

Simulations of Idealized Solid Electrolytes for Solid State Battery Designs*

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*Supported by NSF Grant DMR-1507942, WFU's DEAC cluster, and Center for Energy, Environment, and Sustainability.

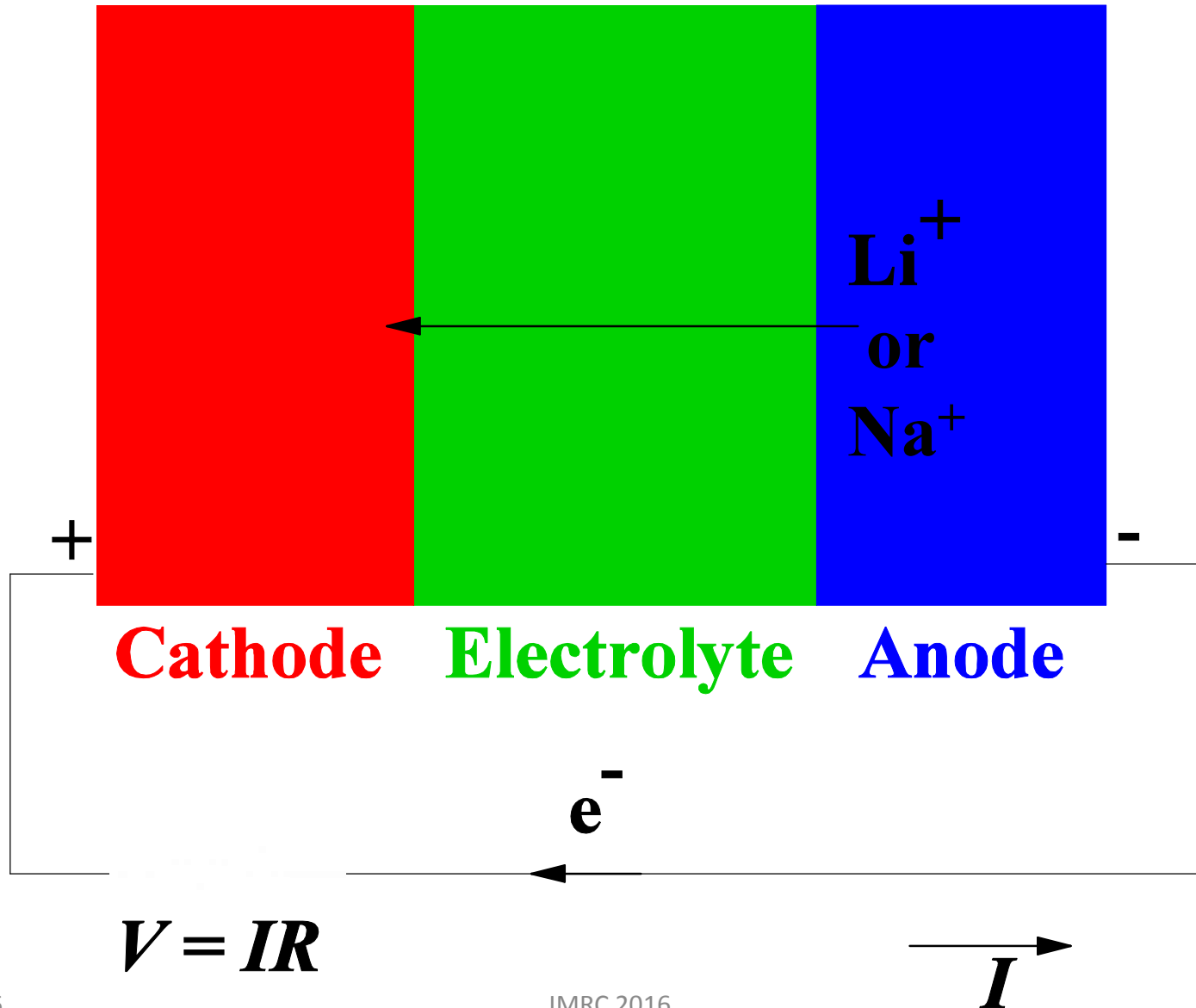
**With help from: Nicholas Lepley (WFU Ph. D., Dec. 2015) , Ahmad Al-Qawasmeh, Jason Howard, and Larry Rush (WFU graduate students), Dr. Yaojun Du (former Postdoc), colleagues from the WFU chemistry department – Dr. Keerthi Senevirathne (currently at Florida A & M U.), Dr. Cynthia Day, Professor Michael Gross, Professor Abdessadek Lachgar, and Zachary Hood (currently at Georgia Tech and ORNL), and Professor Jennifer Aitken from Duquesne U.

Outline

- **Motivation – Why solid electrolytes?**
- **Computational tools & reality checks; “first principles” calculations**
- **Some examples based on crystalline materials**
 - **Li phosphorus oxynitrides (first developed at Oak Ridge National Laboratory)**
 - **Li thiophosphates**
 - **Other examples**
- **Summary and remaining challenges**

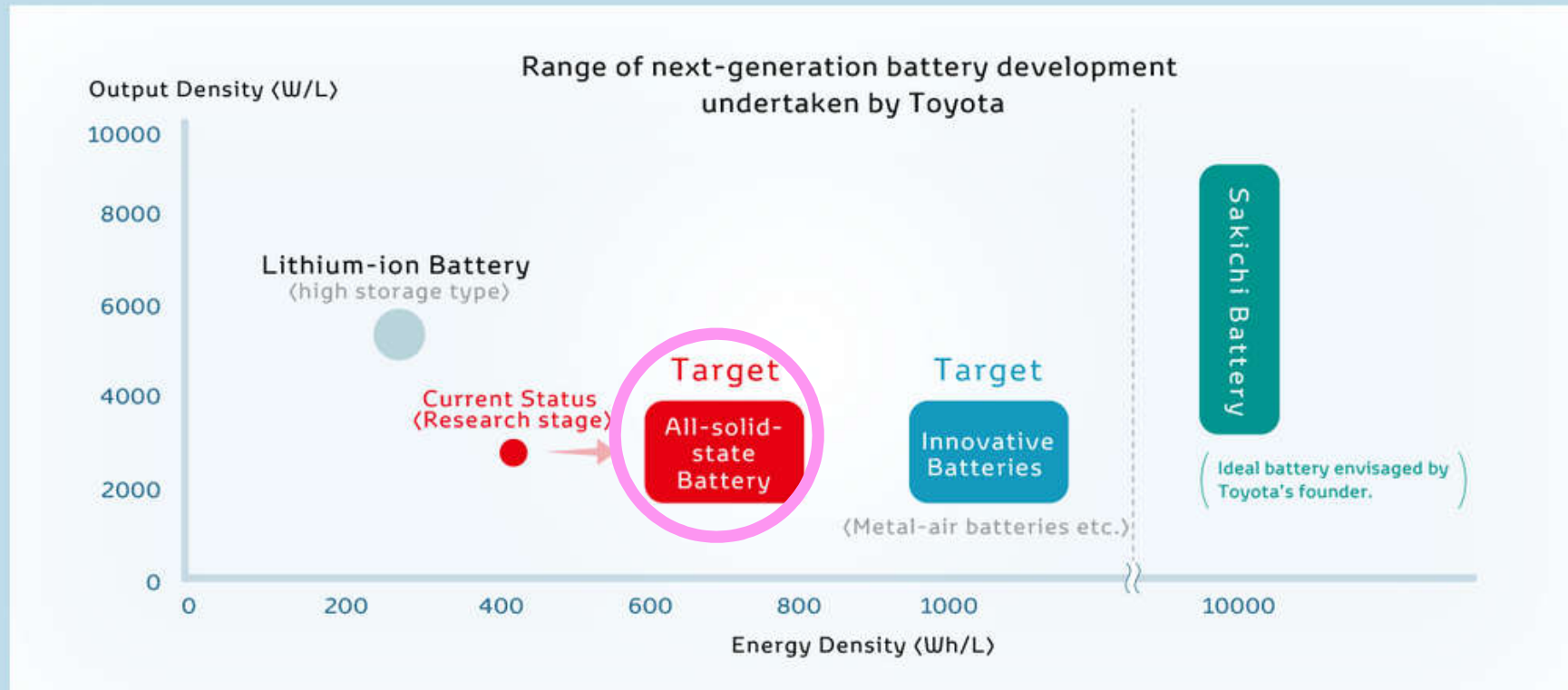
Motivation – Why solid electrolytes?

Materials components of a Li or Na ion battery



From Toyota Motor Company Website:

Expansion of battery development



Toyota is conducting wide-ranging development of next-generation batteries.

http://www.toyota-global.com/innovation/environmental_technology/keytech/

From Oak Ridge National Laboratory:

Solid Electrolyte: the Key for High-Voltage Lithium Batteries

Juchuan Li,* Cheng Ma, Miaofang Chi, Chengdu Liang, and Nancy J. Dudney*

Advantages

- Compatible and stable with high voltage cathodes and with Li metal anodes

Disadvantages

- Relatively low ionic conductivity (Compensated with the use of less electrolyte?)
- Lower total capacity

Demonstrated for $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4/\text{LiPON}/\text{Li}$

- 10^{-6} m LiPON electrolyte layer achieved adequate conductivity
- 10,000 cycles* with 90% capacity retention

*1 cycle per day for 27 years

Motivation: Paper by N. Kayama, *et. al* in *Nature Materials* **10**, 682-686 (2011)

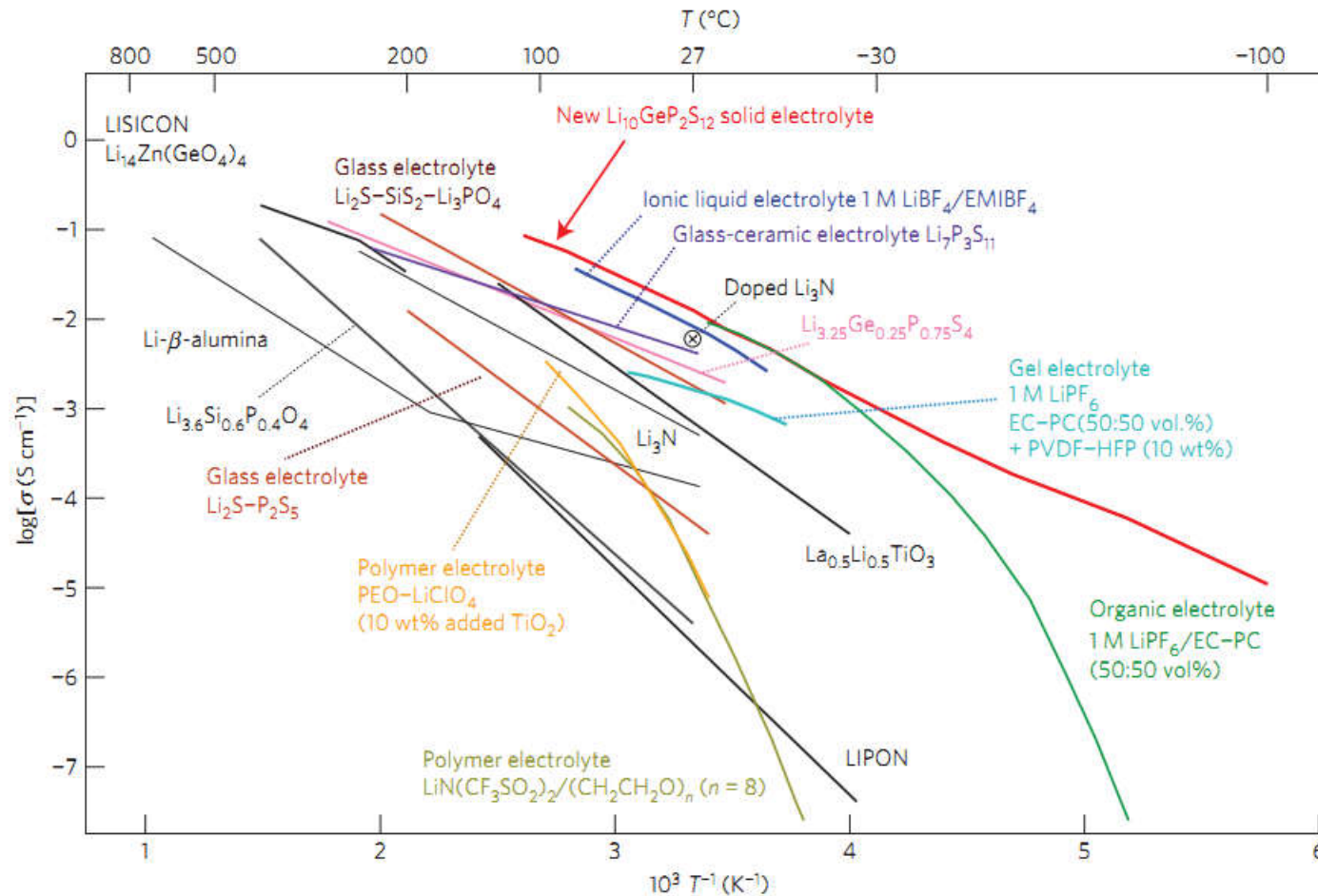


Figure 3 | Thermal evolution of ionic conductivity of the new $\text{Li}_{10}\text{GeP}_2\text{S}_{12}$ phase, together with those of other lithium solid electrolytes, organic liquid electrolytes, polymer electrolytes, ionic liquids and gel electrolytes^{3-8,13-16,20,22}. The new $\text{Li}_{10}\text{GeP}_2\text{S}_{12}$ exhibits the highest lithium ionic conductivity (12 m S cm^{-1} at 27 $^{\circ}\text{C}$) of the solid lithium conducting membranes of inorganic, polymer or composite systems. Because organic electrolytes usually have transport numbers below 0.5, inorganic lithium electrolytes have extremely high conductivities.

