


PHY 712 Electrodynamics

11-11:50 AM MWF Olin 107

Plan for Lecture 19:

Continue reading Chapter 7

- A. Reflectance and transmittance of electromagnetic waves – extension to anisotropy and complexity**
- B. Frequency dependence of dielectric materials**

11	02-11(Mon)	Chap. 5	Magnetostatics	#10
12	02-13(Wed)	Chap. 5	Magnetostatic fields	
13	02-15(Fri)	Chap. 5	Magnetic dipole fields	Exam
14	02-18(Mon)	Chap. 5	Permeable media	Exam
15	02-20(Wed)	Chap. 5	Magnetic susceptibility and permeability	Exam
16	02-22(Fri)	Chap. 6	Maxwell's equations	Exam
17	02-25(Mon)	Chap. 6	Poynting Vector	#11
18	02-27(Wed)	Chap. 7	Reflectance and transmittance of electromagnetic plane waves	#12
 19	02-28(Thur)	Chap. 7	Anisotropic media	#13
20	03-01(Fri)	Chap. 7	Dielectric models; Kramers-Kronig Relations	
21	03-04(Mon)			
22	03-06(Wed)			
23	03-07(Thur)			
24	03-08(Fri)			
	03-11(Mon)	<i>Spring Break</i>		
	03-13(Mon)	<i>Spring Break</i>		
	03-15(Fri)	<i>Spring Break</i>		
	03-18(Mon)	<i>APS Meeting</i>		Exam
	03-20(Wed)	<i>APS Meeting</i>		Exam
	03-22(Fri)	<i>APS Meeting</i>		Exam

Electromagnetic plane waves in isotropic medium with real permeability and permittivity: $\mu \epsilon$.

$$\mathbf{E}(\mathbf{r}, t) = \Re\left(\mathbf{E}_0 e^{i\frac{\omega}{c}(n\hat{\mathbf{k}}\cdot\mathbf{r}-ct)}\right)$$

$$\mathbf{B}(\mathbf{r}, t) = \frac{n}{c} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t) = \sqrt{\mu\epsilon} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}, t)$$

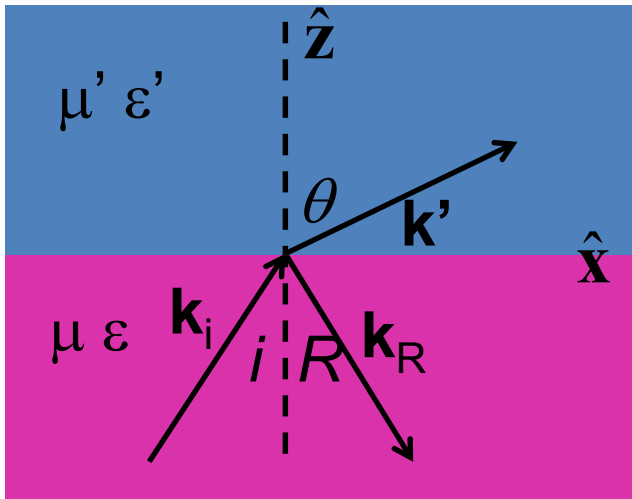
Poynting vector for plane electromagnetic waves:

$$\langle \mathbf{S} \rangle_{avg} = \frac{n|\mathbf{E}_0|^2}{2\mu c} \hat{\mathbf{k}} = \frac{1}{2} \sqrt{\frac{\epsilon}{\mu}} |\mathbf{E}_0|^2 \hat{\mathbf{k}}$$

Energy density for plane electromagnetic waves:

$$\langle u \rangle_{avg} = \frac{1}{2} \epsilon |\mathbf{E}_0|^2$$

Reflection and refraction between two isotropic media



Reflectance, transmittance:

$$R = \frac{\mathbf{S}_R \cdot \hat{\mathbf{z}}}{\mathbf{S}_i \cdot \hat{\mathbf{z}}} = \left| \frac{E_{0R}}{E_{0i}} \right|^2 \quad T = \frac{\mathbf{S}' \cdot \hat{\mathbf{z}}}{\mathbf{S}_i \cdot \hat{\mathbf{z}}} = \left| \frac{E'_0}{E_{0i}} \right|^2 \frac{n' \mu \cos \theta}{n \mu' \cos i}$$

