

PHY 712 Electrodynamics

11-11:50 AM MWF Olin 107

Plan for Lecture 30:

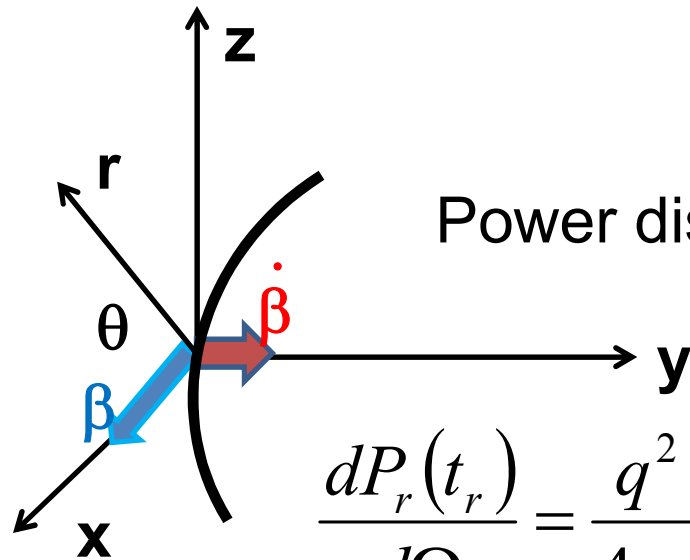
Continue reading Chap. 14 – Synchrotron radiation

- 1. Radiation from electron synchrotron devices**
- 2. Radiation from astronomical objects in circular orbits**

	03-22(Fri)	<i>APS Meeting</i>	(no class)	Exam
25	03-25(Mon)	Chap. 11	Lorentz transformations	#17
26	03-27(Wed)	Chap. 11	Transformations between electromagnetic fields	#18
27	03-28(Thur)	Chap. 11	Liénard-Wiechert potentials revisited	
	03-29(Fri)	<i>Good Friday</i>	(no class)	
28	04-01(Mon)	Chap. 14	Radiation by accelerated charges	#19
29	04-03(Wed)	Chap. 14	Radiation by accelerated charges	#20
30	04-05(Fri)	Chap. 14	Synchrotron radiation spectrum	#21
31	04-08(Mon)			
32	04-10(Wed)			
33	04-12(Fri)			
	04-15(Mon)		(no class -- presentation preparation)	
	04-17(Wed)		(no class -- presentation preparation)	
	04-19(Fri)		(no class -- presentation preparation)	
34	04-22(Mon)			
35	04-24(Wed)			
36	04-26(Fri)			
	04-29(Mon)		Student presentations I	
	05-01(Wed)		Student presentations II	
	05-02(Thurs)		Student presentations III	



Radiation from charged particle in circular path



Power distribution for circular acceleration

$$\frac{dP_r(t_r)}{d\Omega} = \frac{q^2}{4\pi c} \left| \frac{\hat{\mathbf{R}} \times [(\hat{\mathbf{R}} - \boldsymbol{\beta}) \times \dot{\boldsymbol{\beta}}]^2}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{R}})^5} \right|_{t_r = t - R/c}$$

$$= \frac{q^2}{4\pi c} \left| \frac{|\dot{\boldsymbol{\beta}}|^2 (1 - \boldsymbol{\beta} \cdot \hat{\mathbf{R}})^2 - (\hat{\mathbf{R}} \cdot \dot{\boldsymbol{\beta}})^2 (1 - \beta^2)}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{R}})^5} \right|_{t_r = t - R/c}$$

$$P_r(t_r) = \int d\Omega \frac{dP_r(t_r)}{d\Omega} = \frac{2}{3} \frac{q^2}{c^3} |\dot{\mathbf{v}}|^2 \gamma^4$$

Spectral composition of electromagnetic radiation -- continued

When the dust clears, the spectral intensity depends on the following integral :

$$\frac{\partial^2 I}{\partial \omega \partial \Omega} = \frac{q^2 \omega^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt_r [\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \times \boldsymbol{\beta}(t_r))] e^{i\omega(t_r - \hat{\mathbf{r}} \cdot \mathbf{R}_q(t_r)/c)} \right|^2$$