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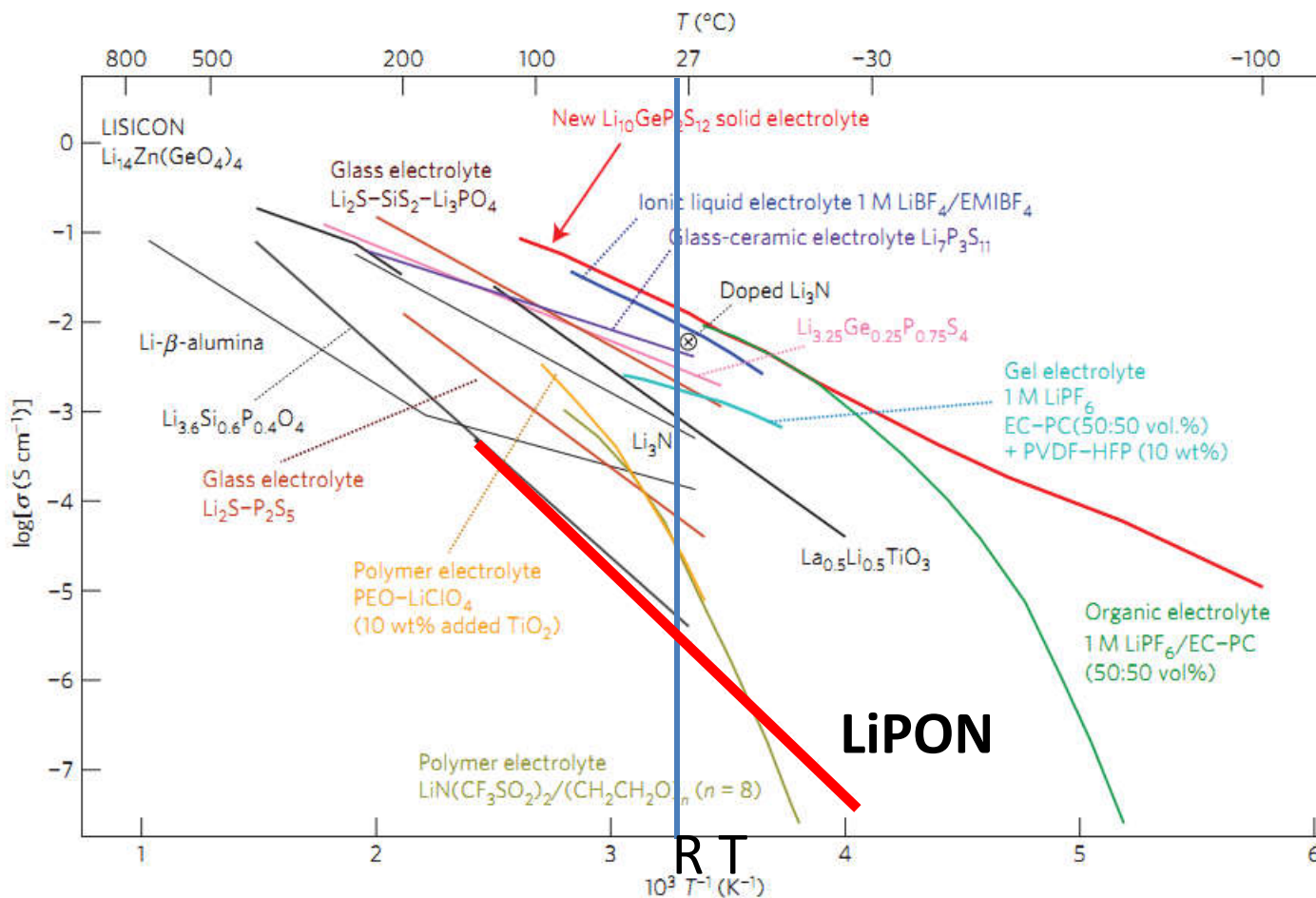


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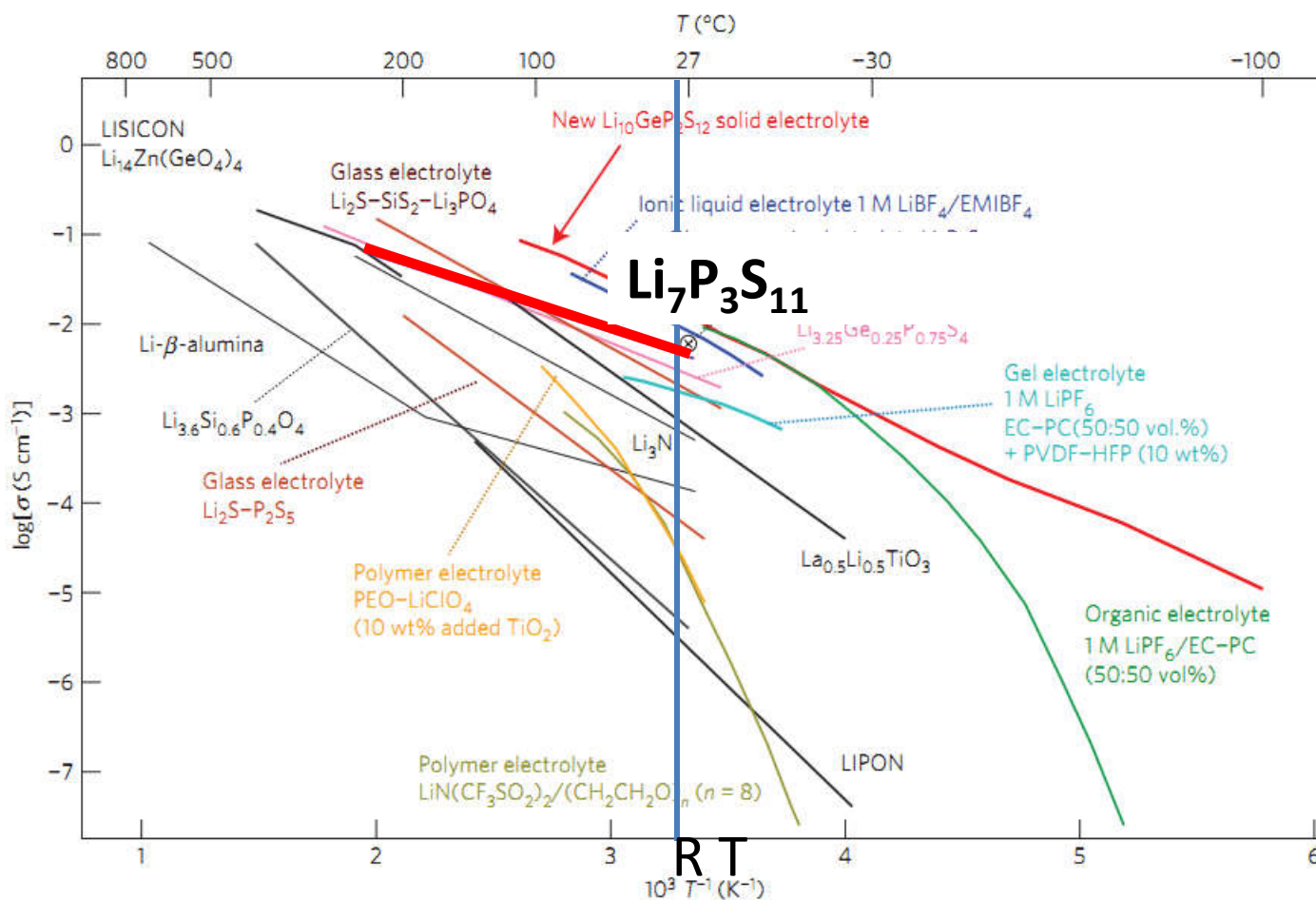


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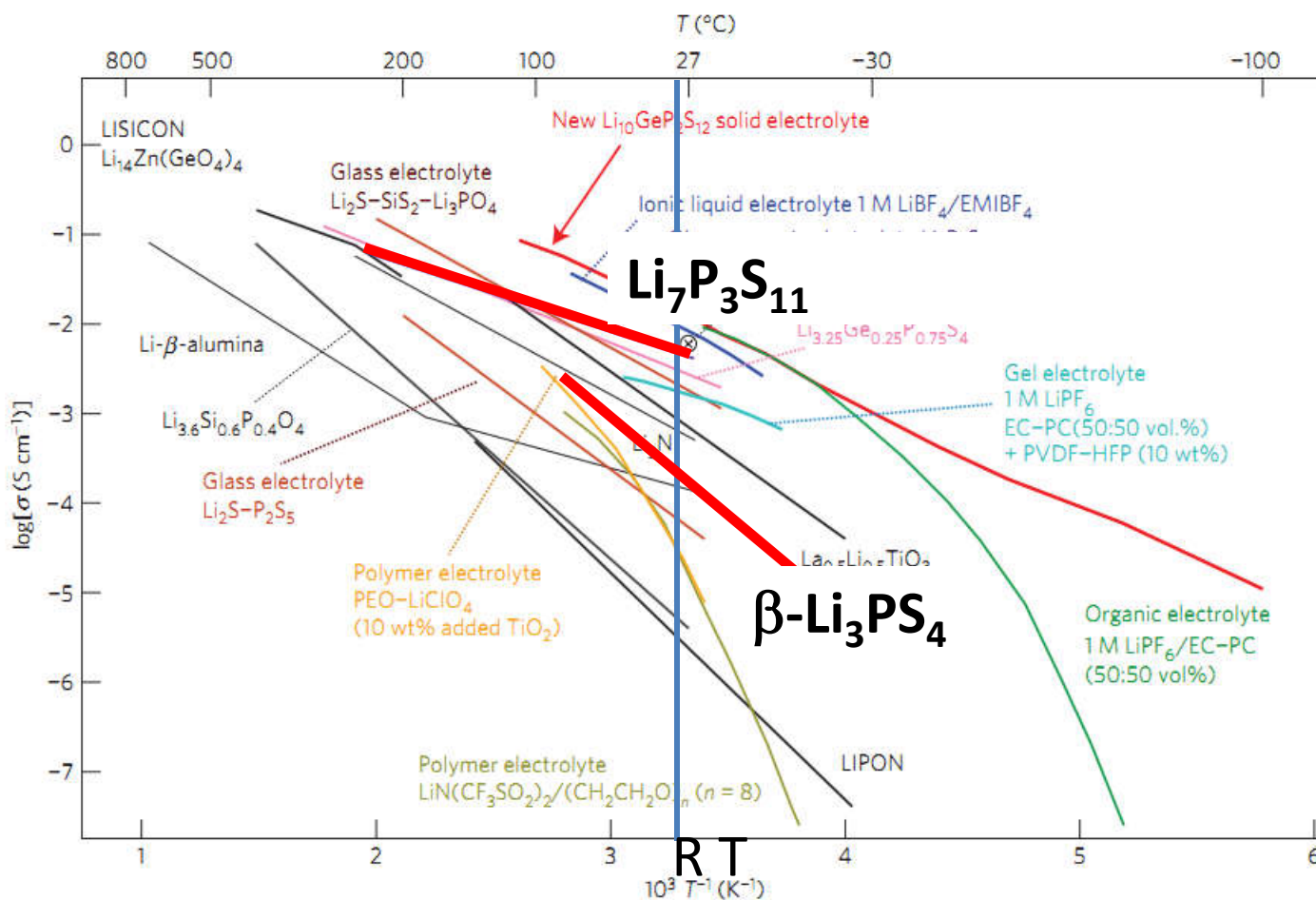


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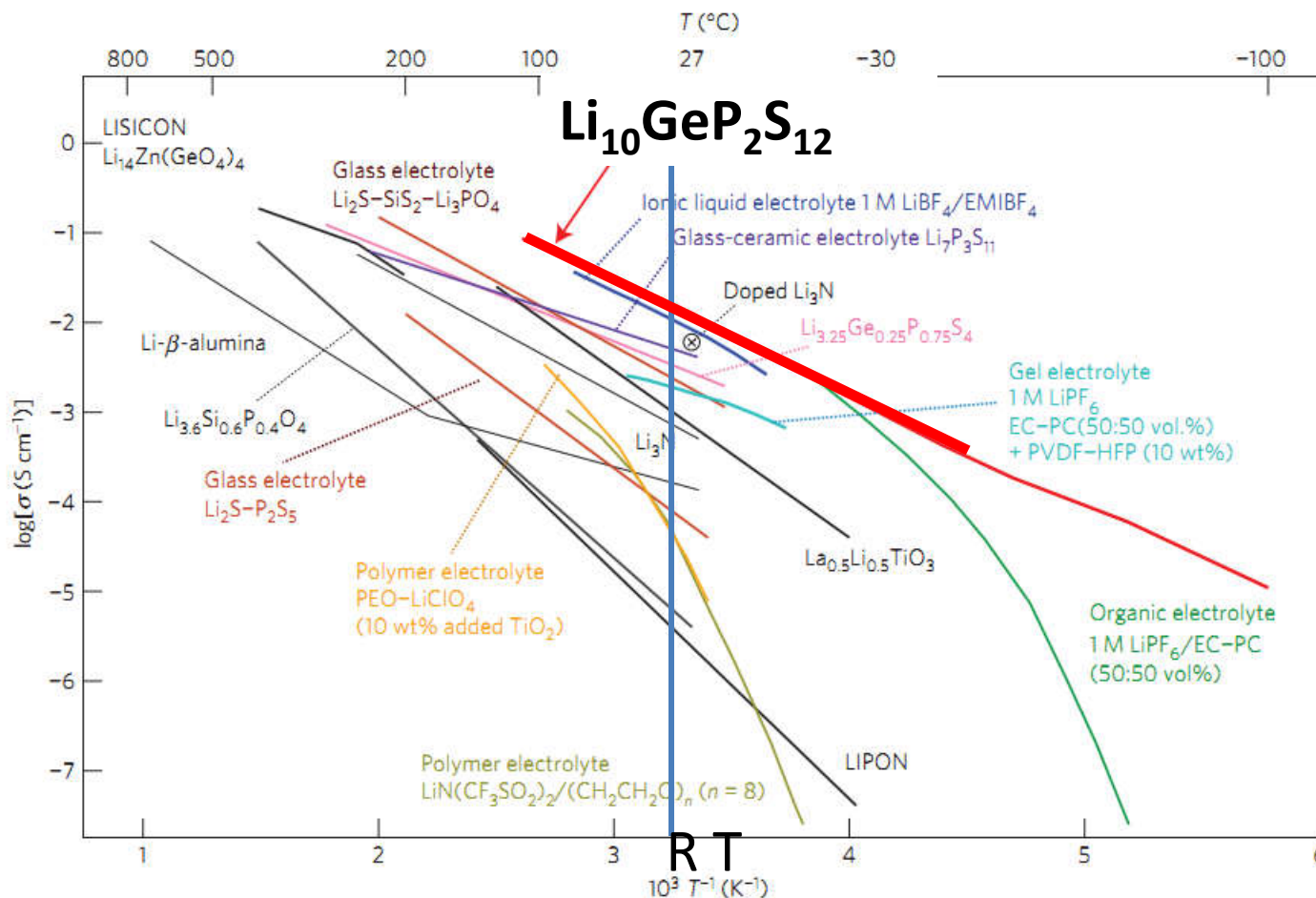
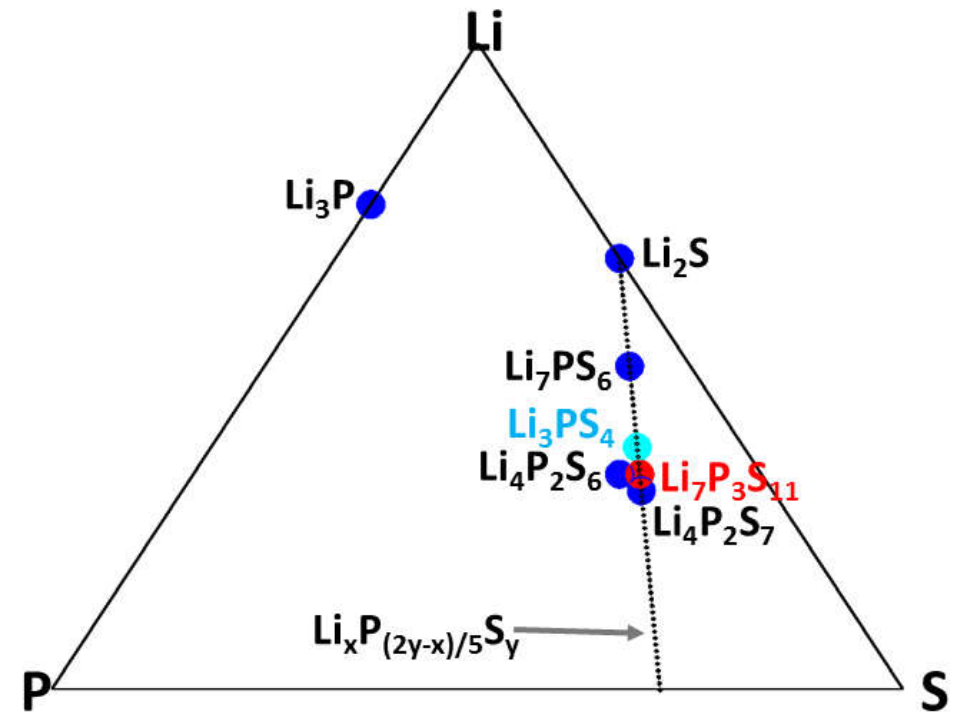
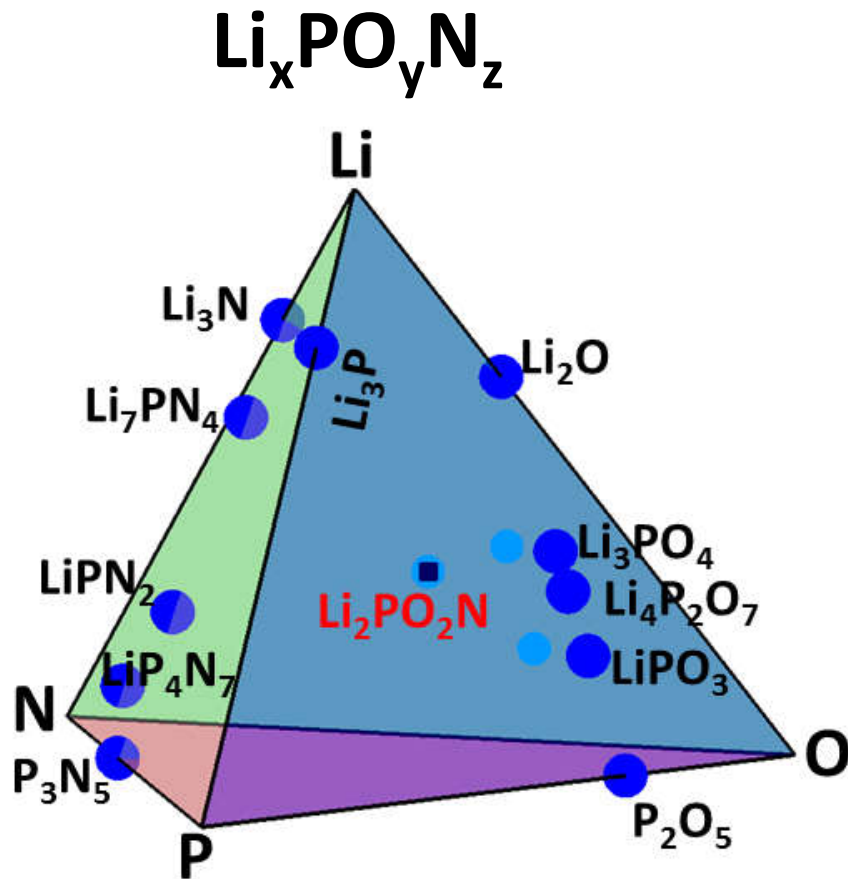