Note that $R + T = 1$

For s-polarization

$$\frac{E_{0R}}{E_{0i}} = \frac{n \cos i - \frac{\mu}{\mu'} n' \cos \theta}{n \cos i + \frac{\mu}{\mu'} n' \cos \theta} \quad \frac{E'_0}{E_{0i}} = \frac{2n \cos i}{n \cos i + \frac{\mu}{\mu'} n' \cos \theta}$$

Note that: $n' \cos \theta = \sqrt{n'^2 - n^2 \sin^2 i}$

For p-polarization

$$\frac{E_{0R}}{E_{0i}} = \frac{\frac{\mu}{\mu'} n' \cos i - n \cos \theta}{\frac{\mu}{\mu'} n' \cos i + n \cos \theta} \quad \frac{E'_0}{E_{0i}} = \frac{2n \cos i}{\frac{\mu}{\mu'} n' \cos i + n \cos \theta}$$

Note that: $n' \cos \theta = \sqrt{n'^2 - n^2 \sin^2 i}$

Special case: normal incidence ($i=0, \theta=0$)

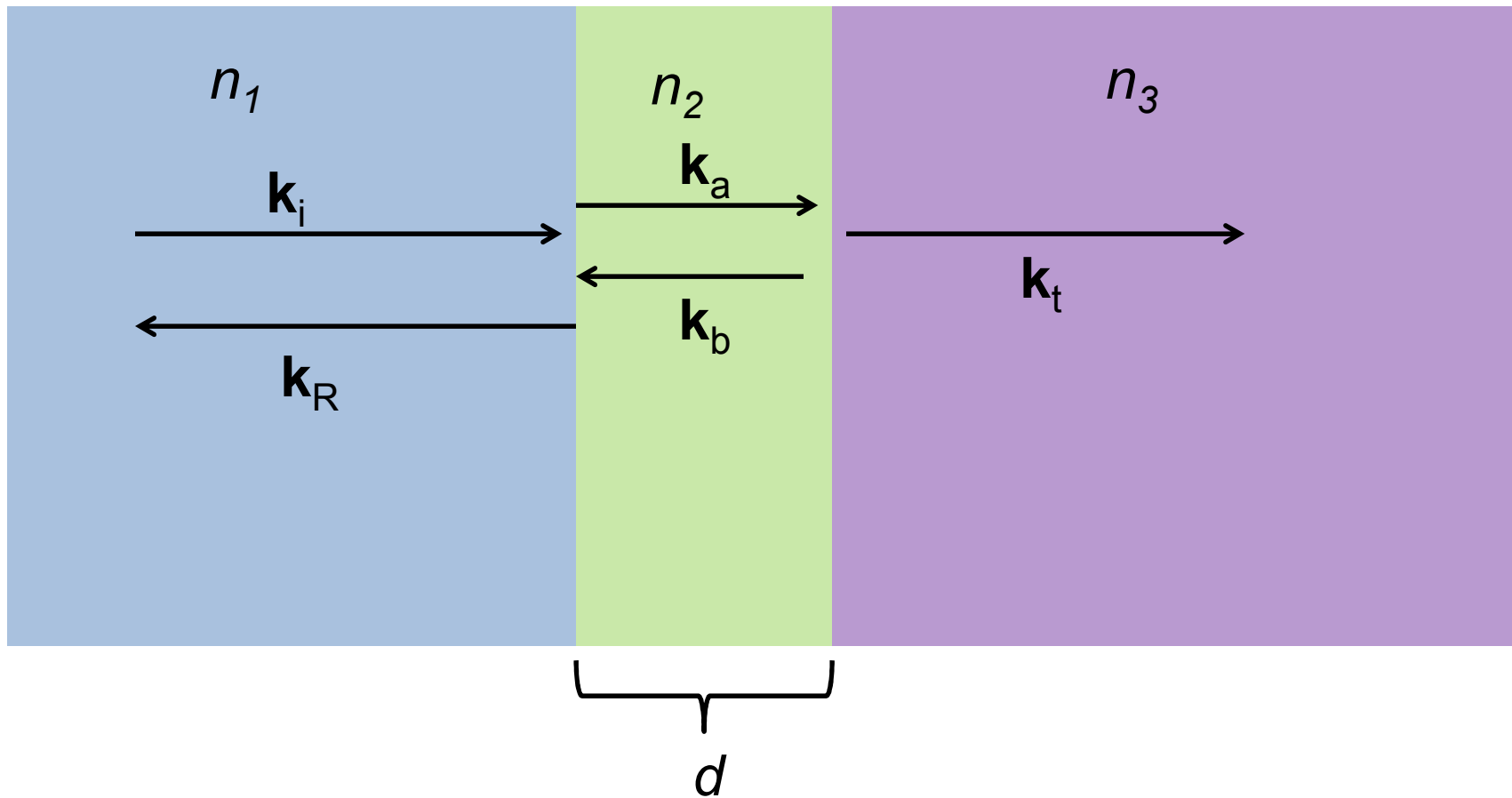
$$\frac{E_{0R}}{E_{0i}} = \frac{\frac{\mu}{\mu'} n' - n}{\frac{\mu}{\mu'} n' + n} \quad \frac{E'_0}{E_{0i}} = \frac{2n}{\frac{\mu}{\mu'} n' + n}$$

