

# **PHY 712 Electrodynamics**

**11-11:50 AM MWF Olin 107**

## **Plan for Lecture 16:**

**Continue reading Chapter 6**

**A. Maxwell's equations**

**B. Vector and scalar potential  
representations**

**C. Green's function solution  
Maxwell's equations**

## Course schedule for 2013

(Preliminary schedule -- subject to frequent adjustment.)

	Date	JDJ Reading	Topic	Assign.
1	01-16(Wed)	Chap. 1	Introduction, units and Poisson equation.	<a href="#">#1</a>
2	01-18(Fri)	Chap. 1	Electrostatic energy calculations	<a href="#">#2</a>
	01-21(Mon)	<i>No class</i>	<i>MKL Holiday</i>	
3	01-23(Wed)	Chap. 1	Poisson Equation and Green's Functions	<a href="#">#3</a>
4	01-25(Fri)	Chap. 1 & 2	Green's Theorem and Functions	<a href="#">#4</a>
5	01-28(Mon)	Chap. 1 & 2	Brief introduction to numerical methods	<a href="#">#5</a>
6	01-30(Wed)	Chap. 2	Method of images	<a href="#">#6</a>
7	02-01(Fri)	Chap. 3	Cylindrical and spherical geometries	<a href="#">#7</a>
8	02-04(Mon)	Chap. 4	Multipole moments	<a href="#">#8</a>
9	02-06(Wed)	Chap. 4	Dipoles and dielectrics	<a href="#">#9</a>
10	02-08(Fri)	Chap. 4	Microscopic and macroscopic polarizability	
11	02-11(Mon)	Chap. 5	Magnetostatics	<a href="#">#10</a>
12	02-13(Wed)	Chap. 5	Magnetostatic fields	
13	02-15(Fri)	Chap. 5	Magnetic dipole fields	Exam
14	02-18(Mon)	Chap. 5	Permeable media	Exam
15	02-20(Wed)	Chap. 5	Magnetic susceptibility and permeability	Exam
16	02-22(Fri)	Chap. 6	Maxwell's equations	Exam



# Maxwell's equations

Coulomb's law :  $\nabla \cdot \mathbf{D} = \rho_{free}$

Ampere - Maxwell's law :  $\nabla \times \mathbf{H} - \frac{\partial \mathbf{D}}{\partial t} = \mathbf{J}_{free}$

Faraday's law :  $\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$

No magnetic monopoles :  $\nabla \cdot \mathbf{B} = 0$

# Maxwell's equations

Microscopic or vacuum form ( $\mathbf{P} = 0$ ;  $\mathbf{M} = 0$ ):

Coulomb's law :  $\nabla \cdot \mathbf{E} = \rho / \epsilon_0$

Ampere - Maxwell's law :  $\nabla \times \mathbf{B} - \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{J}$

Faraday's law :  $\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$

No magnetic monopoles :  $\nabla \cdot \mathbf{B} = 0$

$$\Rightarrow c^2 = \frac{1}{\epsilon_0 \mu_0}$$

# Formulation of Maxwell's equations in terms of vector and scalar potentials

$$\nabla \cdot \mathbf{B} = 0 \quad \Rightarrow \quad \mathbf{B} = \nabla \times \mathbf{A}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0 \quad \Rightarrow \quad \nabla \times \left( \mathbf{E} + \frac{\partial \mathbf{A}}{\partial t} \right) = 0$$

$$\mathbf{E} + \frac{\partial \mathbf{A}}{\partial t} = -\nabla \Phi$$

or  $\mathbf{E} = -\nabla \Phi - \frac{\partial \mathbf{A}}{\partial t}$

## Formulation of Maxwell's equations in terms of vector and scalar potentials -- continued

Analysis of the scalar and vector potential equations :

$$-\nabla^2\Phi - \frac{\partial(\nabla \cdot \mathbf{A})}{\partial t} = \rho / \epsilon_0$$

$$\nabla \times (\nabla \times \mathbf{A}) + \frac{1}{c^2} \left( \frac{\partial(\nabla\Phi)}{\partial t} + \frac{\partial^2 \mathbf{A}}{\partial t^2} \right) = \mu_0 \mathbf{J}$$

Lorentz gauge form -- require  $\nabla \cdot \mathbf{A}_L + \frac{1}{c^2} \frac{\partial\Phi_L}{\partial t} = 0$

$$-\nabla^2\Phi_L + \frac{1}{c^2} \frac{\partial^2\Phi_L}{\partial t^2} = \rho / \epsilon_0$$

$$-\nabla^2\mathbf{A}_L + \frac{1}{c^2} \frac{\partial^2\mathbf{A}_L}{\partial t^2} = \mu_0 \mathbf{J}$$

## Formulation of Maxwell's equations in terms of vector and scalar potentials -- continued

Lorentz gauge form -- require  $\nabla \cdot \mathbf{A}_L + \frac{1}{c^2} \frac{\partial \Phi_L}{\partial t} = 0$

$$-\nabla^2 \Phi_L + \frac{1}{c^2} \frac{\partial^2 \Phi_L}{\partial t^2} = \rho / \epsilon_0$$

$$-\nabla^2 \mathbf{A}_L + \frac{1}{c^2} \frac{\partial^2 \mathbf{A}_L}{\partial t^2} = \mu_0 \mathbf{J}$$

Alternate potentials:  $\mathbf{A}'_L = \mathbf{A}_L + \nabla \Lambda$  and  $\Phi'_L = \Phi_L - \frac{\partial \Lambda}{\partial t}$

Yields same physics provided that:  $\nabla^2 \Lambda - \frac{1}{c^2} \frac{\partial^2 \Lambda}{\partial t^2} = 0$

