

The ATOMPAW generator

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- Brief history
- Validation
- Extensions

^aCollaborators include Alan Tackett (Vanderbilt U.), Marc Torrent and François Jollet (CEA, France), and Xiao Xu (WFU). Yann Pouillon (UPV/EHU) implemented the build scripts. Work Supported by NSF grants DMR-0405456, 0427055, and 0705239; computations were performed on the Wake Forest University DEAC cluster.

Brief History Leading to PAW

- Kleinman and Bylander, PRL **48** 1425 (1982) “*Efficacious Form for Model Pseudopotentials.*”

Separable Form for Non-Local Pseudopotentials $\hat{V}_{NL} = \sum_{LM} |\xi_{LM}\rangle \delta V_L \langle \xi_{LM}|$

- Gonze, Stumpf, and Scheffler, PRB **44** 8503 (1991) “*Analysis of Separable Potentials.*”

Recipes for taming lurking “ghosts”

- Vanderbilt, PRB **41** 7892 (1990) “*Soft Self-Consistent Pseudopotentials in a Generalized Eigenvalue Formalism.*”

Ultra-Soft Pseudopotentials (USPP)

- Blöchl, PRB **50** 17953 (1994) “*Projector augmented-wave method.*”

PAW

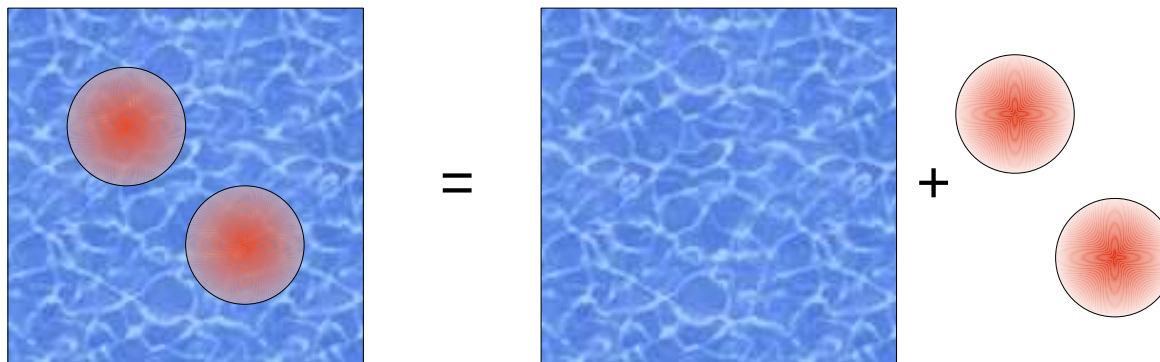
The PAW Formalism

PAW transformation for a valence wavefunction:

$$\Psi_v(\mathbf{r}) = \tilde{\Psi}_v(\mathbf{r}) + \sum_{ai} \left(\phi_i^a(\mathbf{r} - \mathbf{R}^a) - \tilde{\phi}_i^a(\mathbf{r} - \mathbf{R}^a) \right) \langle p_i^a | \tilde{\psi}_v \rangle. \quad (1)$$

Partitioning of evaluation space. For example, the valence electron energy:

$$E_{tot}^{v\text{ale}} = \underbrace{\tilde{E}_{tot}}_{\text{pseudo energy}} + \sum_a \underbrace{\left(E_{tot}^a - \tilde{E}_{tot}^a \right)}_{\text{atom-centered corrections}}. \quad (2)$$



Partitioning of the Coulomb Interactions

Both Hartree and Fock interactions can be expressed in terms of pair densities derived from products of two wavefunctions $\rho_{vw}(\mathbf{r}) \equiv \Psi_v^*(\mathbf{r})\Psi_w(\mathbf{r})$. The PAW transformation the these products take the form:

$$\rho_{vw}(\mathbf{r}) = \tilde{\rho}_{vw}(\mathbf{r}) + \sum_a \{ \rho_{vw}^a(\mathbf{r} - \mathbf{R}^a) - \tilde{\rho}_{vw}^a(\mathbf{r} - \mathbf{R}^a) \}, \quad (3)$$

where

$$\tilde{\rho}_{vw}(\mathbf{r}) \equiv \tilde{\Psi}_v^*(\mathbf{r})\tilde{\Psi}_w(\mathbf{r}), \quad (4)$$

and

$$\rho_{vw}^a(\mathbf{r}) - \tilde{\rho}_{vw}^a(\mathbf{r}) \equiv \sum_{ij} \langle \tilde{\Psi}_v | p_i^a \rangle \langle p_j^a | \tilde{\Psi}_w \rangle \left\{ \phi_i^{a*}(\mathbf{r})\phi_j^a(\mathbf{r}) - \tilde{\phi}_i^{a*}(\mathbf{r})\tilde{\phi}_j^a(\mathbf{r}) \right\}. \quad (5)$$

Equivalently, we can write:

$$\rho_{vw}(\mathbf{r}) = \tilde{\rho}_{vw}(\mathbf{r}) + \hat{\rho}_{vw}(\mathbf{r}) + \sum_a (\rho_{vw}^a(\mathbf{r} - \mathbf{R}^a) - \tilde{\rho}_{vw}^a(\mathbf{r} - \mathbf{R}^a) - \hat{\rho}_{vw}^a(\mathbf{r} - \mathbf{R}^a)), \quad (6)$$

where the “compensation charge”

$$\hat{\rho}_{vw}(\mathbf{r}) \equiv \sum_a \hat{\rho}_{vw}^a(\mathbf{r} - \mathbf{R}^a) \quad (7)$$

is a smooth function localized within the atomic augmentation spheres such that:

$$\int d^3r' \frac{\rho_{vw}^a(\mathbf{r}') - \tilde{\rho}_{vw}^a(\mathbf{r}') - \hat{\rho}_{vw}^a(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} = \begin{cases} V_{vw}^a(\mathbf{r}) & \text{for } |\mathbf{r} - \mathbf{R}^a| \leq r_c^a \\ 0 & \text{otherwise} \end{cases}. \quad (8)$$

ATOMPAW is 10 years old!



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A Projector Augmented Wave (PAW) code for electronic structure calculations, Part I: *atompaw* for generating atom-centered functions

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$$\Rightarrow \{\phi_i(r), \tilde{\phi}_i(r), \tilde{p}_i(r)\} \text{ and } V_{loc}(r)$$

Generate input to *pwaw*, *abinit*, *socorro*, *quantum-espresso*

Database of PAW Atomicdata Files

<http://pwpaw.wfu.edu>

Periodic Table of the Elements for PAW Functions

1 H																2 He	
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba		72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra		104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Uun	111 Uuu	112 Uub		114 Uuq		116 Uuh		

		57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu	
		89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr	

Key

Modified Elements (using \geq atompaw_2.1) Unmodified Elements (2005)



Comments

- The contents of this periodic table were created by John Tumbleston (currently a graduate student at UNC) and modified by Nick Dellaripa, David Harris, and Xiao Xu under the guidance of [N.A.W. Holzwarth](#)

ATOMPAW User's Manual by Marc Torrent

A user's guide for *atompaw* code

Marc Torrent

*Commissariat à l'Energie Atomique et aux Energies Alternatives
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Source code URL : <http://pwpaw.wfu.edu>