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Solid Electrolyte Families Investigated in this Study:



Computational tools

Summary of “first-principles” calculation methods

Exact Schrödinger equation:

$$\mathcal{H}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) \Psi_\alpha(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) = E_\alpha \Psi_\alpha(\{\mathbf{r}_i\}, \{\mathbf{R}^a\})$$

where

$$\mathcal{H}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) = \mathcal{H}^{\text{Nuclei}}(\{\mathbf{R}^a\}) + \mathcal{H}^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\})$$

Born-Oppenheimer approximation

Born & Huang, **Dynamical Theory of Crystal Lattices**, Oxford (1954)

Approximate factorization:

$$\Psi_\alpha(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) = X_\alpha^{\text{Nuclei}}(\{\mathbf{R}^a\}) \Upsilon_\alpha^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\})$$

Treated with classical mechanics

Treated with density
functional theory

Electronic Schrödinger equation:

$$\mathcal{H}^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) \Upsilon_{\alpha}^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) = U_{\alpha}(\{\mathbf{R}^a\}) \Upsilon_{\alpha}^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\})$$

$$\mathcal{H}^{\text{Electrons}}(\{\mathbf{r}_i\}, \{\mathbf{R}^a\}) = -\frac{\hbar^2}{2m} \sum_i \nabla_i^2 - \sum_{a,i} \frac{Z^a e^2}{|\mathbf{r}_i - \mathbf{R}^a|} + \sum_{i<j} \frac{e^2}{|\mathbf{r}_i - \mathbf{r}_j|}$$

For electronic ground state: $\alpha \Rightarrow 0$



Density functional theory

Hohenberg and Kohn, *Phys. Rev.* **136** B864 (1964)

Kohn and Sham, *Phys. Rev.* **140** A1133 (1965)

Mean field approximation: $U_0(\{\mathbf{R}^a\}) \Rightarrow U_0(\{\rho(\mathbf{r})\}, \{\mathbf{R}^a\})$ Electron density

Kohn-Sham construction: $\rho(\mathbf{r}) \approx \rho_{KS}(\mathbf{r}) = \sum_n |\psi_n(\mathbf{r})|^2$

$$\mathcal{H}_{KS}^{\text{Electrons}}(\mathbf{r}, \rho(\mathbf{r}), \{\mathbf{R}^a\}) \psi_n(\mathbf{r}) = \varepsilon_n \psi_n(\mathbf{r})$$

Independent electron wavefunction

More computational details:

$$\mathcal{H}_{\text{KS}}^{\text{Electrons}}(\mathbf{r}, \rho(\mathbf{r}), \{\mathbf{R}^a\}) = -\frac{\hbar^2 \nabla^2}{2m} + \underbrace{\sum_a \frac{-Z^a e^2}{|\mathbf{r} - \mathbf{R}^a|}}_{\text{electron-nucleus}} + \underbrace{e^2 \int d^3 r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}}_{\text{electron-electron}} + \underbrace{V_{xc}(\rho(\mathbf{r}))}_{\text{exchange-correlation}}$$

Exchange-correlation functionals:

LDA: J. Perdew and Y. Wang, Phys. Rev. B **45**, 13244 (1992)

GGA: J. Perdew, K. Burke, and M. Ernzerhof, PRL **77**, 3865 (1996)

HSE06: J. Heyd, G. E. Scuseria, and M. Ernzerhof, JCP **118**, 8207 (2003)

Numerical methods:

“Muffin-tin” construction: Augmented Plane Wave developed by Slater → “linearized” version by Andersen:

J. C. Slater, Phys. Rev. **51** 846 (1937)

O. K. Andersen, Phys. Rev. B **12** 3060 (1975) (LAPW)

Pseudopotential methods:

J. C. Phillips and L. Kleinman, Phys. Rev. **116** 287 (1959) -- original idea

P. Blöchl, Phys. Rev. B. **50** 17953 (1994) – Projector Augmented Wave (PAW) method

Outputs of calculations:

Ground state energy:

$$U_0(\{\rho(\mathbf{r})\}, \{\mathbf{R}^a\}) \Rightarrow \text{Determine formation energies}$$

$$\min_{\{\mathbf{R}^a\}} (U_0(\{\rho(\mathbf{r})\}, \{\mathbf{R}^a\})) \Rightarrow \text{Determine structural parameters}$$

\Rightarrow Stable and meta-stable structures

$$\rho_{KS}(\mathbf{r}) = \sum_n |\psi_n(\mathbf{r})|^2 \Rightarrow \text{Self-consistent electron density}$$

$$\{\varepsilon_n\} \Rightarrow \text{One-electron energies; densities of states}$$

Nuclear Hamiltonian (usually treated classically)

$$\mathcal{H}^{\text{Nuclei}}(\{\mathbf{R}^a\}) = \sum_a \frac{\mathbf{P}^{a2}}{2M^a} + U_0(\{\rho(\mathbf{r})\}, \{\mathbf{R}^a\}) \rightarrow \text{Normal modes of vibration}$$

Codes used for calculations

Function	Code	Website
Generate atomic datasets	ATOMPAW	http://pwpaw.wfu.edu
DFT; optimize structure	PWscf abinit	http://www.quantum-espresso.org http://www.abinit.org
Structural visualization	XCrySDen VESTA	http://ww.xcrysden.org http://jp-minerals.org/vesta/en/

ATOMPAW Code for generating atomic datasets for PAW calculations

Holzwarth, Tackett, and Matthews, CPC 135 329 (2001) <http://pwpaw.wfu.edu>

ATOMPAW

INFO

DATASETS

CONTRIBUTERS

CONTACT INFO

NAWH Web

PHYSICS Web

WFU Web

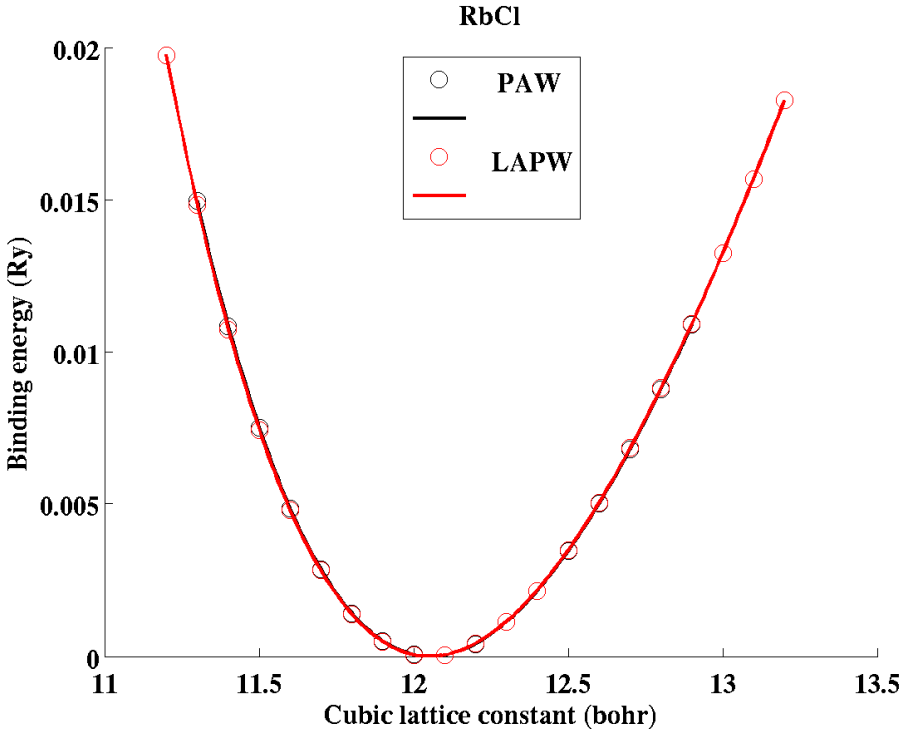
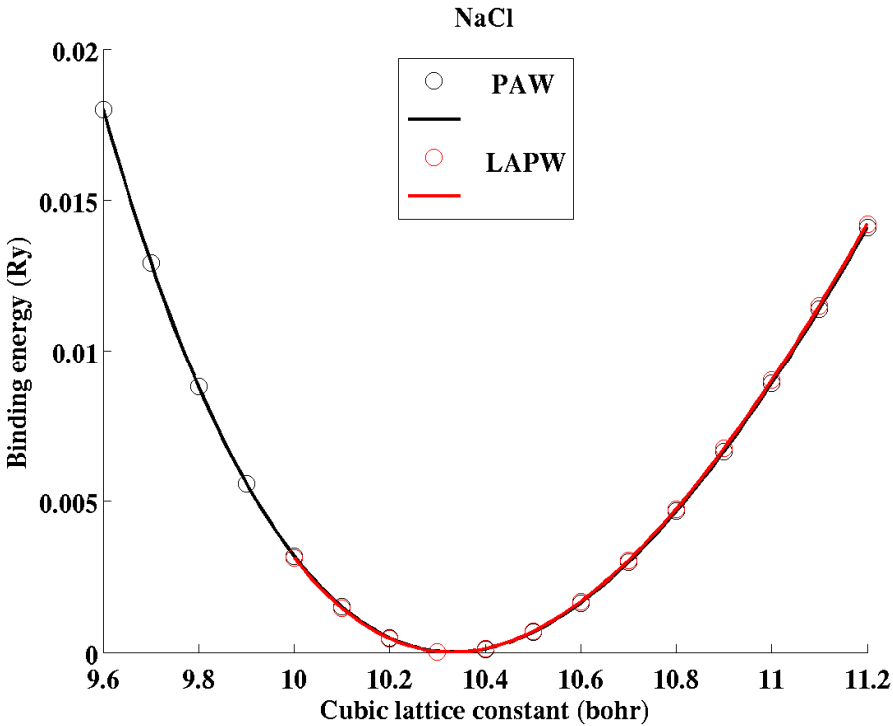
ATOMPAW

Download source code and example files:

- [atompaw-4.0.0.13.tar.gz](#) (5.4mb) 2/15/2016 Marc Torrent introduced improvements to XML output for abinit and new option for command line determination of version. MT also prepared some notes [in text format](#) or [in PDF format](#) to install atompaw under MacOS. In addition, NH made some slight changes to the UPF output for Quantum Espresso towards compatibility with libxc.
- [atompaw-3.1.0.3.tar.gz](#) (3.8mb) - January 2014 - Older version of atompaw with contributions from Marc Torrent and Francois Jollet as well as several others.
- [pwpaw_2.4.tgz](#) (0.2 mb) Updated 05/12/2010 version of *pwpaw* with very minor changes to accomodate changes to input files generated by new *atompaw* output files; also includes a BSD license file.