Synchrotron radiation light source installations

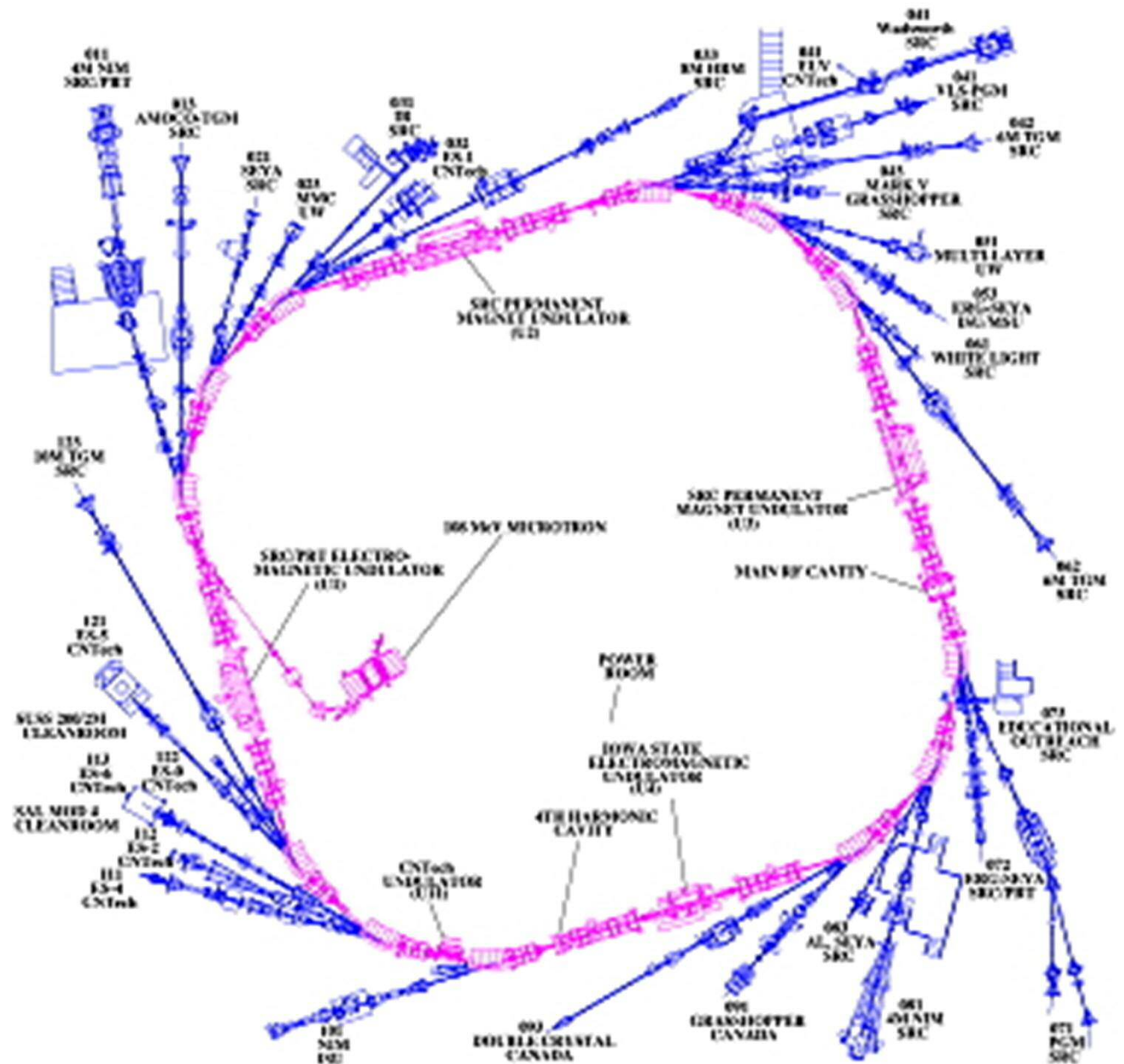
Synchrotron radiation center in Madison, Wisconsin