Revised November 4th, 2010

*Compatible with **atompaw v3.0** and later*

Consistency of *abinit* and *pwpaw*

Computer Physics Communications 181 (2010) 1862–1867

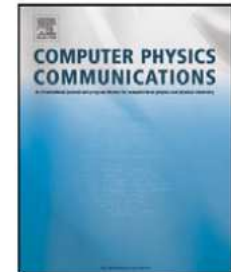


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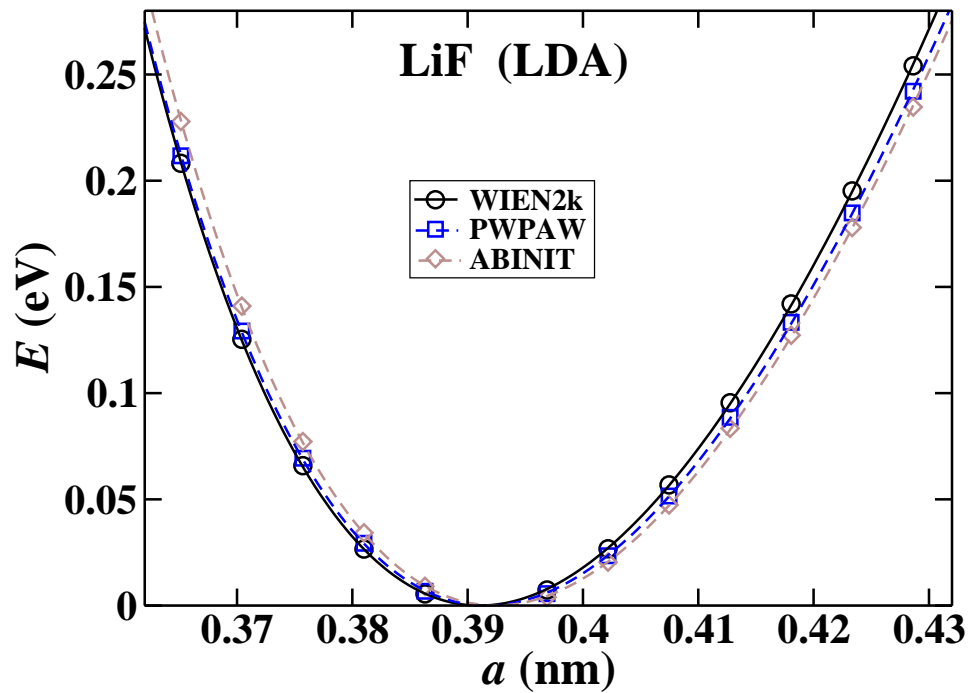
Electronic structure packages: Two implementations of the projector augmented wave (PAW) formalism

Marc Torrent^a, N.A.W. Holzwarth^{b,*}, François Jollet^a, David Harris^b, Nicholas Lepley^b, Xiao Xu^b

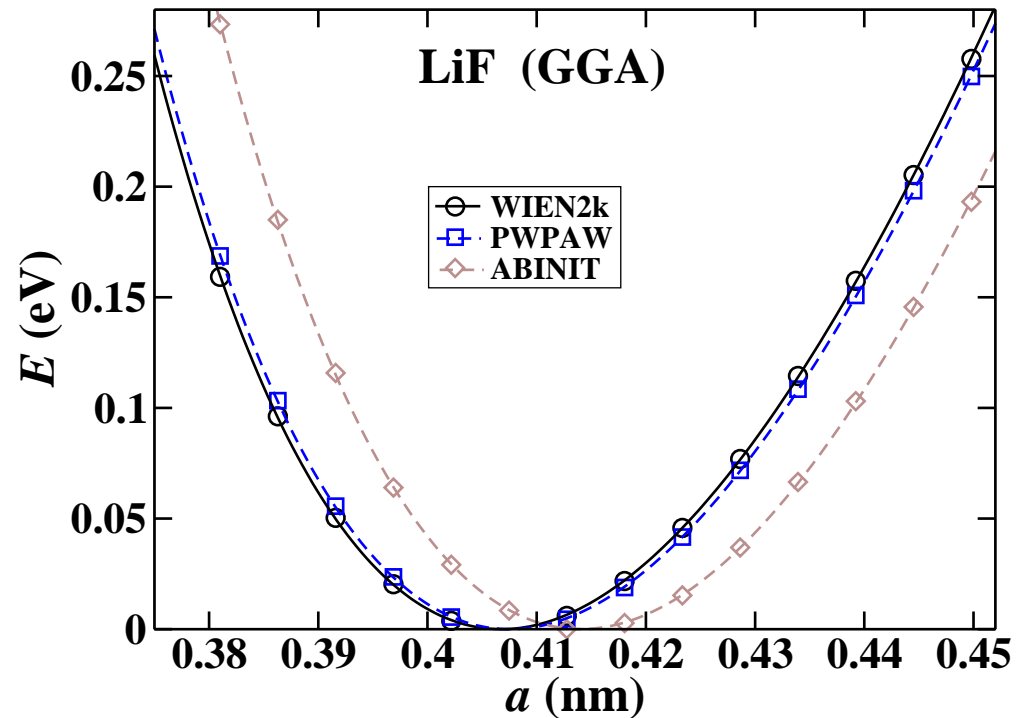
^a CEA, DAM, DIF, F-91297 Arpajon, France

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Mystery < 2010



LDA binding energy curve for LiF



GGA binding energy curve for LiF

Mystery Traced to Exchange-Correlation Terms

Blöchl's formulation

(PRB **50**, 17953 (1994))

$$E_{xc}^B = E_{xc}[\tilde{n} + \tilde{n}_c] + \sum_a (E_{xc}^a[n^a + n_c^a] - E_{xc}^a[\tilde{n}^a + \tilde{n}_c^a]).$$

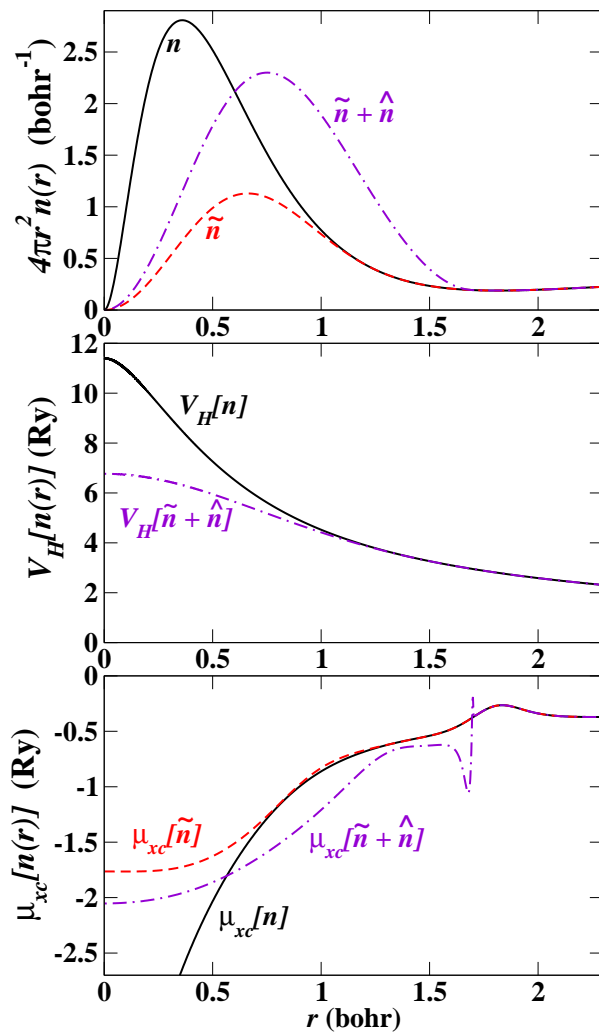
Kresse's formulation

(PRB **59**, 1758 (1999))