Atomic PAW datasets:

Comparison with LAPW results for binding energy curves --








Validation

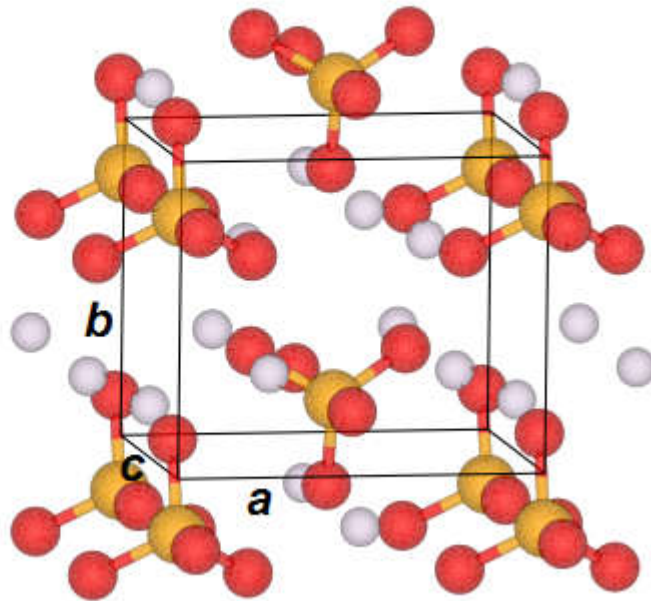
Li₃PO₄ crystals

γ -Li₃PO₄

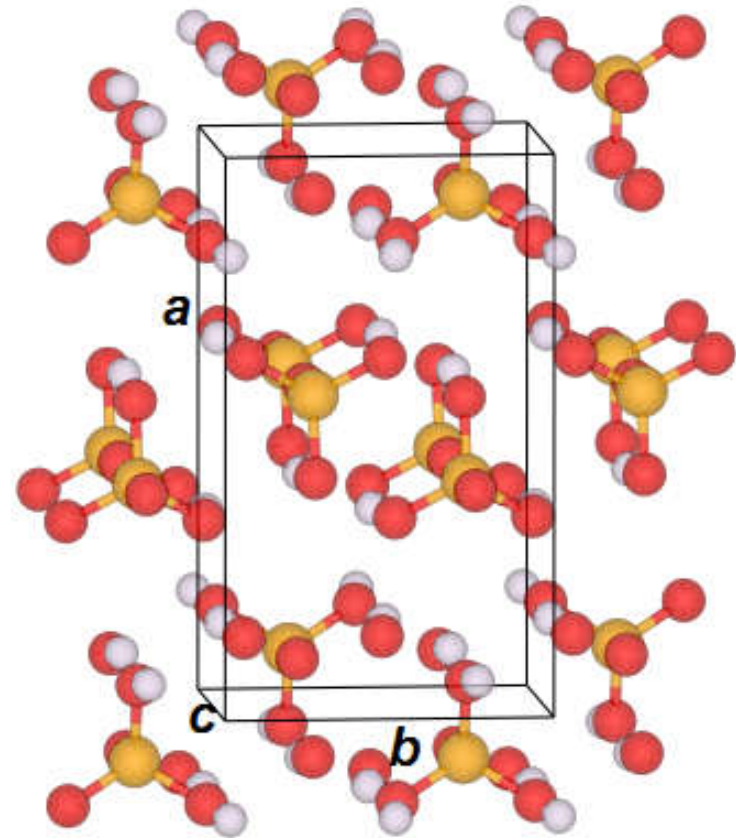
Key

-  Li
-  N
-  O
-  P
-  S

β -Li₃PO₄



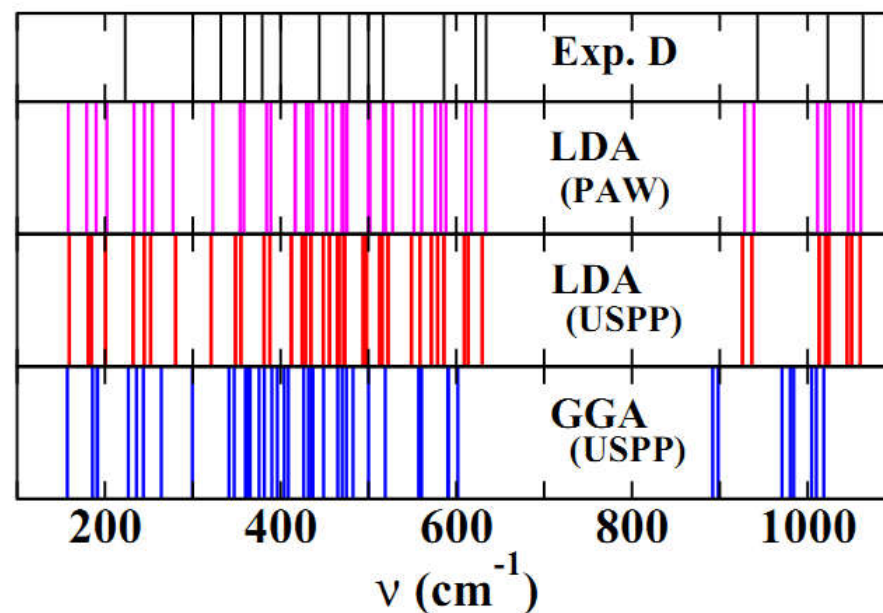
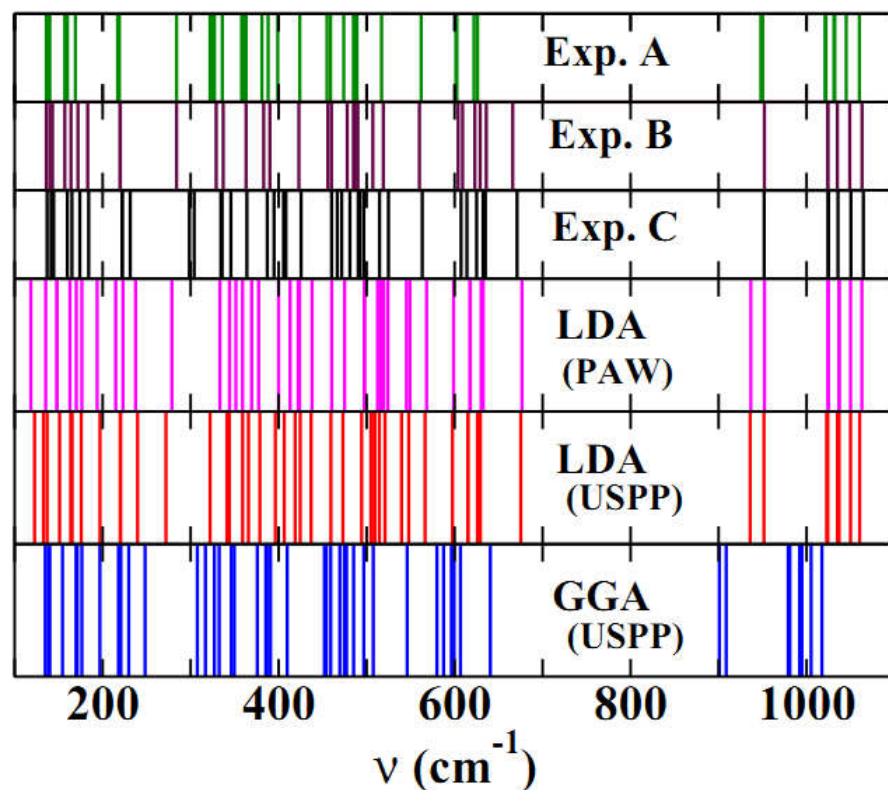
(*Pmn2*₁)



(*Pnma*)

Validation of calculations

Raman spectra – Experiment & Calculation



A: B. N. Mavrin et al, J. Exp. Theor. Phys. **96**,53 (2003); B: F. Harbach and F. Fischer, Phys. Status Solidi B **66**, 237 (1974) – room temp. C: Ref. B at liquid nitrogen temp.; D: L. Popović et al, J. Raman Spectrosc. **34**,77 (2003).

Heats of formation – Experiment & Calculation



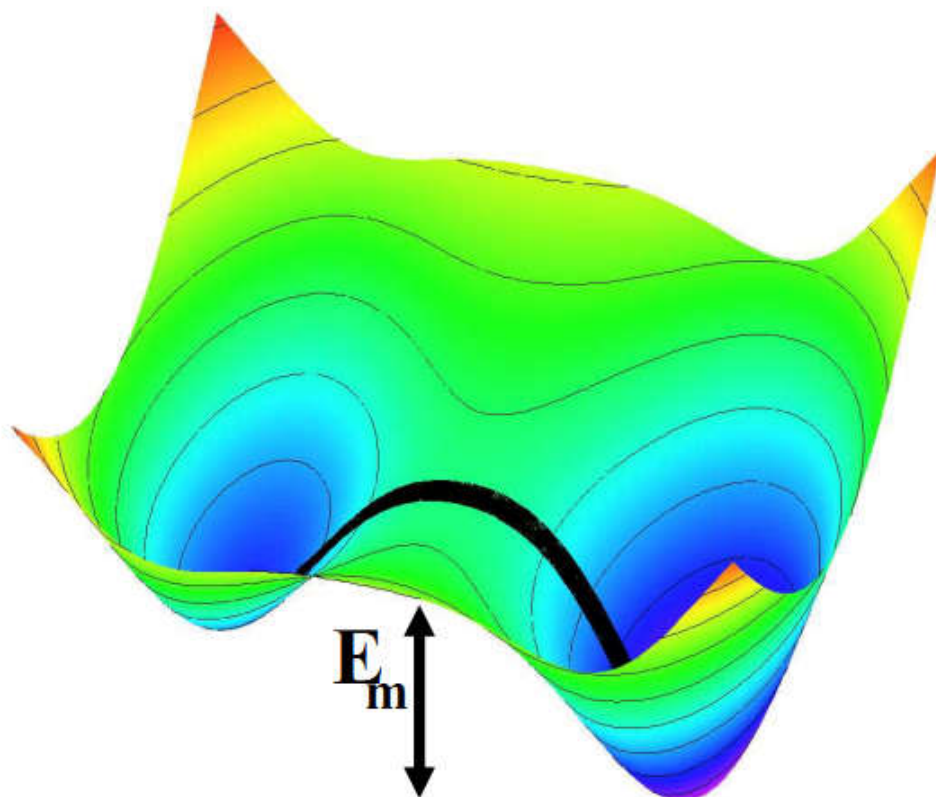
Table 1. Calculated heats of formation for Li phosphates, phospho-nitrides, and thiophosphates and related materials. The structural designation uses the notation defined in the International Table of Crystallography⁸⁵ based on structural information reported in the International Crystal Structure Database.⁸⁶ The heats of formation ΔH (eV/FU) are given in units of eV per formula unit. When available from Ref. [31] and [32] experiment values are indicated in parentheses. Those indicated with “*” were used fitting the O and N reference energies as explained in the text.

Material	Structure	ΔH (eV/FU)	Material	Structure	ΔH (eV/FU)
β -Li ₃ PO ₄	<i>Pmn</i> 2 ₁ (#31)	-21.23	N ₂ O ₅	<i>P6</i> ₃ / <i>mmc</i> (#194)	- 0.94 (- 0.45*)
γ -Li ₃ PO ₄	<i>Pnma</i> (#62)	-21.20 (-21.72*)	P ₃ N ₅	<i>C2/c</i> (#15)	- 3.02 (- 3.32*)
γ -Li ₃ PS ₄	<i>Pmn</i> 2 ₁ (#31)	- 8.37	<i>h</i> -P ₂ O ₅	<i>R3c</i> (#161)	-15.45 (-15.53*)
β -Li ₃ PS ₄	<i>Pnma</i> (#62)	- 8.28	α -P ₂ O ₅	<i>Fdd</i> 2 (#43)	-15.78
Li ₄ P ₂ O ₆	<i>P</i> $\bar{3}1m (#162)$	-29.72	P ₂ S ₅	<i>P</i> $\bar{1}$ (#2)	- 1.93
Li ₄ P ₂ O ₇	<i>P</i> $\bar{1}$ (#2)	-33.97	P ₄ S ₃	<i>Pnma</i> (#62)	- 2.45 (- 2.33)
Li ₅ P ₂ O ₆ N	<i>P</i> $\bar{1}$ (#2)	-33.18	SO ₃	<i>Pna</i> 2 ₁ (#33)	- 4.84 (- 4.71*)
Li ₄ P ₂ S ₆	<i>P</i> $\bar{3}$ 1 <i>m</i> (#162)	-12.42	Li ₃ N	<i>P6/mmm</i> (#191)	- 1.60 (- 1.71*)
Li ₄ P ₂ S ₇	<i>P</i> $\bar{1}$ (#2)	-11.59	Li ₂ O	<i>Fm</i> $\bar{3}$ <i>m</i> (#225)	- 6.10 (- 6.20*)
Li ₇ P ₃ O ₁₁	<i>P</i> $\bar{1}$ (#2)	-54.84	Li ₂ O ₂	<i>P6</i> ₃ / <i>mmc</i> (#194)	- 6.35 (- 6.57*)
Li ₇ P ₃ S ₁₁	<i>P</i> $\bar{1}$ (#2)	-20.01	Li ₃ P	<i>P6</i> ₃ / <i>mmc</i> (#194)	- 3.47
LiPO ₃	<i>P2/c</i> (#13)	-12.75	Li ₂ S	<i>Fm</i> $\bar{3}$ <i>m</i> (#225)	- 4.30 (- 4.57)
LiPN ₂	<i>I</i> $\bar{4}$ 2 <i>d</i> (#122)	- 3.65	Li ₂ S ₂	<i>P6</i> ₃ / <i>mmc</i> (#194)	- 4.09
<i>s1</i> -Li ₂ PO ₂ N	<i>Pbcm</i> (#57)	-12.35	LiNO ₃	<i>R</i> $\bar{3}$ <i>c</i> (#167)	- 5.37 (- 5.01*)
<i>SD</i> -Li ₂ PO ₂ N	<i>Cmc</i> 2 ₁ (#36)	-12.47	Li ₂ SO ₄	<i>P2</i> ₁ / <i>c</i> (#14)	-14.63 (-14.89*)
<i>SD</i> -Li ₂ PS ₂ N	<i>Cmc</i> 2 ₁ (#36)	- 5.80			

Estimate of ionic conductivity assuming activated hopping

Schematic diagram of minimal energy path

Approximated using NEB algorithm^a
– “Nudged Elastic Band”



^aHenkelman and Jónsson, *JCP* 113, 9978 (2000)

Arrhenius relation

$$\sigma \cdot T = K e^{-E_A/kT}$$

From: Ivanov-Shitz and co-workers,
Cryst. Reports 46, 864 (2001):

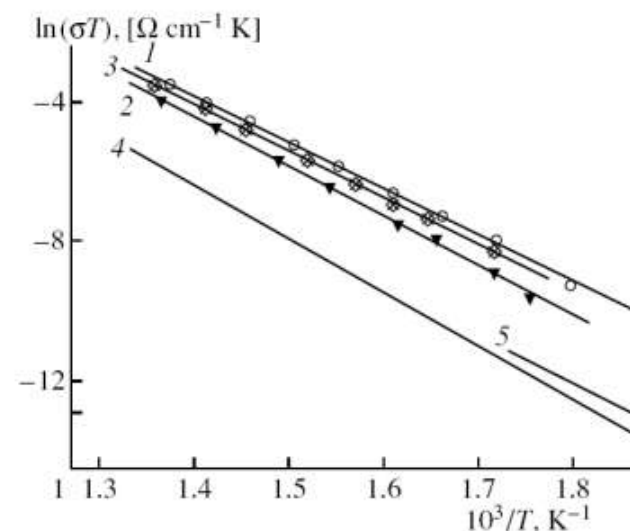


Fig. 2. Temperature dependences of conductivity in $\gamma\text{-Li}_3\text{PO}_4$; (1-3) for single crystals measured along the (1) *a*-axis, (2) *b*-axis, (3) *c*-axis and (4, 5) for a polycrystal (4) according to [4, 5] and (5) according to [7].