$E_c = 0.5 \text{ GeV}$ and 1 GeV ; $\lambda_c = 20 \text{ \AA}$ and 10 \AA

<http://www.src.wisc.edu/>

SRC –
 “Aladdin” --
 Madison
 Wisconsin

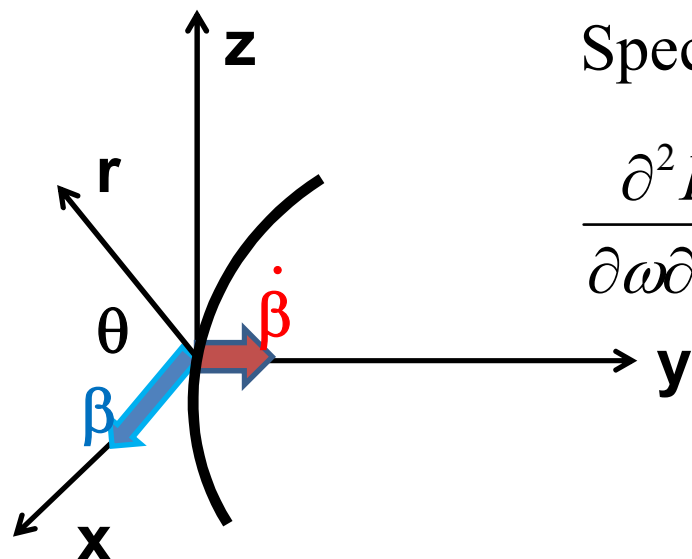


Brookhaven National Laboratory – National Light Source



$$E_c = 0.6 \text{ GeV}; \quad \lambda_c = 20 \text{ \AA}$$

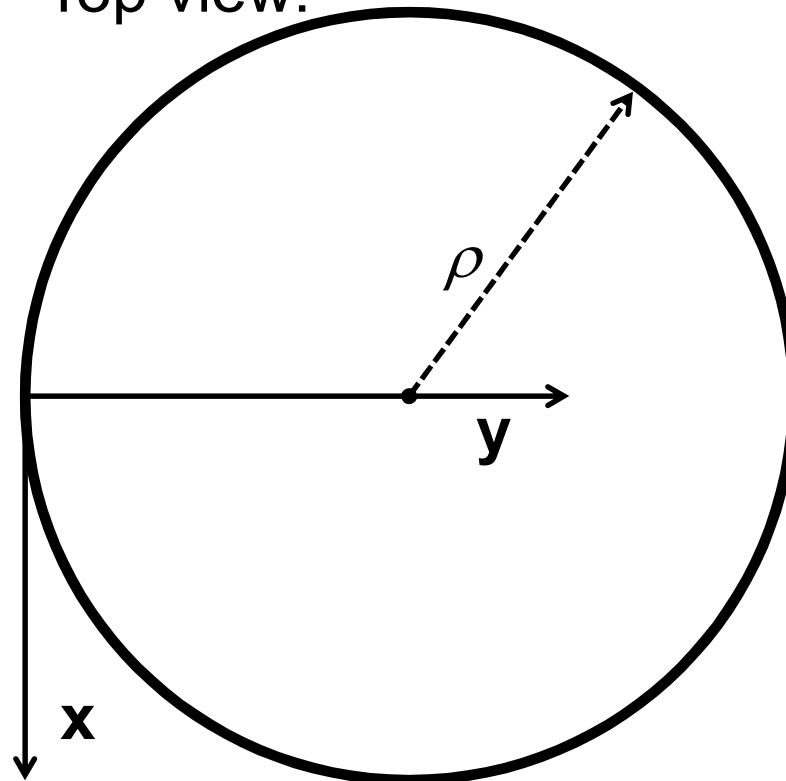
<http://www.bnl.gov/ps/>



Spectral intensity relationship :

$$\frac{\partial^2 I}{\partial \omega \partial \Omega} = \frac{q^2 \omega^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt_r [\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \times \boldsymbol{\beta}(t_r))] e^{i\omega(t_r - \hat{\mathbf{r}} \cdot \mathbf{R}_q(t_r)/c)} \right|^2$$

Top view:



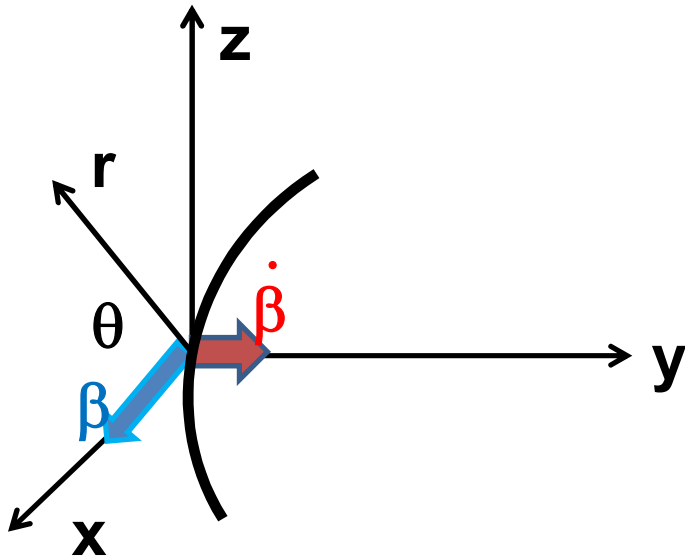
$$\mathbf{R}_q(t_r) = \rho \hat{\mathbf{x}} \sin(vt_r / \rho)$$

$$+ \rho \hat{\mathbf{y}} (1 - \cos(vt_r / \rho))$$

$$\boldsymbol{\beta}(t_r) = \beta (\hat{\mathbf{x}} \cos(vt_r / \rho) + \hat{\mathbf{y}} \sin(vt_r / \rho))$$

For convenience, choose:

$$\hat{\mathbf{r}} = \hat{\mathbf{x}} \cos \theta + \hat{\mathbf{z}} \sin \theta$$



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For convenience, choose:

$$\hat{\mathbf{r}} = \hat{\mathbf{x}} \cos \theta + \hat{\mathbf{z}} \sin \theta$$

Note that we have previously shown that in the radiation zone, the Poynting vector is in the $\hat{\mathbf{r}}$ direction; we can then choose to analyze two orthogonal polarization directions:

$$\boldsymbol{\varepsilon}_{\parallel} = \hat{\mathbf{y}} \quad \boldsymbol{\varepsilon}_{\perp} = -\hat{\mathbf{x}} \sin \theta + \hat{\mathbf{z}} \cos \theta$$

$$\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \times \boldsymbol{\beta}) = \beta (-\boldsymbol{\varepsilon}_{\parallel} \sin(vt_r / \rho) + \boldsymbol{\varepsilon}_{\perp} \sin \theta \cos(vt_r / \rho))$$

Modified Bessel functions

$$K_{1/3}(\xi) = \sqrt{3} \int_0^{\infty} dx \cos\left[\frac{3}{2} \xi \left(x + \frac{1}{3} x^3\right)\right] \quad K_{2/3}(\xi) = \sqrt{3} \int_0^{\infty} dx x \sin\left[\frac{3}{2} \xi \left(x + \frac{1}{3} x^3\right)\right]$$

Exponential factor

$$\omega(t_r - \hat{\mathbf{r}} \cdot \mathbf{R}_q(t_r)) = \omega\left(t_r - \frac{\rho}{c} \cos\theta \sin(vt_r / \rho)\right)$$

In the limit of $t_r \approx 0$, $\theta \approx 0$, $v \approx c\left(1 - \frac{1}{2\gamma^2}\right)$

$$\omega(t_r - \hat{\mathbf{r}} \cdot \mathbf{R}_q(t_r)) \approx \frac{\omega t_r}{2\gamma^2} (1 + \gamma^2 \theta^2) + \frac{\omega c^2 t_r^3}{6\rho^2} = \frac{3}{2} \xi \left(x + \frac{1}{3} x^3\right)$$

$$\text{where } \xi = \frac{\omega\rho}{3c\gamma^3} (1 + \gamma^2 \theta^2)^{3/2} \quad \text{and } x = \frac{c\gamma t_r}{\rho(1 + \gamma^2 \theta^2)^{1/2}}$$