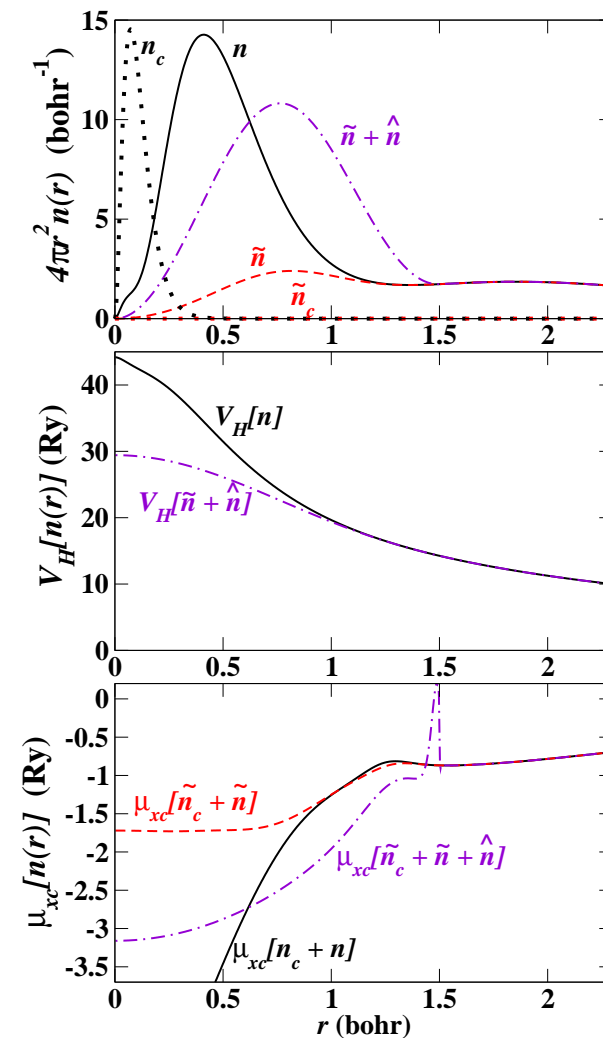
$$E_{xc}^K = E_{xc}[\tilde{n} + \tilde{n}_c + \hat{n}] + \sum_a (E_{xc}^a[n^a + n_c^a] - E_{xc}^a[\tilde{n}^a + \tilde{n}_c^a + \hat{n}^a]).$$

Example Exchange-Correlation Potentials

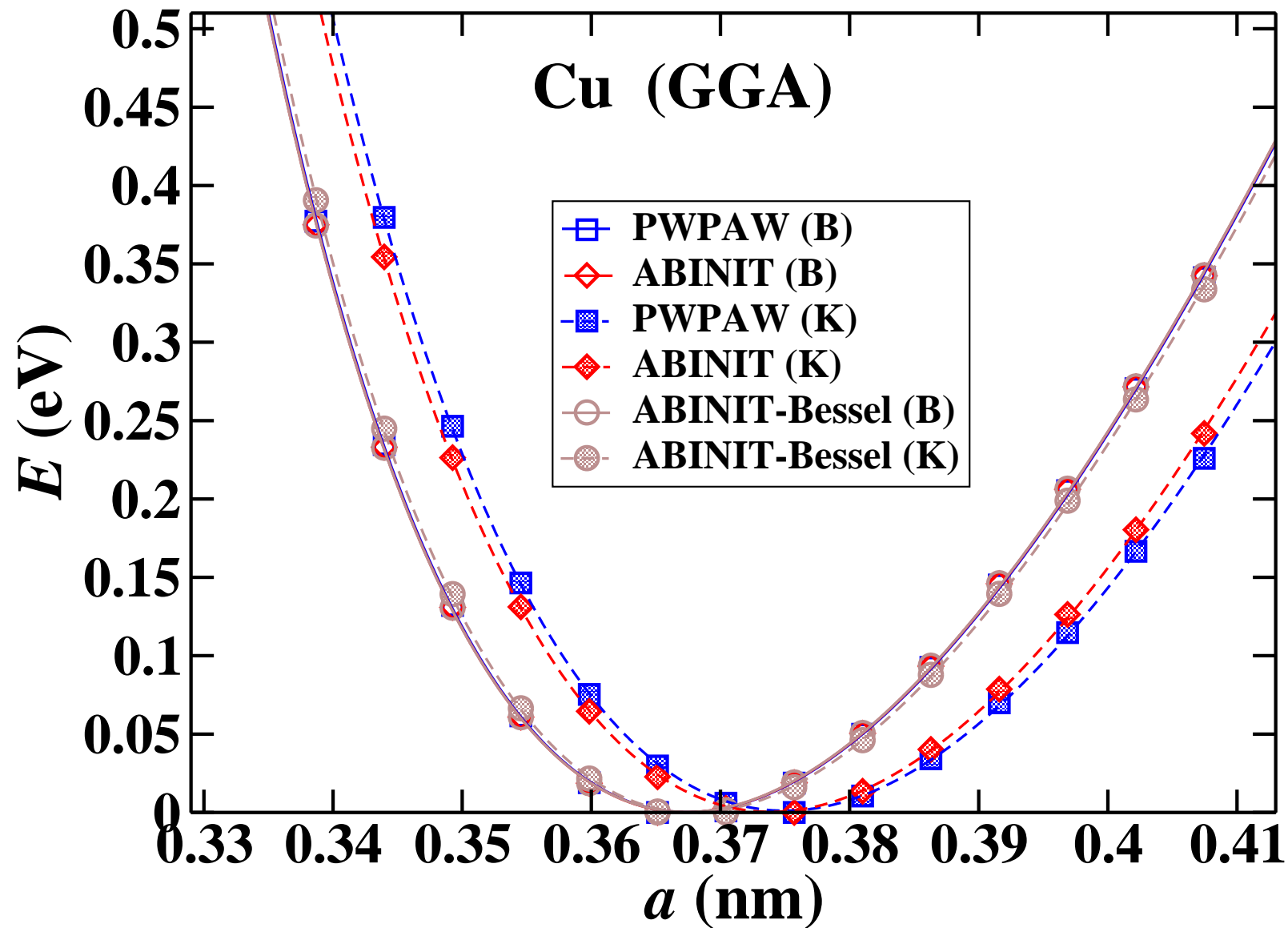
Li GGA ($1s^2 2s^1$)



Si GGA ($2s^2 3s^2 2p^6 3p^2$)

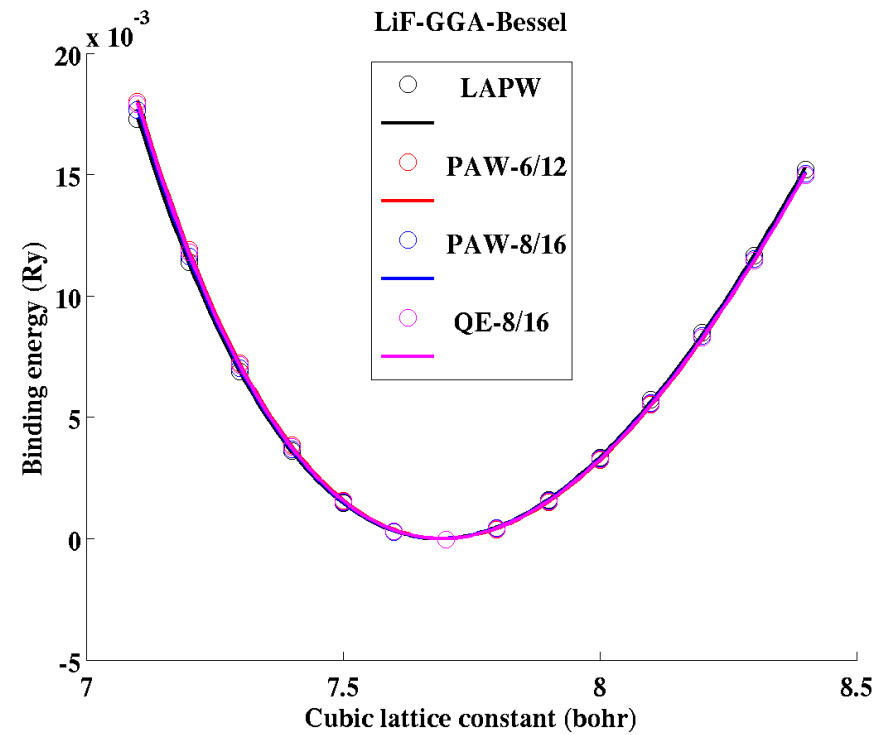
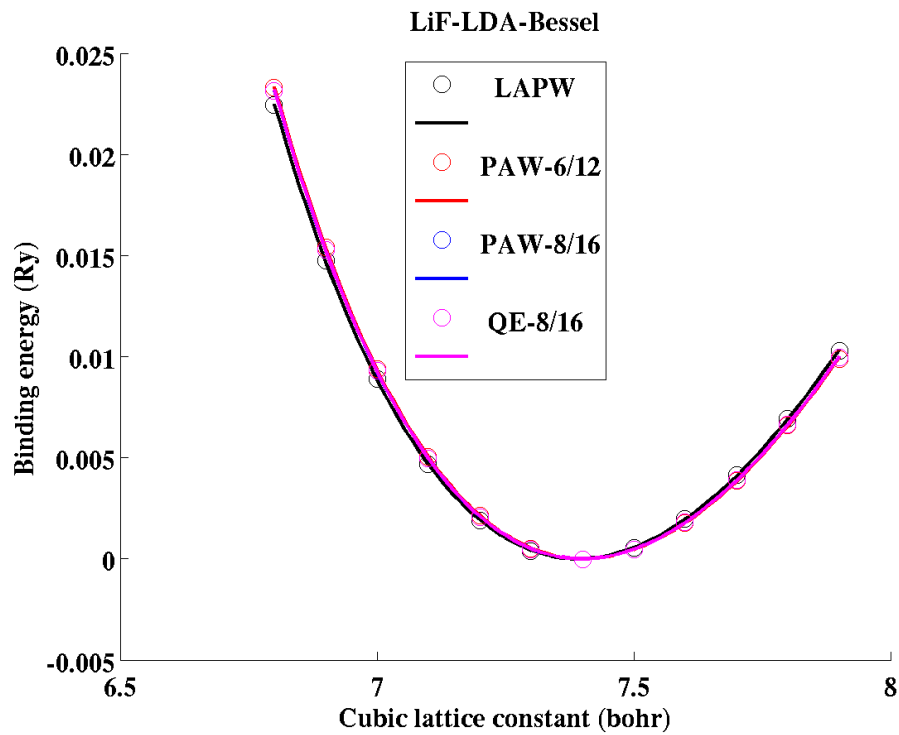