$E_A = 1.14, 1.23, 1.14, 1.31, 1.24$ eV for 1,2,3,4,5, respectively.

Arrhenius activation energies – Experiment and Calculation

Table 3. Calculated migration energies (E_m^{cal}) for Li ion vacancies (vac) and interstitials (int), vacancy-interstitial formation energies (E_f^{cal}), and corresponding the activation energies (E_A^{cal}) for crystalline materials computed using the NEB method in idealized supercells. When available, experimental activation energies E_A^{exp} are also listed together with additional information including the literature reference indicated in [] brackets. For γ - Li_3PO_4 , results for different crystallographic directions are quoted to compare with single crystal experiment; in other cases, only the minimum energies are given. All energies are given in eV.

Material	vac E_m^{cal}	int E_m^{cal}	E_f^{cal}	$E_A^{\text{cal}} \approx E_m^{\text{cal}} + \frac{1}{2} E_f^{\text{cal}}$		Reference
				E_A^{cal}	E_A^{exp}	
β - Li_3PO_4	0.7	0.4	2.1	1.4		
γ - Li_3PO_4	0.7, 0.7	0.4, 0.3	1.7	1.3, 1.1	1.23, 1.14	(sngl. cryst.) [100]
$\text{Li}_{2.88}\text{PO}_{3.73}\text{N}_{0.14}$					0.97	(poly cryst.) [58]
$\text{Li}_{3.3}\text{PO}_{3.9}\text{N}_{0.17}$					0.56	(amorphous) [58]
$\text{Li}_{1.35}\text{PO}_{2.99}\text{N}_{0.13}$					0.60	(amorphous) [101]
LiPO_3	0.6	0.7	1.2	1.2	1.4	(poly cryst.) [97]
					0.76-1.2	(amorphous) [97]
LiPN_2	0.4		2.5	1.7	0.6	(poly cryst.) [99]
$SD\text{-Li}_2\text{PO}_2\text{N}$	0.4	0.8	2.0	1.4	0.6	(poly cryst.) [52]
γ - Li_3PS_4	0.3		0.8	0.7	0.5	(poly cryst.) [102]
β - Li_3PS_4	0.2		0.0	0.2	0.4	(nano cryst.) [103]
$\text{Li}_7\text{P}_3\text{S}_{11}$	0.2	0.5	0.0	0.2	0.1	(poly cryst.) [76]

➤ What is meant by “first principles”?

A series of well-controlled approximations

- Born-Oppenheimer Approximation
- Density Functional Approximation
- Local density Approximation (LDA)
- Numerical method: Projector Augmented Wave

Validation

- Lattice vibration modes
- Heats of formation
- Activation energies for lattice migration

How can computer simulations contribute to the development of materials?

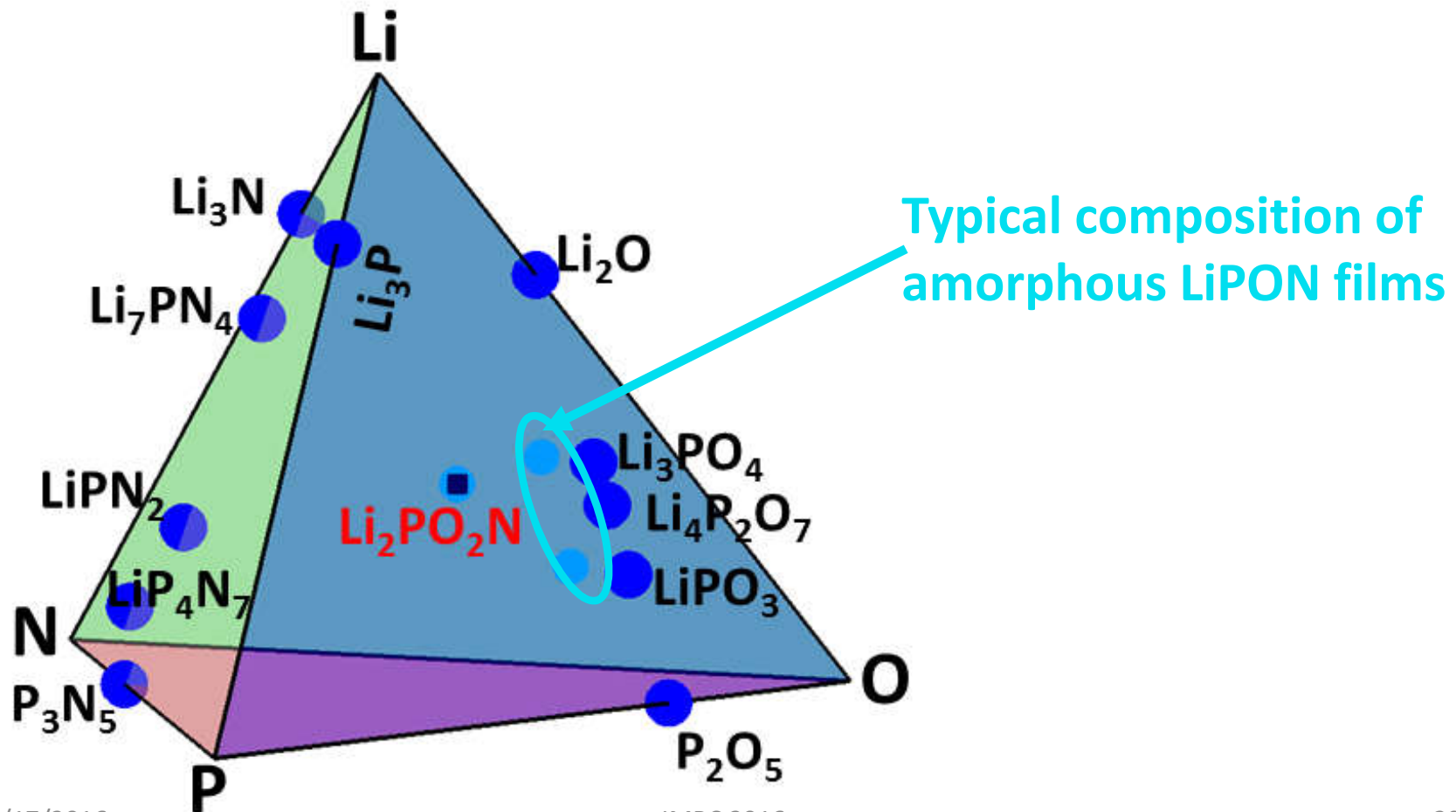


- Computationally examine known materials and predict new materials and their properties
 - Structural forms
 - Relative stabilities
 - Direct comparisons of simulations and experiment
 - Investigate properties that are difficult to realize experimentally

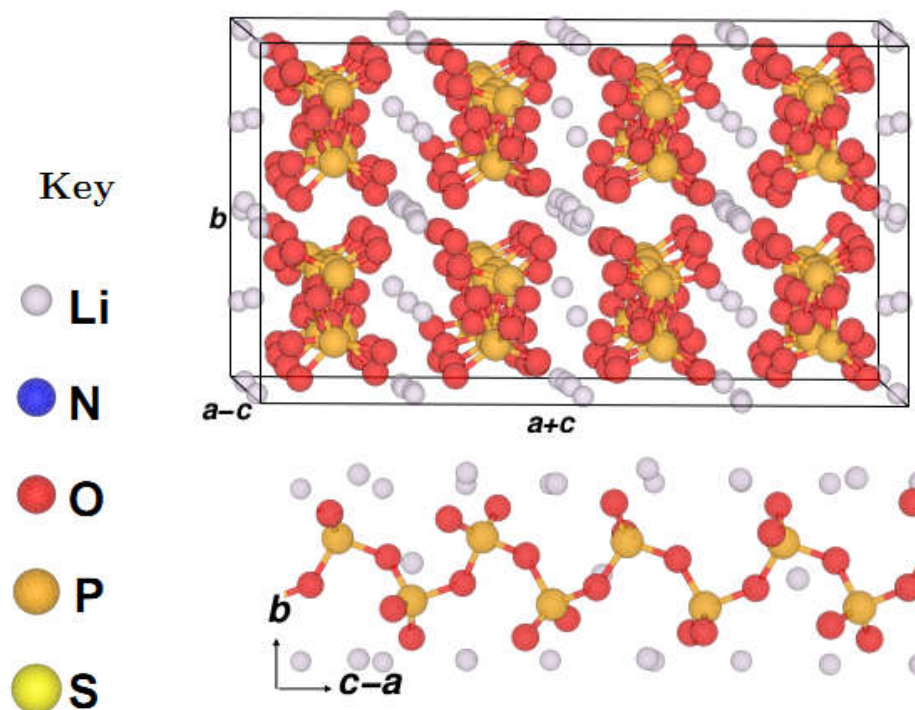
Of particular interest in battery materials --

- Model ion migration mechanisms
 - Vacancy migration
 - Interstitial migration
 - Vacancy-interstitial formation energies
- Model ideal electrolyte interfaces with anodes

**Systematic study of LiPON materials – $\text{Li}_x\text{PO}_y\text{N}_z$ –
(Yaojun A. Du and N. A. W. Holzwarth, Phys. Rev. B 81, 184106 (2010))**



Experimentally known structure



Computationally predicted structure

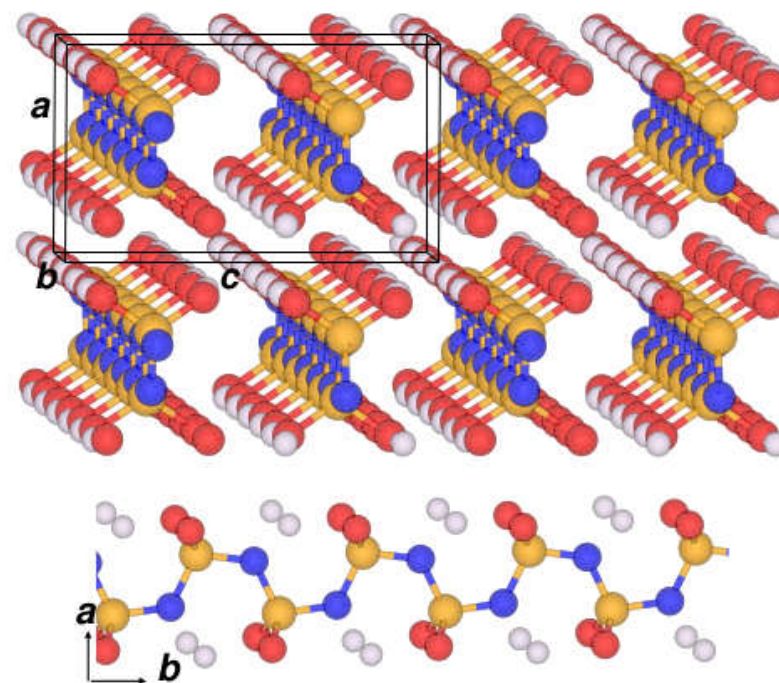


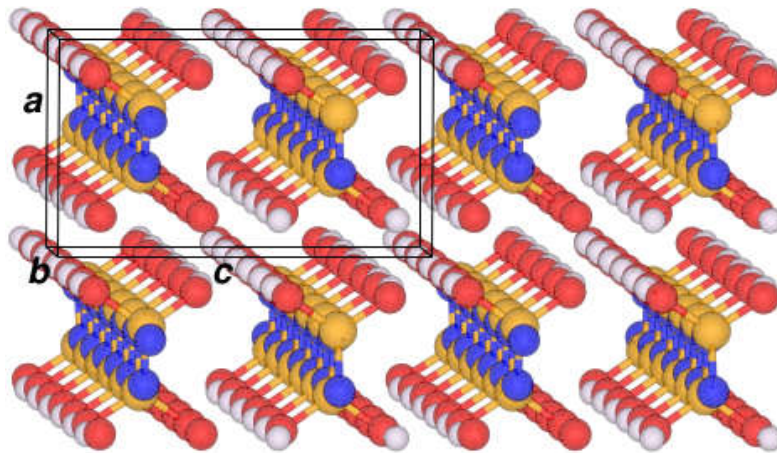
Fig. 7. Ball and stick diagrams for LiPO_3 in the $P2/c$ structure (20 formula units per unit cell) and $s_1\text{-Li}_2\text{PO}_2\text{N}$ in the $Pbcm$ structure (4 formula units per unit cell) from the calculated results. For each crystal diagram, a view of a horizontal chain axis is also provided for a single phosphate or phospho-nitride chain.

Computationally predicted structure

Key

s_1 -Li₂PO₂N *Pbcm*

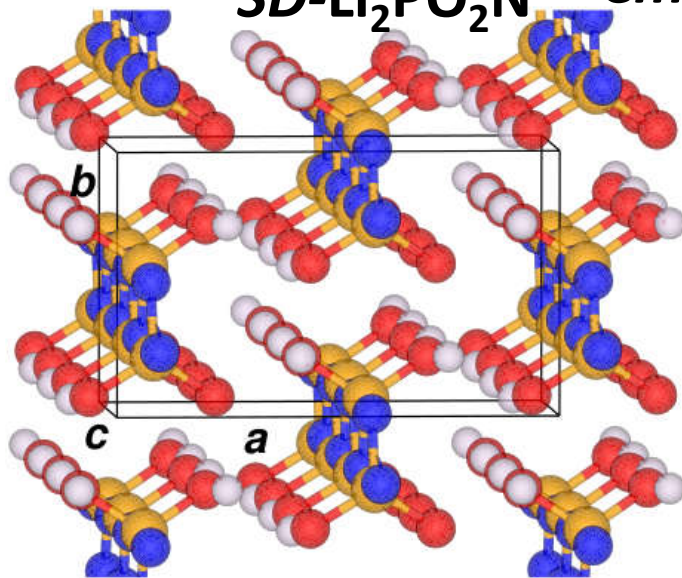
- Li
- N
- O
- P
- S



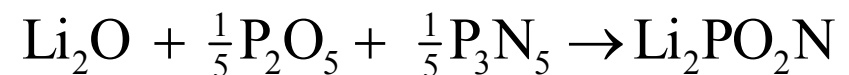
Calculations have now verified that the *SD* structure is more stable than the s_1 structure by 0.1 eV/FU.