Reflectance, transmittance:

$$R = \left| \frac{E_{0R}}{E_{0i}} \right|^2 = \left| \frac{\frac{\mu}{\mu'} n' - n}{\frac{\mu}{\mu'} n' + n} \right|^2$$

$$T = \left| \frac{E'_0}{E_{0i}} \right|^2 \frac{n' \mu}{n \mu'} = \left| \frac{2n}{\frac{\mu}{\mu'} n' + n} \right|^2 \frac{n' \mu}{n \mu'}$$

Multilayer dielectrics (Problem #7.2)



Extension of analysis to anisotropic media --

Extension of analysis to complex dielectric functions

For simplicity assume that $\mu = \mu_0$

Suppose the dielectric function is complex :

$$\epsilon = \epsilon_R + i\epsilon_I \quad \frac{\epsilon}{\epsilon_0} = (n_R + in_I)^2 \equiv \alpha + i\beta$$

$$n_R = \left(\frac{\sqrt{\alpha^2 + \beta^2} + \alpha}{2} \right)^{1/2} \quad n_I = \left(\frac{\sqrt{\alpha^2 + \beta^2} - \alpha}{2} \right)^{1/2}$$

$$\mathbf{E}(\mathbf{r}, t) = \Re\left(\mathbf{E}_0 e^{i\frac{\omega}{c}(n\hat{\mathbf{k}}\cdot\mathbf{r} - ct)}\right) = \Re\left(\mathbf{E}_0 e^{i\frac{\omega}{c}(n_R\hat{\mathbf{k}}\cdot\mathbf{r} - ct)}\right) e^{-\frac{\omega}{c}n_I\hat{\mathbf{k}}\cdot\mathbf{r}}$$

Paul Karl Ludwig Drude 1863-1906



http://photos.aip.org/history/Thumbnails/drude_paul_a1.jpg

Drude model:

Vibrations of charged particles near equilibrium:

$$m\delta \ddot{\mathbf{r}} = q\mathbf{E}_0 e^{-i\omega t} - m\omega_0^2 \delta \mathbf{r} - m\gamma \delta \dot{\mathbf{r}}$$

$$\text{For } \delta \mathbf{r} \equiv \delta \mathbf{r}_0 e^{-i\omega t}, \quad \delta \mathbf{r}_0 = \frac{q\mathbf{E}_0}{m} \frac{1}{\omega_0^2 - \omega^2 - i\omega\gamma}$$

Induced dipole:

$$\mathbf{p} = q \delta \mathbf{r} = \frac{q^2 \mathbf{E}_0}{m} \frac{1}{\omega_0^2 - \omega^2 - i\omega\gamma} e^{-i\omega t}$$

Displacement field:

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P}$$

$$\mathbf{P} = \sum_i \mathbf{p}_i \delta^3(\mathbf{r} - \mathbf{r}_i)$$

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