Mystery Solved – *abinit* and *pwpaw* now agree



Bind energy curves for fcc Cu using GGA, comparing Blöchl and Kresse formulations.

Validation: Consistency between various codes



Heats of formation ΔH_{cal} (eV per formula unit) calculated for the lithium (thio)phosphate and related materials, comparing USPP, PAW and experiment (CRC or NIST).

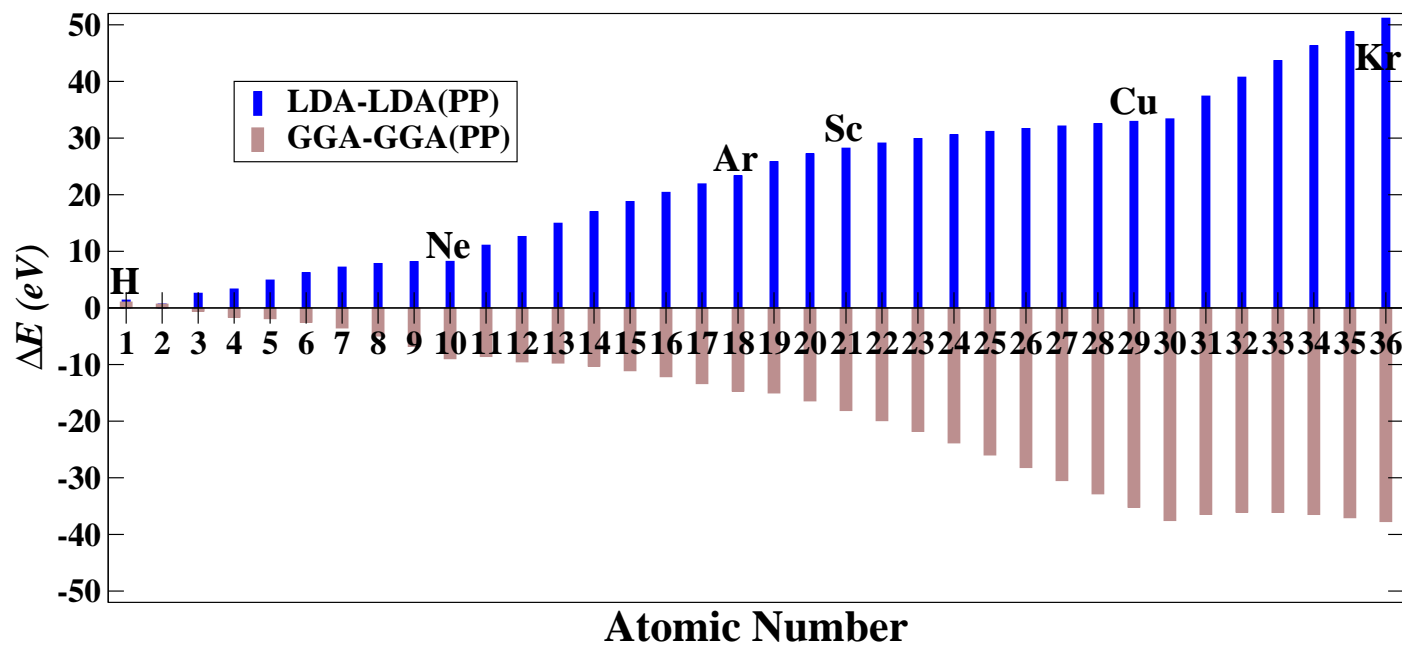
Material	(USPP)	(PAW)	exp
Li ₂ O	-6.18	-6.19	-6.20
Li ₂ O ₂	-6.53	-6.52	-6.57
β -Li ₃ PO ₄	-21.41	-21.39	
γ -Li ₃ PO ₄	-21.38	-21.36	-21.72
Li ₄ P ₂ O ₆	-30.02	-29.93	
Li ₄ P ₂ O ₇	-34.25	-34.21	
Li ₇ P ₃ O ₁₁	-55.26	-55.26	
Li ₂ S	-4.30	-4.30	-4.57
Li ₂ S ₂	-4.10	-4.10	
β -Li ₃ PS ₄	-8.39	-8.35	
γ -Li ₃ PS ₄	-8.19	-8.16	
Li ₄ P ₂ S ₆	-12.45	-12.38	
Li ₄ P ₂ S ₇	-11.62	-11.54	
Li ₇ P ₃ S ₁₁	-20.06	-19.94	
SO ₃	-4.83	-4.86	-4.71
Li ₂ SO ₄	-14.74	-14.76	-14.89

Ongoing work – extension of PAW formalism to orbital dependent functionals

- Motivation
- Treatment of core electrons
- Hartree-Fock equations
- Approximate OEP equations (KLI)

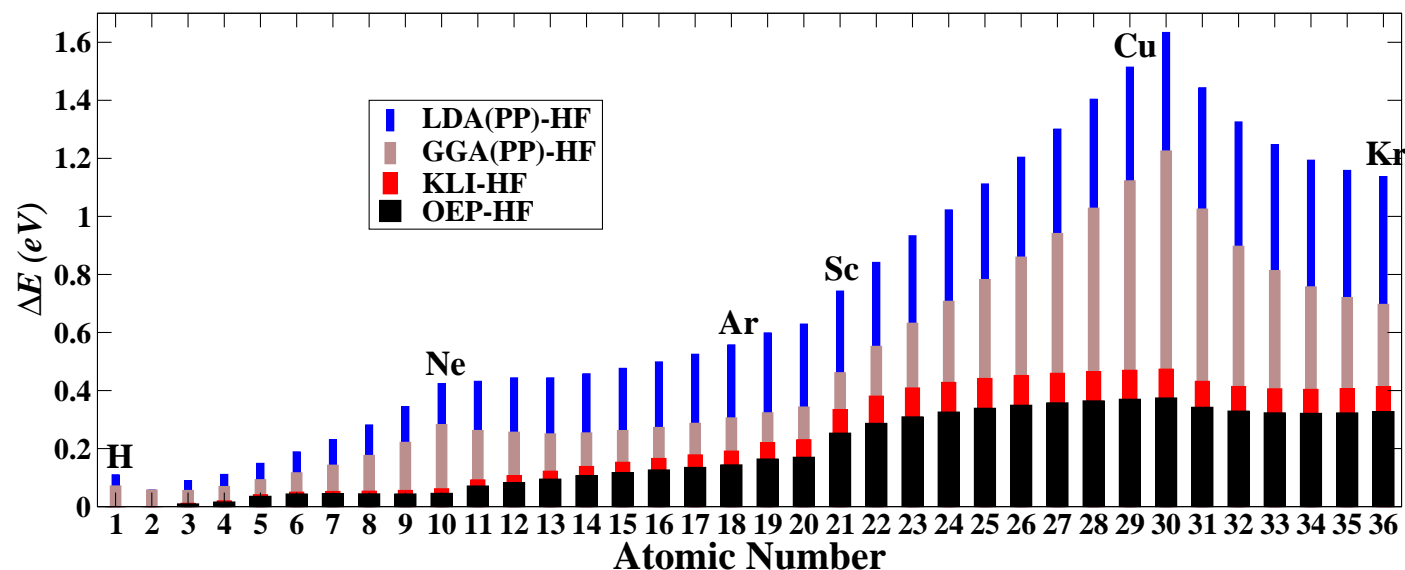
“Post-processing” estimate of Fock exchange relative to self-consistent LDA or GGA Results