Experimentally realized structure

SD-Li₂PO₂N *Cmc2₁*

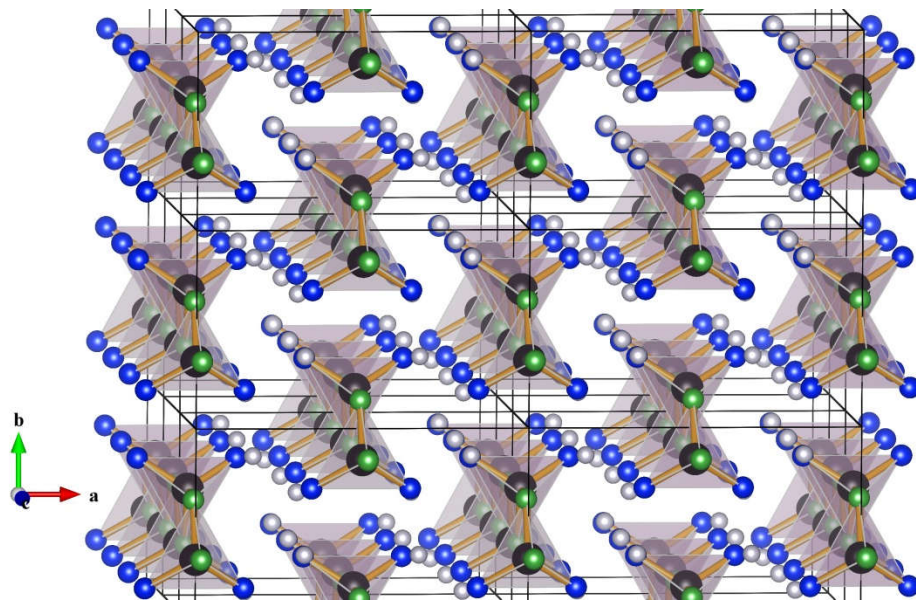


Synthesis of Li₂PO₂N by Keerthi Senevirathne, Cynthia Day, Michael Gross, and Abdessadek Lachgar (SSI 233, 95-101 (2013))
High temperature solid state synthesis using reaction:



Comparison of synthesized $\text{Li}_2\text{PO}_2\text{N}$ with Li_2SiO_3

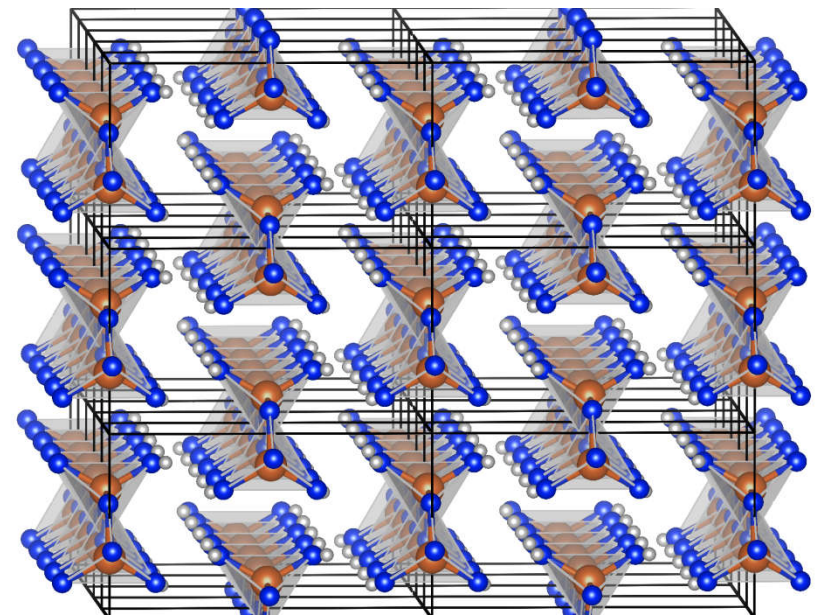
$\text{SD-Li}_2\text{PO}_2\text{N}$ ($Cmc2_1$)



$a=9.07 \text{ \AA}$, $b=5.40 \text{ \AA}$, $c=4.60 \text{ \AA}$

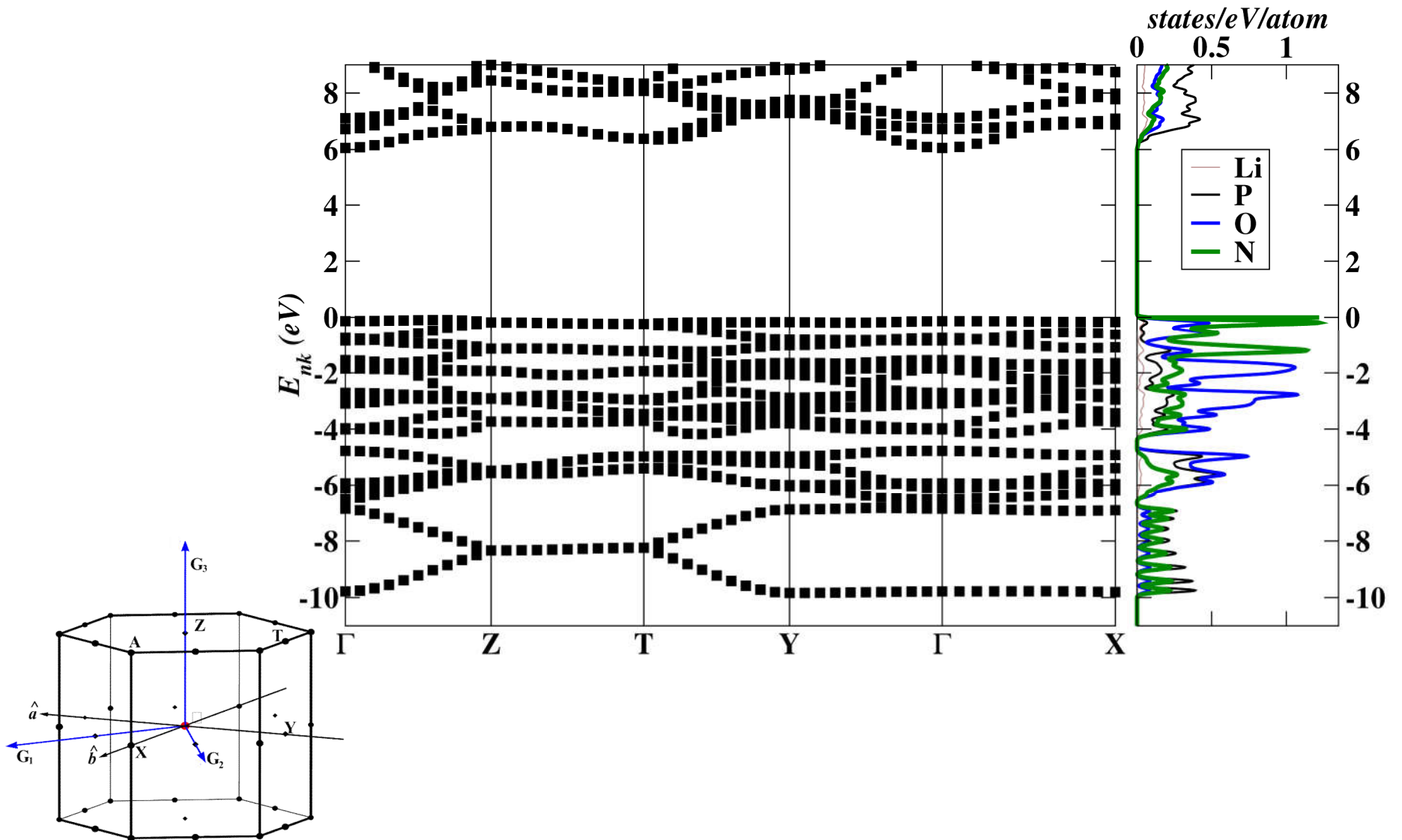


Li_2SiO_3 ($Cmc2_1$)



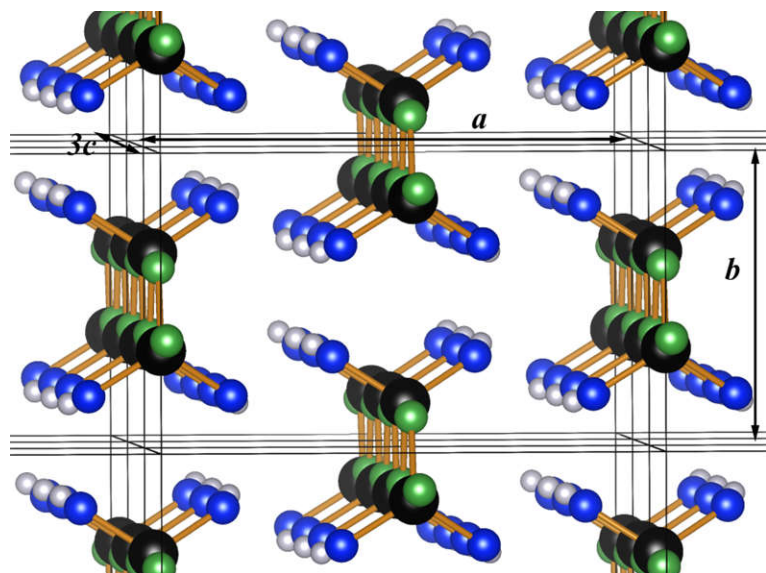
$a=9.39 \text{ \AA}$, $b=5.40 \text{ \AA}$, $c=4.66 \text{ \AA}$
K.-F. Hesse, Acta Cryst. B33, 901 (1977)