## Solution of Maxwell's equations in the Lorentz gauge

$$\nabla^2 \Phi_L - \frac{1}{c^2} \frac{\partial^2 \Phi_L}{\partial t^2} = -\rho / \epsilon_0$$

$$\nabla^2 \mathbf{A}_L - \frac{1}{c^2} \frac{\partial^2 \mathbf{A}_L}{\partial t^2} = -\mu_0 \mathbf{J}$$

Consider the general form of the 3 - dimensional wave equation :

$$\nabla^2 \Psi - \frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = -4\pi f$$

$\Psi(\mathbf{r}, t) \Rightarrow$  wave field

$f(\mathbf{r}, t) \Rightarrow$  source



## Solution of Maxwell's equations in the Lorentz gauge -- continued

$$\nabla^2 \Psi(\mathbf{r}, t) - \frac{1}{c^2} \frac{\partial^2 \Psi(\mathbf{r}, t)}{\partial t^2} = -4\pi f(\mathbf{r}, t)$$

Green's function :

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) G(\mathbf{r}, t; \mathbf{r}', t') = -4\pi \delta^3(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

Formal solution for field  $\Psi(\mathbf{r}, t)$ :

$$\Psi(\mathbf{r}, t) = \Psi_{f=0}(\mathbf{r}, t) + \int d^3 r' \int dt' G(\mathbf{r}, t; \mathbf{r}', t') f(\mathbf{r}', t')$$

## Solution of Maxwell's equations in the Lorentz gauge -- continued

Determination of the form for the Green's function :

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) G(\mathbf{r}, t; \mathbf{r}', t') = -4\pi \delta^3(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

For the case of isotropic boundary values at infinity :

$$G(\mathbf{r}, t; \mathbf{r}', t') = \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta\left( t' - \left( t - \frac{1}{c} |\mathbf{r} - \mathbf{r}'| \right) \right)$$

Formal solution for field  $\Psi(\mathbf{r}, t)$ :

$$\Psi(\mathbf{r}, t) = \Psi_{f=0}(\mathbf{r}, t) + \int d^3 r' \int dt' \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta\left( t' - \left( t - \frac{1}{c} |\mathbf{r} - \mathbf{r}'| \right) \right) f(\mathbf{r}', t')$$

Solution of Maxwell's equations in the Lorentz gauge -- continued  
Analysis of the Green's function :

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) G(\mathbf{r}, t; \mathbf{r}', t') = -4\pi \delta^3(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

Fourier analysis in the time domain -- note that

$$\delta(t - t') = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')}$$

Define :

$$G(\mathbf{r}, t; \mathbf{r}', t') = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')} \tilde{G}(\mathbf{r}, \mathbf{r}', \omega)$$

$$\Rightarrow \left( \nabla^2 + \frac{\omega^2}{c^2} \right) \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) = -4\pi \delta^3(\mathbf{r} - \mathbf{r}')$$

## Solution of Maxwell's equations in the Lorentz gauge -- continued

Analysis of the Green's function (continued) :

$$\left( \nabla^2 + \frac{\omega^2}{c^2} \right) \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) = -4\pi\delta^3(\mathbf{r} - \mathbf{r}')$$

For the case of isotropic boundary values at infinity :

$$\tilde{G}(\mathbf{r}, \mathbf{r}', \omega) = \tilde{G}(\mathbf{r} - \mathbf{r}', \omega)$$

Further assuming that  $\tilde{G}(\mathbf{r} - \mathbf{r}', \omega)$  is isotropic in  $|\mathbf{r} - \mathbf{r}'| \equiv R$  :

$$\left( \frac{1}{R} \frac{d^2}{dR^2} R + \frac{\omega^2}{c^2} \right) \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) = -4\pi\delta^3(\mathbf{r} - \mathbf{r}')$$

$$\text{Solution : } \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) = \frac{1}{R} e^{\pm i\omega R/c}$$

Solution of Maxwell's equations in the Lorentz gauge -- continued  
 Analysis of the Green's function (continued):

$$\begin{aligned}
 \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) &= \frac{1}{|\mathbf{r} - \mathbf{r}'|} e^{\pm i\omega|\mathbf{r} - \mathbf{r}'|/c} \\
 G(\mathbf{r}, t; \mathbf{r}', t') &= \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')} \tilde{G}(\mathbf{r}, \mathbf{r}', \omega) \\
 &= \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')} \frac{1}{|\mathbf{r} - \mathbf{r}'|} e^{\pm i\omega|\mathbf{r} - \mathbf{r}'|/c} \\
 &= \frac{1}{|\mathbf{r} - \mathbf{r}'|} \left( \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t-t' \pm |\mathbf{r} - \mathbf{r}'|/c)} \right) \\
 &= \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta(t - t' \pm |\mathbf{r} - \mathbf{r}'|/c) = \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta(t' - t \mp |\mathbf{r} - \mathbf{r}'|/c)
 \end{aligned}$$

## Solution of Maxwell's equations in the Lorentz gauge -- continued

$$G(\mathbf{r}, t; \mathbf{r}', t') = \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta\left(t' - \left(t \pm |\mathbf{r} - \mathbf{r}'| / c\right)\right)$$

Solution for field  $\Psi(\mathbf{r}, t)$ :

$$\Psi(\mathbf{r}, t) = \Psi_{f=0}(\mathbf{r}, t) +$$

$$\int d^3 r' \int dt' \frac{1}{|\mathbf{r} - \mathbf{r}'|} \delta\left(t' - \left(t - \frac{1}{c} |\mathbf{r} - \mathbf{r}'|\right)\right) f(\mathbf{r}', t')$$