$$E_{xc}^{\text{GGA/LDA}} - E_x^{\text{Fock}}$$



Comparison of Self-Consistent and “post-processing” Treatments of Fock exchange

$$E_{tot}^{OEP/KLI/PP} - E_{tot}^{HF}$$



“KLI” reference: Krieger, Li, and Iafrate, PRA **45**, 101 (1992)

More recent perspective on OEP and KLI: Bulat and Levy, PRB **80**, 052510 (2009)

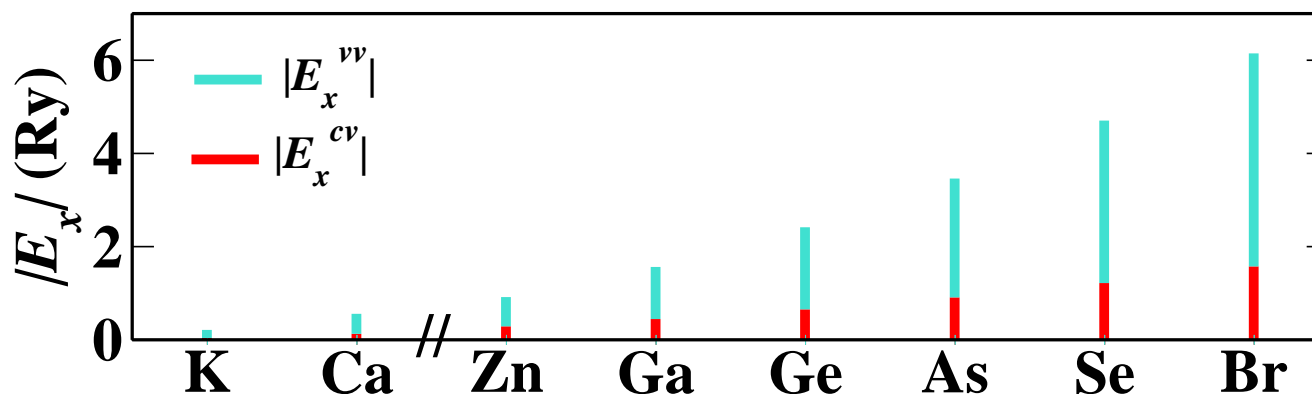
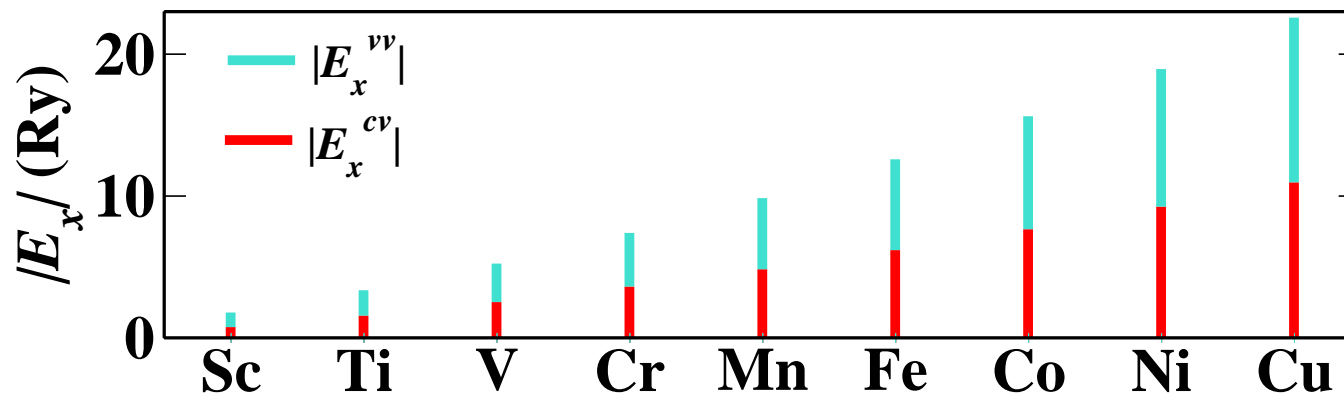
Orbital Dependent Functionals

Fock Exchange Energy

$$E_x \equiv -e^2 \sum_{pq} ' \int \int d^3r d^3r' \frac{\Psi_p^*(\mathbf{r}) \Psi_q(\mathbf{r}) \Psi_q^*(\mathbf{r}') \Psi_p(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Core and Valence Electron Contributions to Fock Exchange Energy

$$E_x \equiv E_x^{vv} + E_x^{cv} + E_x^{cc}$$



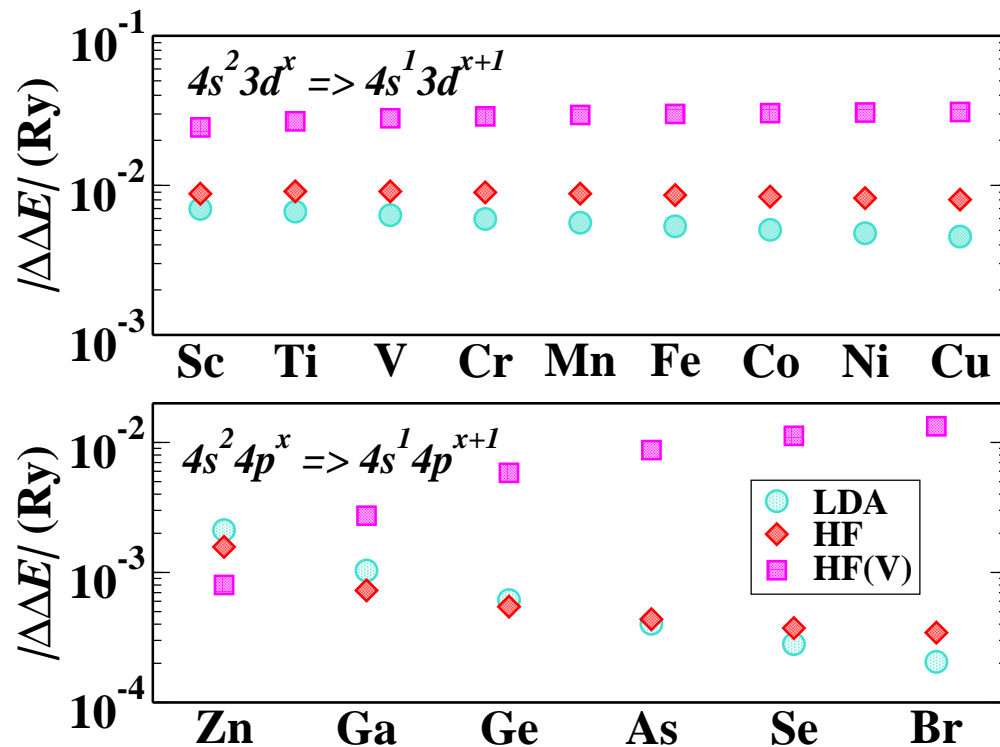
Core-Valence contributions in Fock Exchange

✓ Frozen Core *Orbital* Approximation to Fock Exchange Energy

$$E_x^{cv} \approx -e^2 \sum_{vc} \int \int d^3r d^3r' \frac{\Psi_v^*(\mathbf{r}) \bar{\Psi}_c(\mathbf{r}) \bar{\Psi}_c^*(\mathbf{r}') \Psi_v(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Frozen Core *Potential* Approximation to Fock Exchange Energy

$$E_x^{cv} \approx \int d^3r V_x^{cv}(\mathbf{r}) \rho_v(\mathbf{r})$$



Frozen core *orbital* and frozen core *potential* errors in various excitation energies.

