      | $\left(\frac{\partial X}{\partial N_j}\right)_{T, Y, \{N_{i \neq j}\}} = -\left(\frac{\partial \mu_j}{\partial Y}\right)_{T, \{N_j\}}$                  |
|                      | $\left(\frac{\partial \mu_i}{\partial N_j}\right)_{T, Y, \{N_{i \neq j}\}} = \left(\frac{\partial \mu_j}{\partial N_i}\right)_{T, Y, \{N_{i \neq j}\}}$ |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

17

---

---

---

---

---

---

---

---

---

---

---

---

Thermodynamic potentials – Grand Potential

| Grand potential      | $\Omega(T, X, \{\mu_j\})$   |
|----------------------|---|
| Total Differential   | $d\Omega = -S dT + Y dX - \sum_{j=1}^r N_j d\mu_j$  |
| Fundamental Equation | $\Omega = U - TS - \sum_{j=1}^r \mu_j N_j = XY$   |
| Equations of State   | $S = -\left(\frac{\partial \Omega}{\partial T}\right)_{X, \{\mu_j\}}$   |
|                      | $Y = \left(\frac{\partial \Omega}{\partial X}\right)_{T, \{\mu_j\}}$  |
|                      | $N_j = -\left(\frac{\partial \Omega}{\partial \mu_j}\right)_{T, X, \{\mu_{i \neq j}\}}$   |
| Maxwell Relations    |   |
|                      | $\left(\frac{\partial S}{\partial X}\right)_{T, \{\mu_j\}} = -\left(\frac{\partial Y}{\partial T}\right)_{X, \{\mu_j\}}$                                    |
|                      | $\left(\frac{\partial S}{\partial \mu_j}\right)_{T, Y, \{\mu_{i \neq j}\}} = \left(\frac{\partial N_j}{\partial T}\right)_{X, \{\mu_j\}}$                   |
|                      | $\left(\frac{\partial Y}{\partial \mu_j}\right)_{T, X, \{\mu_{i \neq j}\}} = -\left(\frac{\partial N_j}{\partial X}\right)_{T, \{\mu_j\}}$                  |
|                      | $\left(\frac{\partial N_i}{\partial \mu_j}\right)_{T, X, \{\mu_{i \neq j}\}} = \left(\frac{\partial N_j}{\partial \mu_i}\right)_{T, X, \{\mu_{i \neq j}\}}$ |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

18

---

---

---

---

---

---

---

---

---

---

---

---

## Summary of thermodynamic potentials

| Potential | Variables                  | Total Diff                            | Fund. Eq.                        |
|-----------|----------------------------|---------------------------------------|----------------------------------|
| <b>U</b>  | <b>S, X, N<sub>i</sub></b> | $dU = TdS + YdX + \sum_i \mu_i dN_i$  | $U = TS + YX + \sum_i \mu_i N_i$ |
| <b>H</b>  | <b>S, Y, N<sub>i</sub></b> | $dH = TdS - XdY + \sum_i \mu_i dN_i$  | $H = U - YX$                     |
| <b>A</b>  | <b>T, X, N<sub>i</sub></b> | $dA = -SdT + YdX + \sum_i \mu_i dN_i$ | $A = U - TS$                     |
| <b>G</b>  | <b>T, Y, N<sub>i</sub></b> | $dG = -SdT - XdY + \sum_i \mu_i dN_i$ | $G = U - TS - YX$                |
| <b>Ω</b>  | <b>T, X, μ<sub>i</sub></b> | $dΩ = -SdT + YdX - \sum_i \mu_i dN_i$ | $Ω = U - TS - \sum_i \mu_i N_i$  |

1/16/2014

PHY 770 Spring 2014 – Lecture 2

19

---

---

---

---

---

---

---

---

---

---

## Some examples

Ideal gas with  $N$  particles of the same type and with  
 $X \rightarrow V$                        $Y \rightarrow -P$

$$k \equiv k_B \text{ Boltzmann constant; } \gamma \equiv C_p / C_v$$

$$PV = NkT$$

$$U = \frac{Nk}{\gamma-1} T = \frac{PV}{\gamma-1}$$

$$H = U - XY = U + PV$$

$$\Rightarrow H = \frac{Nk\gamma}{\gamma-1} T$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

20

---

---

---

---

---

---

---

---

---

---

## Some examples -- continued

Ideal gas with  $N$  particles of the same type and with  
 $X \rightarrow V$                        $Y \rightarrow -P$

$$k \equiv k_B \text{ Boltzmann constant; } \gamma \equiv C_p / C_v$$

$$PV = NkT$$

$$U = \frac{Nk}{\gamma-1} T = \frac{PV}{\gamma-1}$$

$$A = U - TS$$

Previously, we have shown :

$$S(T, V) = \frac{Nk}{\gamma-1} \left( \ln \left( \frac{TV^{\gamma-1}}{T_0 V_0^{\gamma-1}} \right) \right) + S_0$$

1/16/2014

PHY 770 Spring 2014 – Lecture 2

21

---

---

---

---

---

---

---

---

---

---

Some examples -- continued

Ideal gas with  $N$  particles of the same type and with

$X \rightarrow V$        $Y \rightarrow P$

$$U = \frac{Nk}{\gamma-1}T \quad A = U - TS$$

$$\text{with } S(T, V) = \frac{Nk}{\gamma-1} \left( \ln \left( \frac{TV^{\gamma-1}}{T_0 V_0^{\gamma-1}} \right) \right) + S_0$$

$$A = \frac{Nk}{\gamma-1}T \left( 1 - \ln \left( \frac{TV^{\gamma-1}}{T_0 V_0^{\gamma-1}} \right) \right) - S_0 T$$

$$= -NkT \left( 1 + \ln \left( \frac{\frac{1}{T^{\gamma-1}V}}{T_0^{\frac{1}{\gamma-1}V_0}} \right) \right) - \left( S_0 - \frac{\gamma Nk}{\gamma-1} \right) T$$

1/16/2014

PHY 770 Spring 2014 -- Lecture 2

22

---



---



---



---



---



---



---



---