Electronic band structure of $SD\text{-Li}_2\text{PO}_2\text{N}$

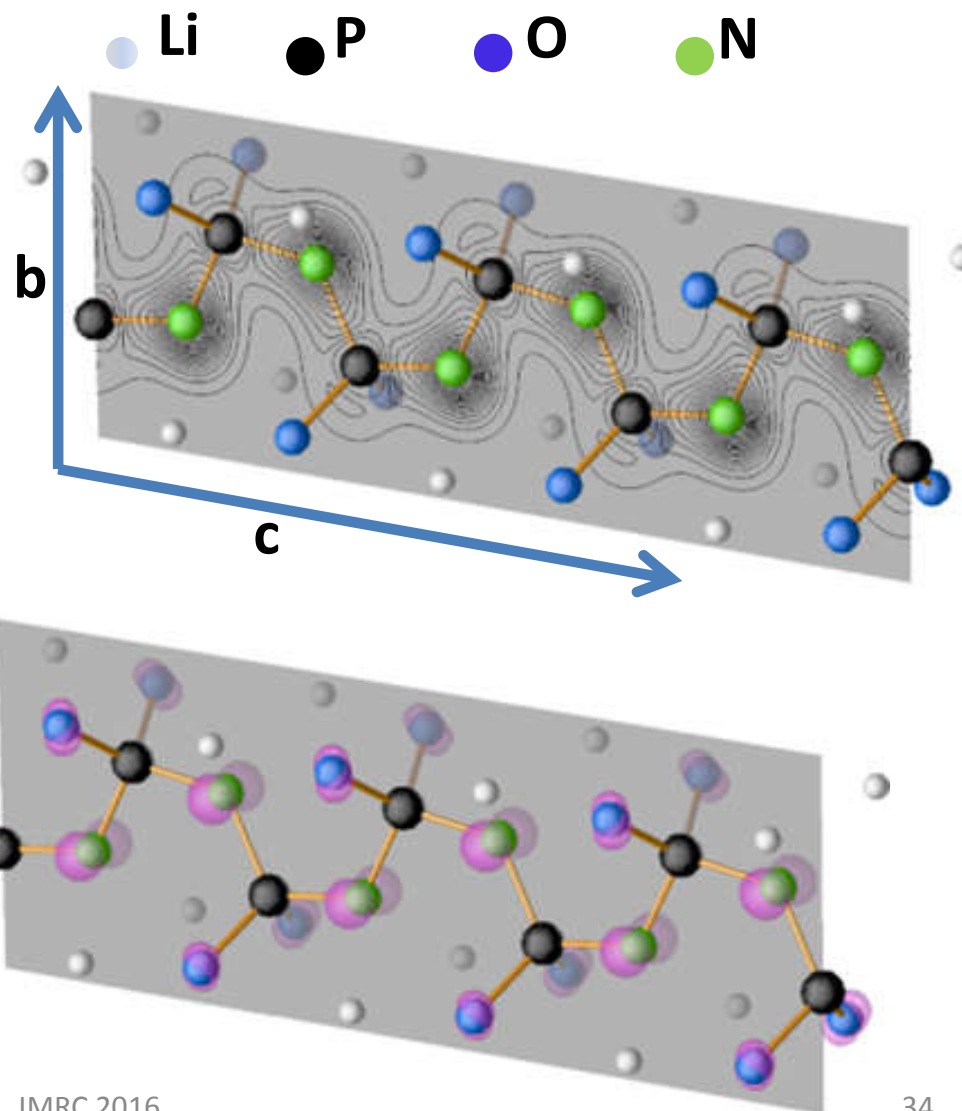


More details of *SD*-Li₂PO₂N structure

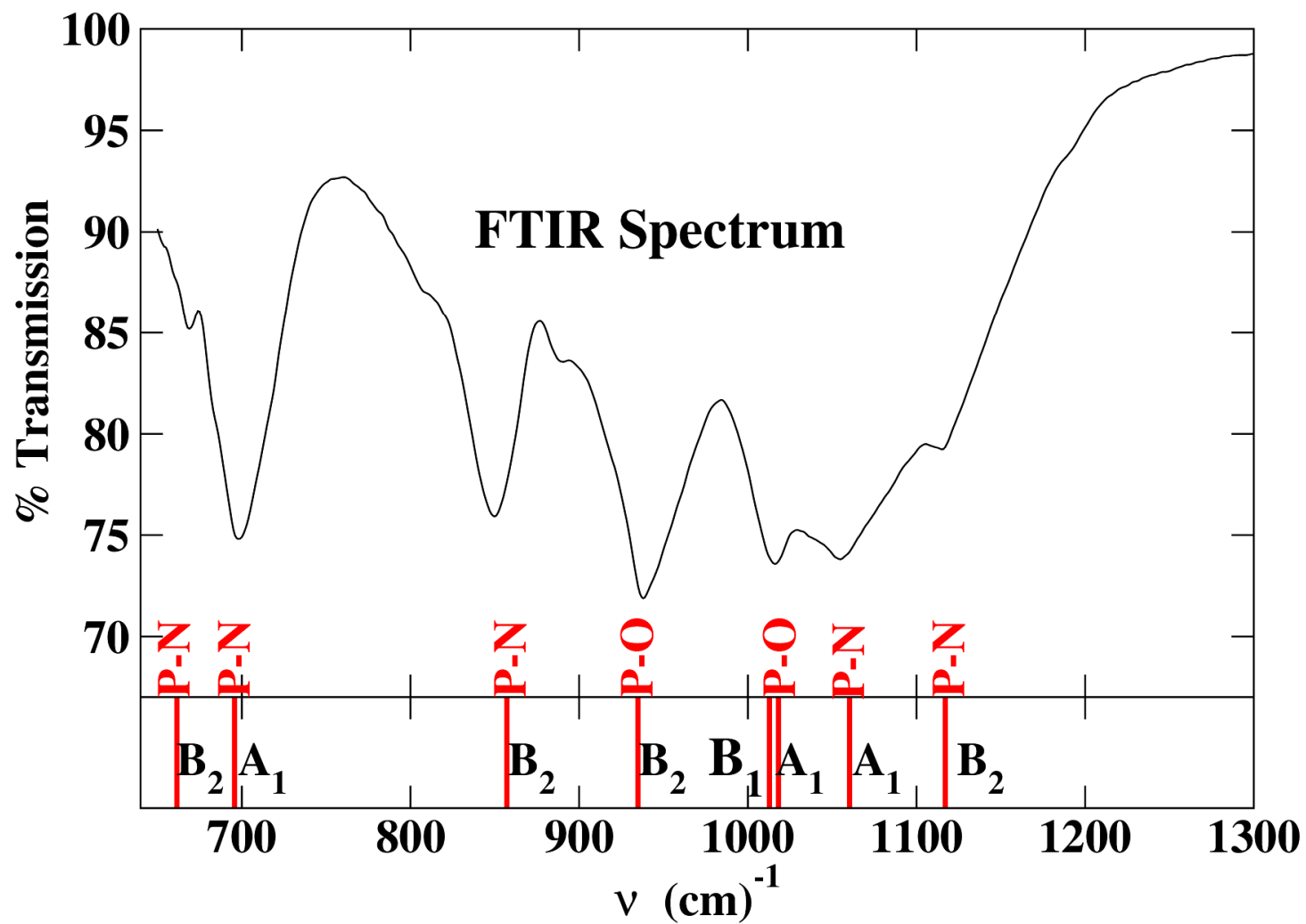
Ball and stick model



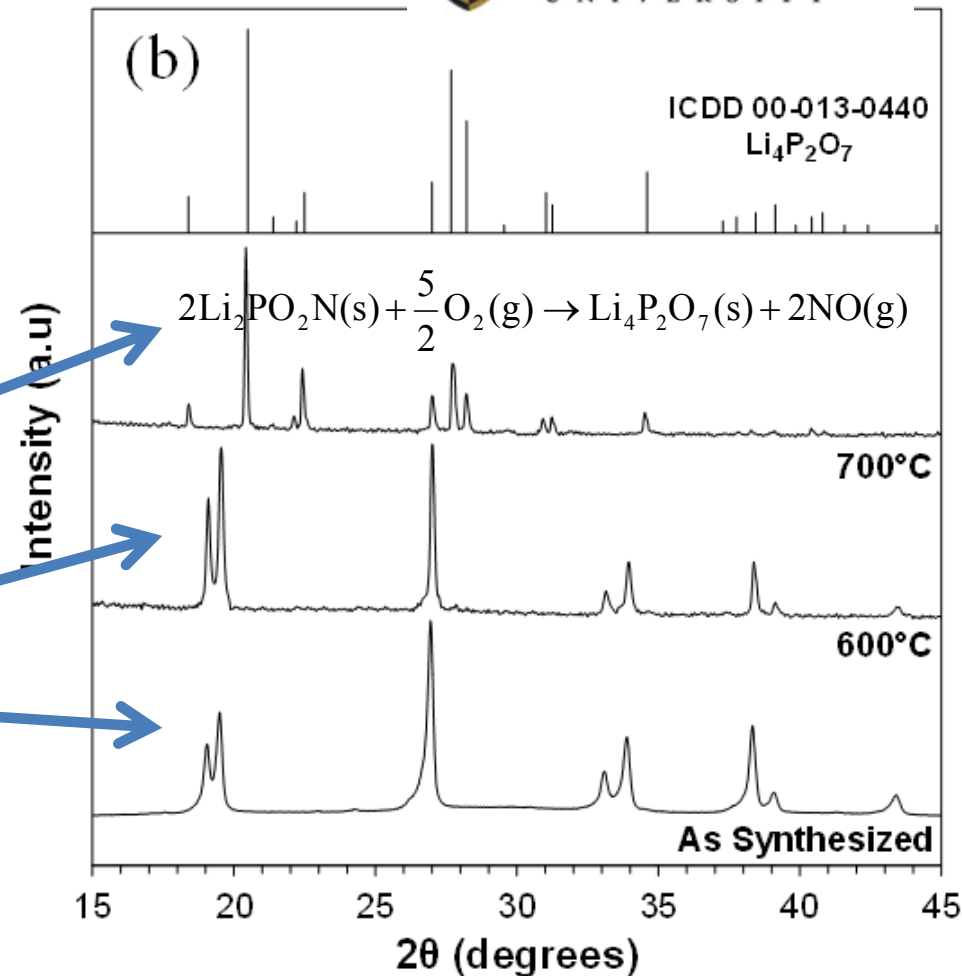
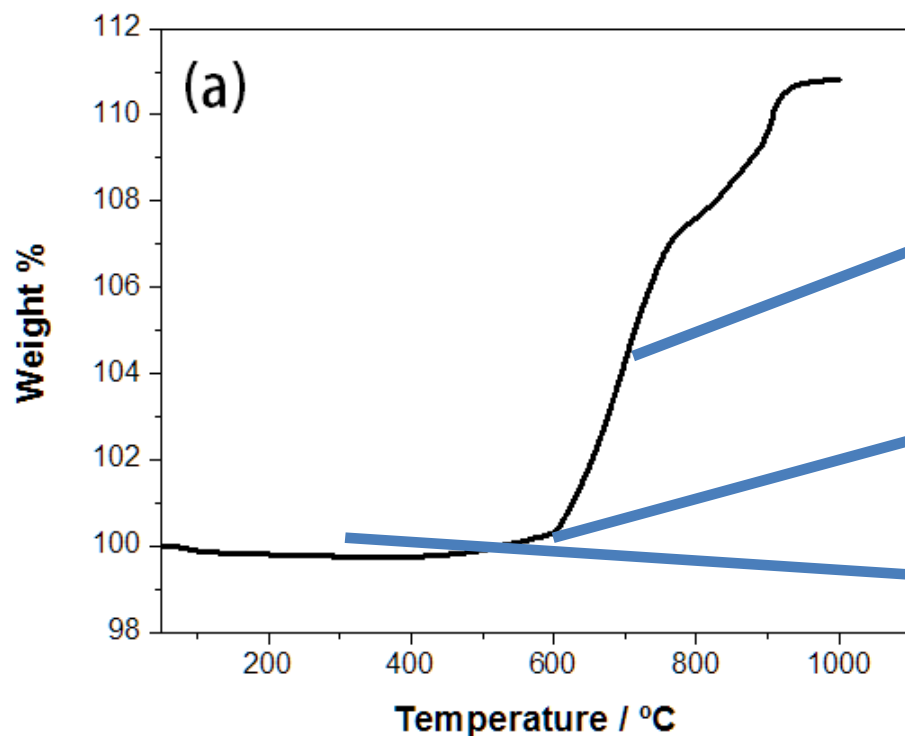
Isosurfaces (maroon) of charge density of states at top of valence band, primarily π states on N.



Vibrational spectrum of $SD\text{-Li}_2\text{PO}_2\text{N}$



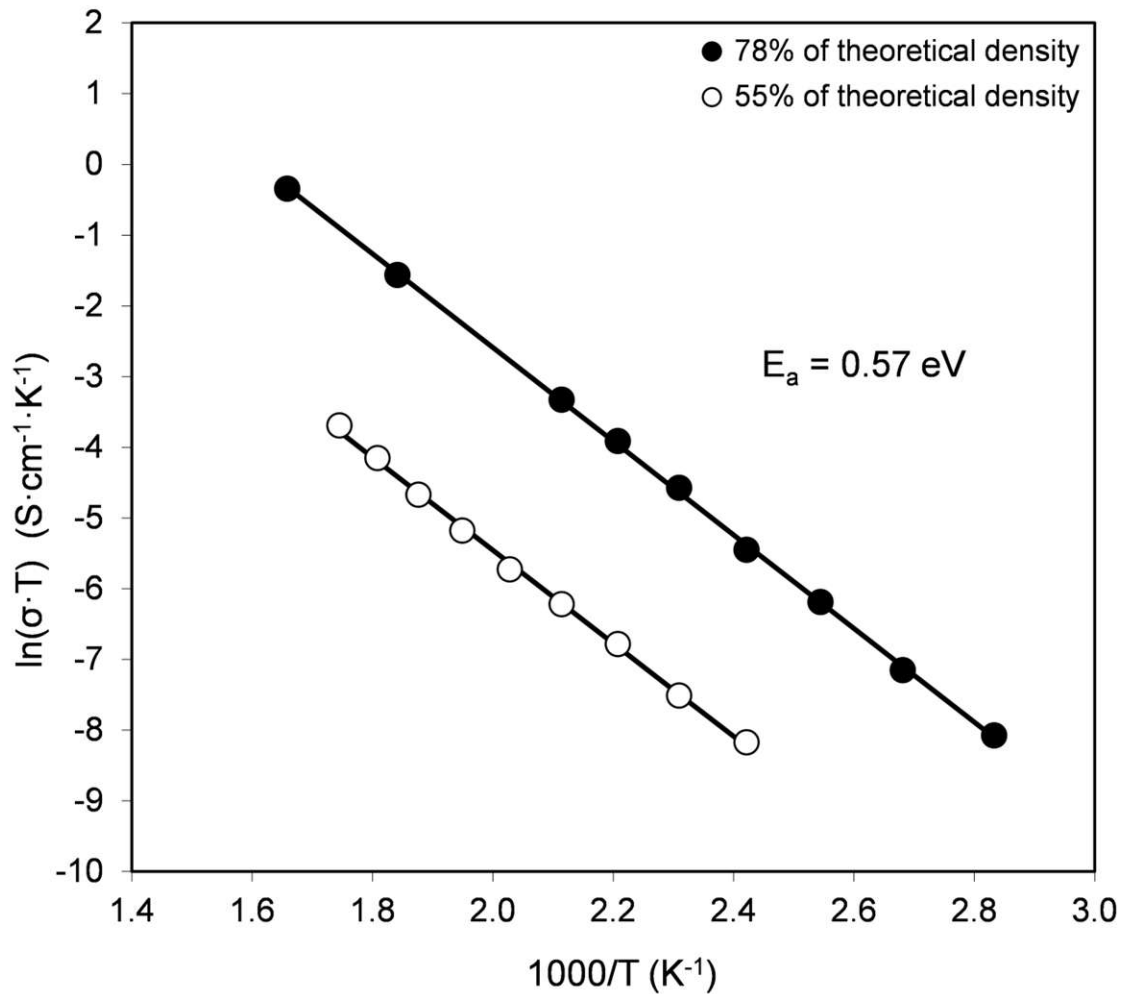
Stability of $SD\text{-Li}_2\text{PO}_2\text{N}$ in air



**Thermogravimetric analysis
curve in air**

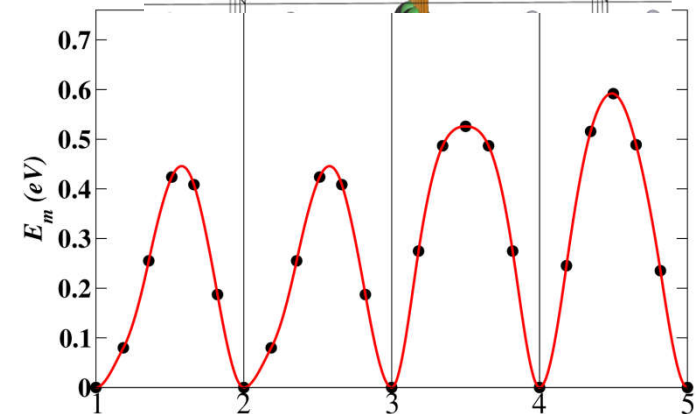
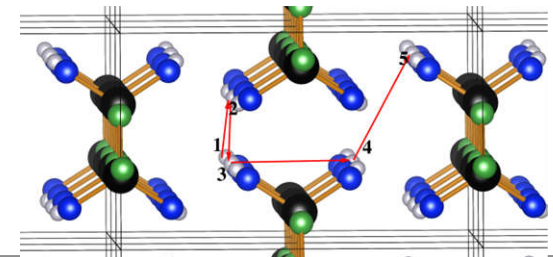
**Note: no structural changes were observed while heating in
vacuum up to 1050° C.**

Ionic conductivity of $SD\text{-Li}_2\text{PO}_2\text{N}$



$\sigma \approx 10^{-6} \text{ S/cm}$ at 80° C

NEB analysis of E_m (vacancy mechanism)



$$E_m \approx 0.4 \text{ eV}; E_f \approx 2 \text{ eV}$$

$$\Rightarrow E_A = E_m + \frac{1}{2} E_f \approx 1.4 \text{ eV}$$

➔ Sample has appreciable population of vacancies

Summary of the $\text{Li}_2\text{PO}_2\text{N}$ story

- ❑ Predicted on the basis of first principles theory
- ❑ Subsequently, experimentally realized by Keerthi Seneviranthe and colleagues; generally good agreement between experiment and theory
- ❑ Ion conductivity properties not (yet) competitive
- ❑ Crystalline *SD*- $\text{Li}_2\text{PO}_2\text{N}$ ($Cmc2_1$) is quite different from the amorphous LiPON electrolyte developed at ORNL

Other electrolyte materials -- thiophosphate

LiPON and $\text{LiS}_2\text{-P}_2\text{S}_5$ conductivities

X. Yu, J. B. Bates, G. E. Jellison, Jr., and F. X. Hart, *J. Electrochem. Soc.* **144** 524-532 (1997):

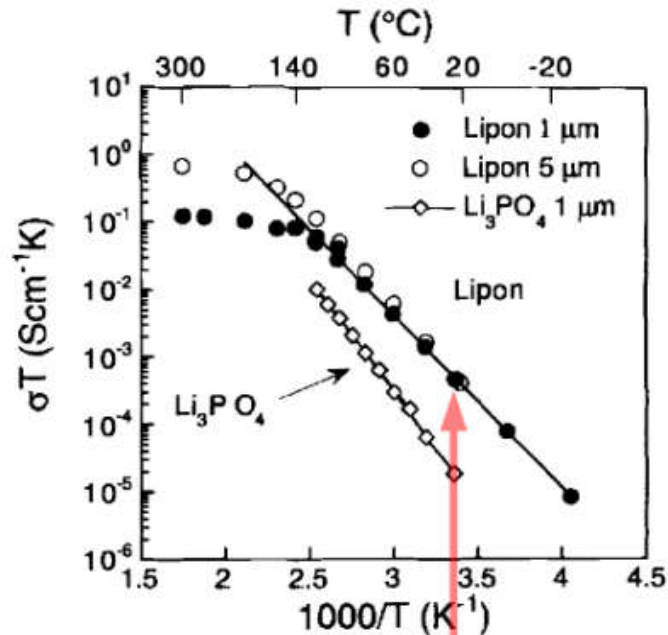


Fig. 3. Arrhenius plot of ionic conductivity of Lipon and Li_3PO_4 vs. temperature.