Hartree Fock Equations for Atoms

Xu and Holzwarth, PRB **81**, 245105 (2010)

$$H^{\text{HF}}(r)\psi_p^{\text{HF}}(r) + X_p(r) - \sum_{q; N_q > 0} \lambda_{qp} \psi_q^{\text{HF}}(r) = 0.$$

Here,

$$H^{\text{HF}}(r) \equiv \mathcal{K} + V_N(r) + V_H(r). \quad (9)$$

The exchange kernel function is given by

$$X_p(r) \equiv \frac{1}{N_p} \frac{\delta E_x}{\delta \psi_p^*} = - \sum_q \sum_{L=|l_p-l_q|}^{l_p+l_q} \frac{1}{N_p} \Theta_{pq}^L W_{qp}^L(r) \psi_q^{\text{HF}}(r),$$

where

$$W_{qp}^L(r) \equiv e^2 \int dr' \frac{r_{<}^L}{r_{>}^{L+1}} \psi_q^{\text{HF}*}(r') \psi_p^{\text{HF}}(r').$$

Hartree Fock PAW Equations for Atoms

$$\mathcal{H}_{\text{HF}}^{\text{PAW}}(\mathbf{r})\tilde{\Psi}_v^{\text{HF}}(r) + X_v^{\text{PAW}}(\mathbf{r}) - \sum_q \lambda_{qv} \mathcal{O}_{\text{HF}}^{\text{PAW}} \tilde{\Psi}_q^{\text{HF}}(\mathbf{r}) = 0.$$

The single particle term takes the form

$$\mathcal{H}_{\text{HF}}^{\text{PAW}}(\mathbf{r}) = \tilde{\mathcal{H}}^{\text{HF}} + \sum_{aij} |\tilde{P}_i^a\rangle D_{ij}^{a\text{HF}} \langle \tilde{P}_j^a|,$$

where the pseudo Hamiltonian-like terms depend on the pseudopotentials due to the nuclear and Hartree interactions:

$$\tilde{\mathcal{H}}^{\text{HF}}(\mathbf{r}) \equiv \mathcal{K} + \tilde{V}_N(\mathbf{r}) + \tilde{V}_H(\mathbf{r}). \quad (10)$$

The exchange function term takes the form

$$X_v^{\text{PAW}}(\mathbf{r}) = \tilde{X}_v(\mathbf{r}) + \sum_{ai} |\tilde{P}_i^a\rangle X_{iv}^a,$$

where the pseudo-exchange kernel function takes the form:

$$\tilde{X}_v(r) \equiv - \sum_q \sum_{L=|l_v-l_q|}^{l_v+l_q} \frac{1}{N_v} \Theta_{vq}^L \tilde{W}_{qv}^L(r) \tilde{\psi}_q^{\text{HF}}(r).$$

$$\tilde{W}_{qv}^L(r) \equiv e^2 \int dr' \frac{r_{<}^L}{r_{>}^{L+1}} \left(\tilde{\psi}_q^{\text{HF}*}(r') \tilde{\psi}_v^{\text{HF}}(r') + M_{qv}^L(r') \right).$$

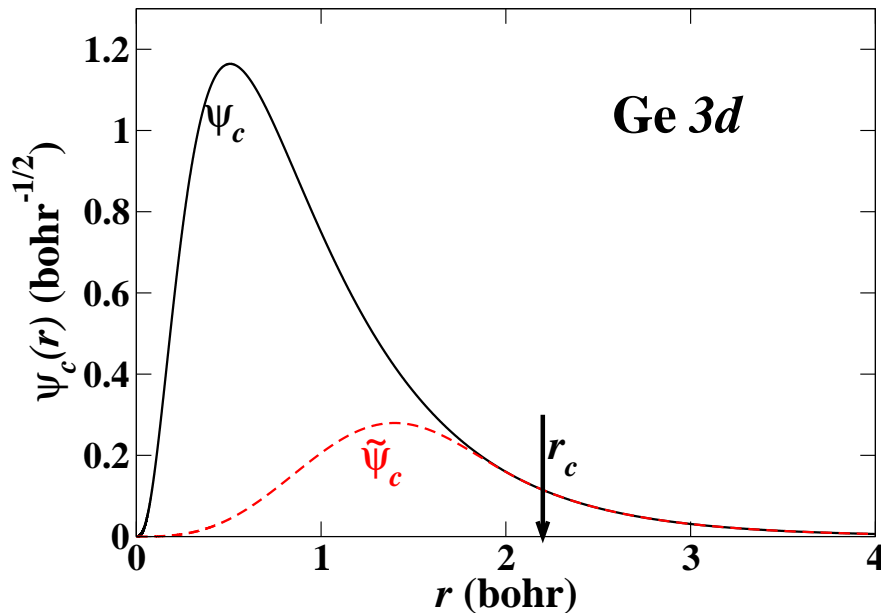
Hartree Fock PAW Equations for Atoms – (continued)

$$\begin{aligned}
 X_{iv}^a = & - \sum_{v'} \sum_{L=|l_v-l_{v'}|}^{l_v+l_{v'}} \frac{1}{N_v} \Theta_{vv'}^L \\
 & \times \left(\sum_{jkl} \langle \tilde{P}_j^a | \tilde{\Psi}_{v'}^{\text{HF}} \rangle \langle \tilde{\Psi}_{v'}^{\text{HF}} | \tilde{P}_k^a \rangle \langle \tilde{P}_l^a | \tilde{\Psi}_v^{\text{HF}} \rangle \left(R_{ij;kl}^{aL} - \tilde{R}_{ij;kl}^{aL} \right) \right. \\
 & \quad \left. + \sum_j \langle \tilde{P}_j^a | \tilde{\Psi}_{v'}^{\text{HF}} \rangle Z_{v'v;ij}^{aL} \right) \\
 & - \sum_c \sum_{L=|l_v-l_c|}^{l_v+l_c} \frac{1}{N_v} \Theta_{vc}^L \left(\sum_j \langle \tilde{P}_j^a | \tilde{\Psi}_v^{\text{HF}} \rangle \left(R_{ic;cj}^{aL} - \tilde{R}_{ic;cj}^{aL} \right) \right. \\
 & \quad \left. + Z_{cv;ic}^{aL} \right),
 \end{aligned}$$

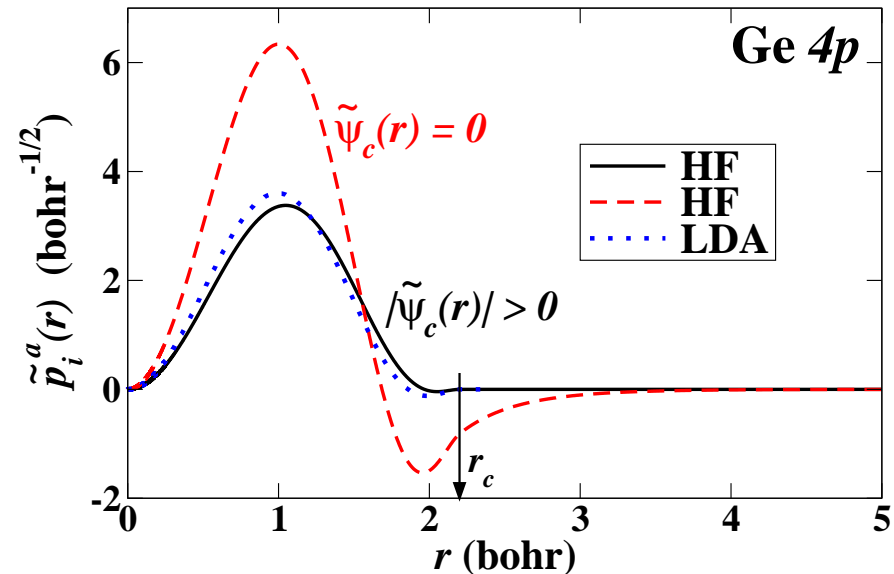
where

$$Z_{qv;ij}^{aL} \equiv \int dr \tilde{W}_{qv}^L(r) \tilde{m}_{ij}^{aL}(r). \quad (11)$$

Some details about “extended” core states



Radial Hartree-Fock wavefunction $\psi_c(r)$ for the $3d$ core state of Ge compared with the constructed pseudo-wavefunction $\tilde{\psi}_c(r)$.



Radial PAW projector functions for $4p$ states of Ge, comparing effects of including and excluding pseudo-core orbital functions $\tilde{\psi}_c(r)$ within the Hartree-Fock (HF) formulation and also comparing the corresponding LDA projector.

Some details about Kohn-Sham vs Hartree-Fock projector functions

Kohn-Sham

$$\left(\tilde{\mathcal{H}}^{\text{KS}}(r) - \varepsilon_i^a\right) \tilde{\Phi}_i^a(\mathbf{r}) = \sum_j \tilde{P}_j^a(\mathbf{r}) \langle \tilde{\Phi}_j^a | \tilde{\mathcal{H}}^{\text{KS}} - \varepsilon_i^a | \tilde{\Phi}_i^a \rangle$$

Hartree-Fock

$$\tilde{H}^{\text{HF}}(\mathbf{r}) \tilde{\Phi}_i^a(\mathbf{r}) + \tilde{X}_i(\mathbf{r}) - \sum_{q; N_q > 0} \lambda_{qi} \tilde{\Psi}_q^{\text{HF}}(\mathbf{r}) =$$
$$\sum_j \tilde{P}_j^a(\mathbf{r}) \left(\langle \tilde{\Phi}_j^a | \tilde{H}^{\text{HF}} | \tilde{\Phi}_i^a \rangle + \langle \tilde{\Phi}_j^a | \tilde{X}_i \rangle - \sum_{q; N_q > 0} \lambda_{qi} \langle \tilde{\Phi}_j^a | \tilde{\Psi}_q^{\text{HF}} \rangle \right),$$

Orbital Dependent Functionals + Kohn-Sham = Optimized Effective Potential (OEP) Formalism

Fock Exchange Energy

$$E_x \equiv -e^2 \sum_{pq} ' \int \int d^3r d^3r' \frac{\Psi_p^*(\mathbf{r}) \Psi_q(\mathbf{r}) \Psi_q^*(\mathbf{r}') \Psi_p(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Kohn-Sham Equations

$$H^{KS}(\mathbf{r}) \Psi_p(\mathbf{r}) = \varepsilon_p \Psi_p(\mathbf{r})$$

$$H^{KS}(\mathbf{r}) \equiv \mathcal{K} + V_N(\mathbf{r}) + V_H(\mathbf{r}) + V_x(\mathbf{r})$$

$$V_x(\mathbf{r}) = \frac{\delta E_x}{\delta \rho(\mathbf{r})}$$

KLI* approximation to Optimized Effective Potential (OEP) Formalism

* Krieger, Li, and Iafrate, PRA **45**, 101 (1992)

$$V_x^{\text{KLI}}(r) n(r) = \sum_p N_p \psi_p(r) X_p^{\text{KLI}}(r) + \sum_p N_p |\psi_p(r)|^2 (\bar{V}_{x p}^{\text{KLI}} - \bar{U}_{x p}^{\text{KLI}})$$

where

$$n(r) \equiv \sum_p N_p |\psi_p(r)|^2$$

$$\sum_{q \neq 0} [\delta_{pq} - \Gamma_{pq} N_q] \bar{V}_{x q}^{\text{KLI}} = \Xi_p - \sum_{q \neq 0} \Gamma_{pq} N_q \bar{U}_{x q}^{\text{KLI}}.$$

Here

$$\Gamma_{pq} \equiv \int dr \frac{|\psi_p(r)|^2 |\psi_q(r)|^2}{n(r)},$$

and

$$\Xi_p \equiv \int dr \frac{|\psi_p(r)|^2 \sum_q N_q \psi_q(r) X_q^{\text{KLI}}(r)}{n(r)}.$$

Relationship between KLI approximation to OEP

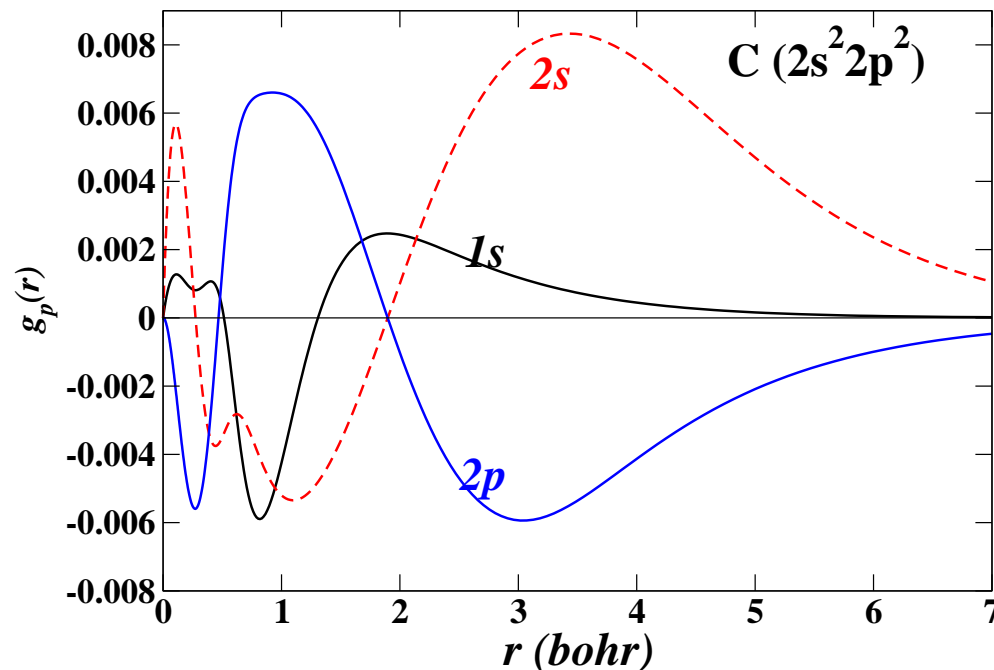
Bulat and Levy, PRB **80**, 052510 (2009)

$$V_x^{\text{OEP}}(r) = \underbrace{V_x^{\text{occ}}(r)}_{\text{slight extension to KLI}} + f(\{\phi_q^{\text{unocc}}(r)\}).$$

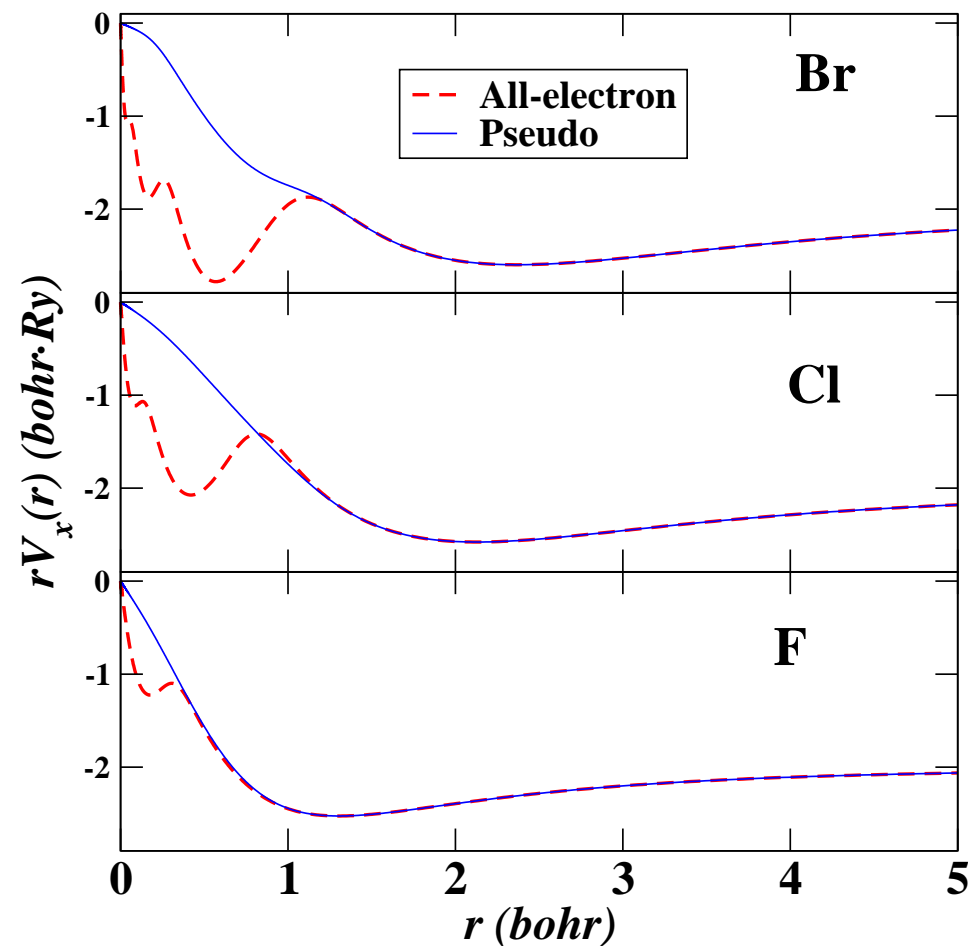
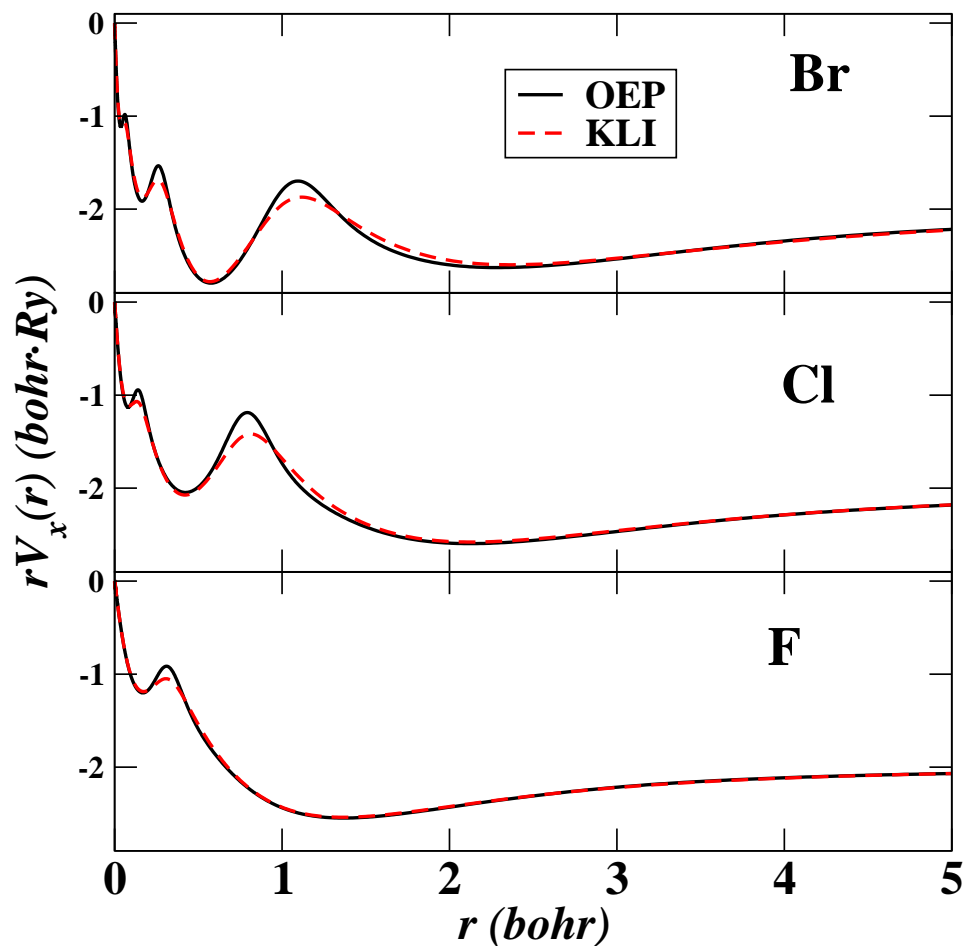
slight extension to KLI

Orbital shift function of full OEP

$$(\mathcal{H} - \epsilon_p) g_p(r) = X_p(r) - V_x(r) \psi_p(r) - (\bar{U}_{x p} - \bar{V}_{x p}) \psi_p(r); \quad \sum_p N_p g_p(r) \psi_p(r) = 0.$$



Comparison of $V_x(r)$ AE-OEP, AE-KLI, and PAW-KLI



PAW-KLI formalism

$$V_x^{\text{KLI}}(r) n(r) \Rightarrow \tilde{V}_x^{\text{KLI}}(r) \tilde{n}(r) + \left(V_x^{a\text{KLI}}(r) n^a(r) - \tilde{V}_x^{a\text{KLI}}(r) \tilde{n}^a(r) \right).$$

\Rightarrow Assume that each term satisfies the KLI equations in their respective spacial and functional domains.

$$\tilde{V}_x^{\text{KLI}}(r) \tilde{n}(r) = \sum_v N_v \tilde{\psi}_v(r) \tilde{X}_v^{\text{KLI}}(r) + \sum_v N_v \left| \tilde{\psi}_v(r) \right|^2 \left(\bar{V}_{x v}^{\text{KLI}} - \bar{U}_{x v}^{\text{KLI}} \right).$$

$$V_x^{a\text{KLI}}(r) n^a(r) = \sum_p N_p \psi_p^a(r) X_p^{a\text{KLI}}(r) + \sum_p N_p \left| \psi_p^a(r) \right|^2 \left(\bar{V}_{x p}^{\text{KLI}} - \bar{U}_{x p}^{\text{KLI}} \right),$$

$$\tilde{V}_x^{a\text{KLI}}(r) \tilde{n}^a(r) = \sum_v N_v \tilde{\psi}_v^a(r) \tilde{X}_v^{a\text{KLI}}(r) + \sum_v N_v \left| \tilde{\psi}_v^a(r) \right|^2 \left(\bar{V}_{x v}^{\text{KLI}} - \bar{U}_{x v}^{\text{KLI}} \right).$$

Summary and Conclusions

- Thanks to the *abinit* community for continuously providing excellent computational tools, productive framework for collaboration, . . .
- *atompaw* now produces atomic datasets for *abinit*, *quantum espresso*, *pwpaw*, *socorro*
- For Fock exchange, we argue that the frozen core *orbital* approach offers better accuracy and is compatible with the PAW formalism.
- We have formulated the PAW equations for Hartree-Fock and KLI and are currently testing them in *pwpaw*.
- Future work –
 - Finish testing Hartree-Fock and KLI in *pwpaw*.
 - Extend approach to hybrid functionals ??
 - Possibly extend Hartree-Fock and KLI in other codes.