$$\sigma = 2 \times 10^{-6} \text{ S/cm}$$

$$E_a = 0.5 \text{ eV}$$

M. Tatsumisago and A. Hayashi, *J. Non-Cryst. Solids* **354** 1411-1417 (2008):

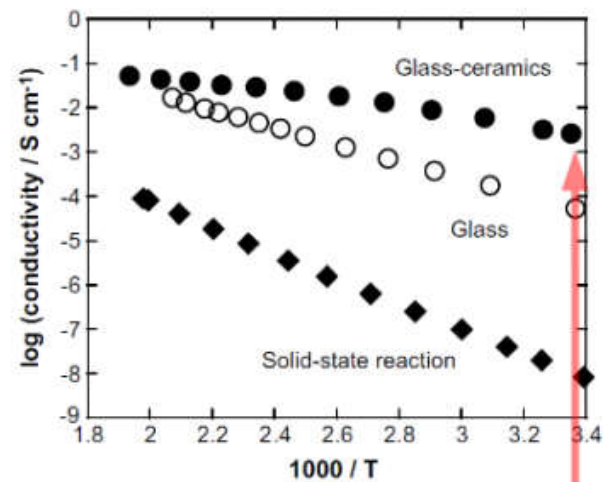
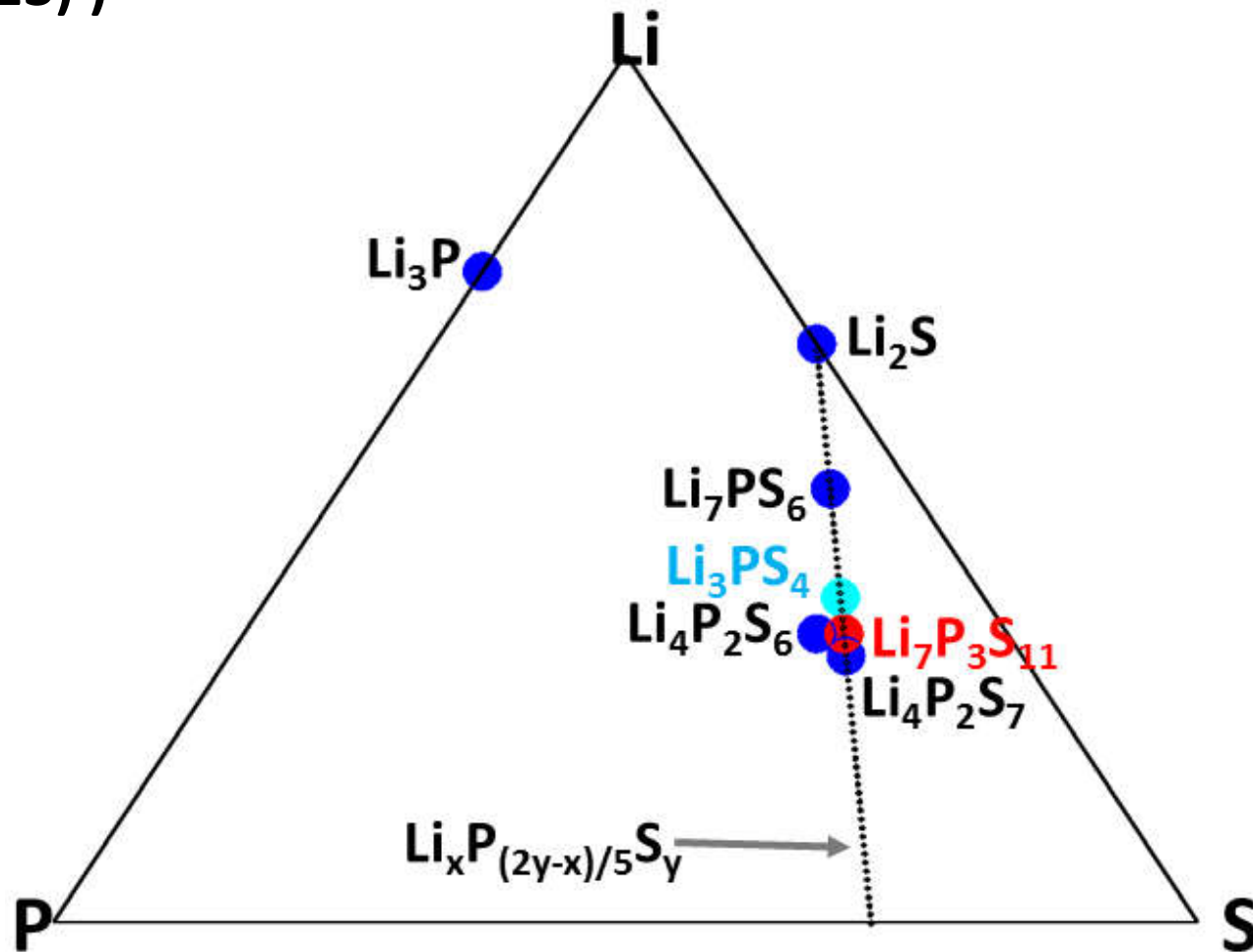


Fig. 5. Temperature dependences of the conductivities for the $70\text{Li}_2\text{S} \cdot 30\text{P}_2\text{S}_5$ glass and glass-ceramics. The conductivity data for the sample prepared by solid-state reaction are also shown.

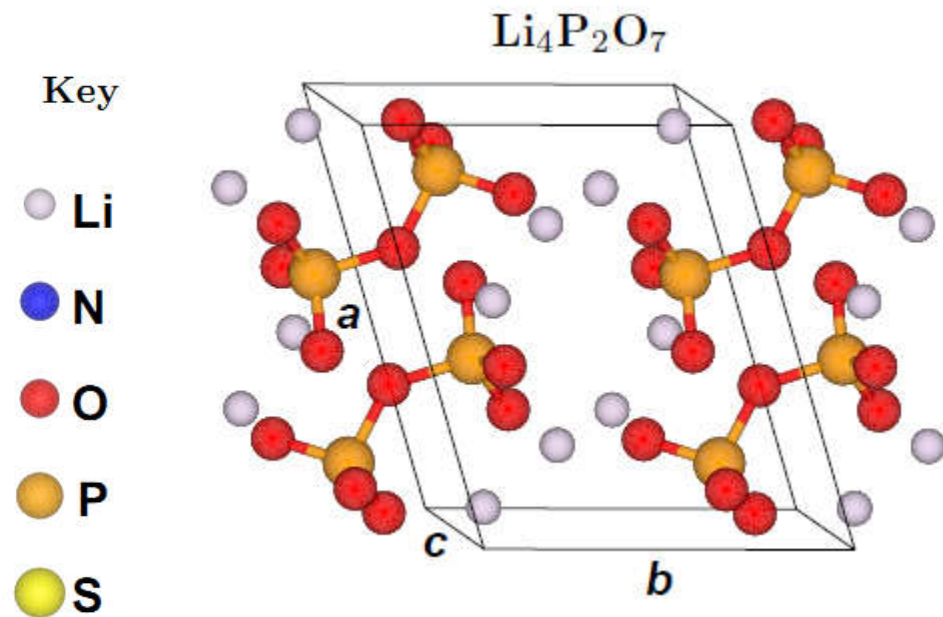
$$\sigma = 3 \times 10^{-3} \text{ S/cm}$$

$$E_a = 0.1 \text{ eV}$$

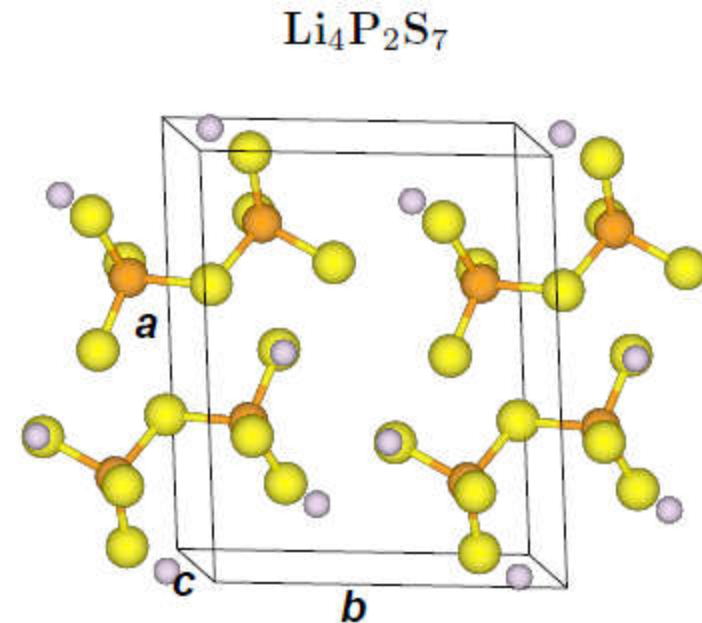
Systematic study of Li_xPS_y materials – (N. D. Lepley and N. A. W. Holzwarth, J. Electrochem. Soc. 159, A538 (2012), Phys. Rev. B 88, 104103 (2013))



Comparison of some lithium phosphates and thiophosphates

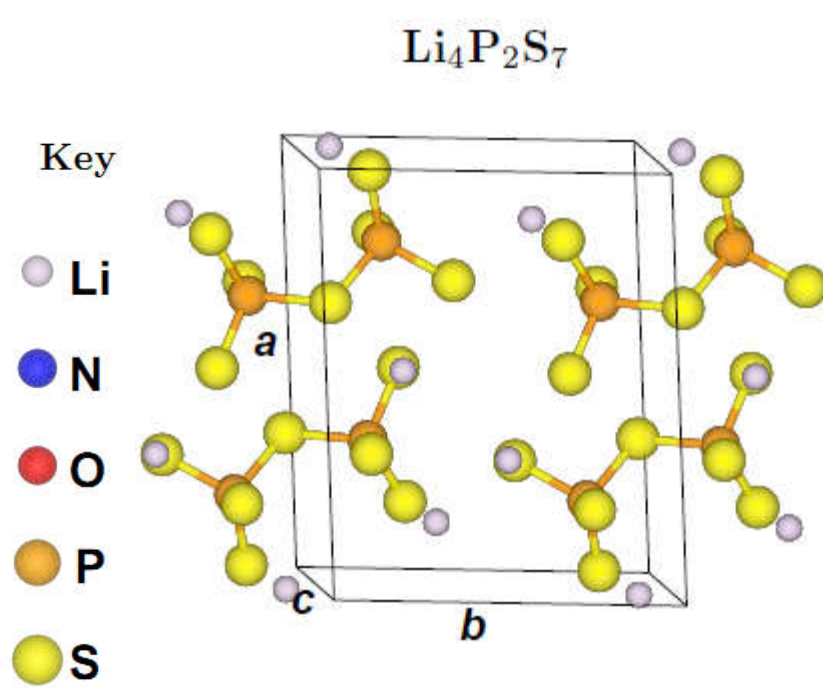


Crystallizes (experimentally and computationally) into $P\bar{1}$ structure

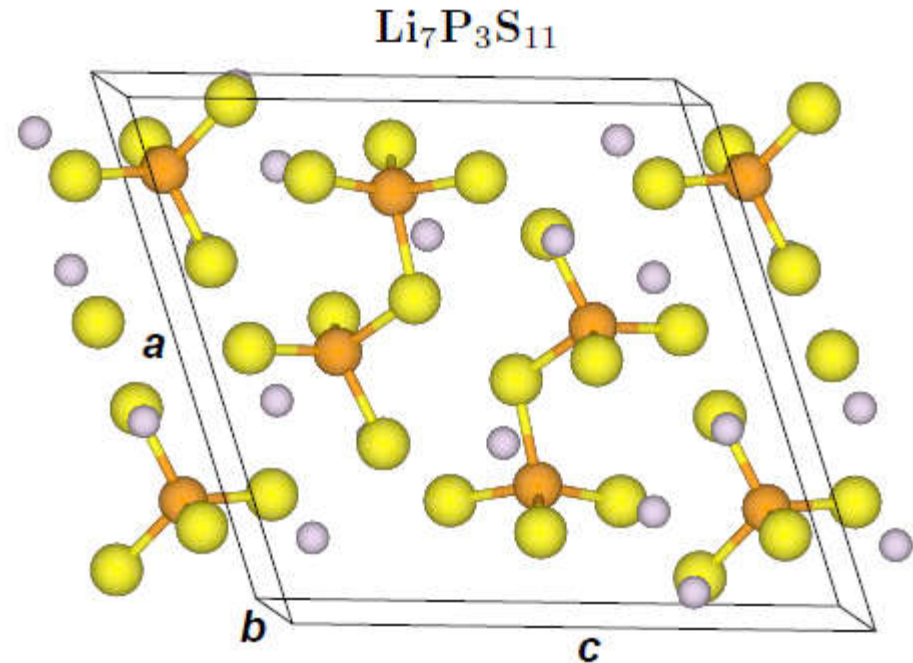


Experimentally amorphous; computationally metastable in $P\bar{1}$ structure

Some lithium thiophosphate crystal structures



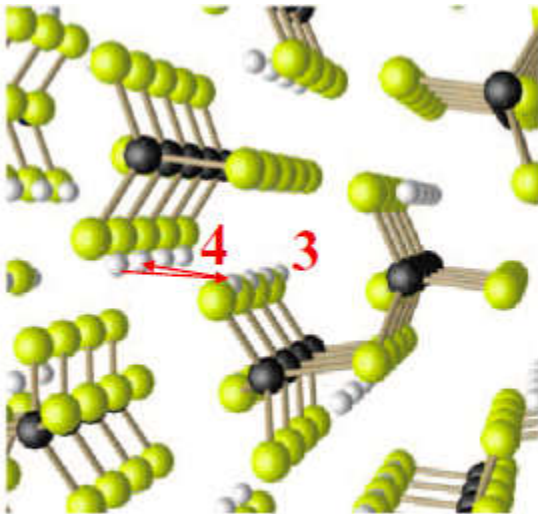
Experimentally amorphous;
computationally metastable
in $P\bar{1}$ structure



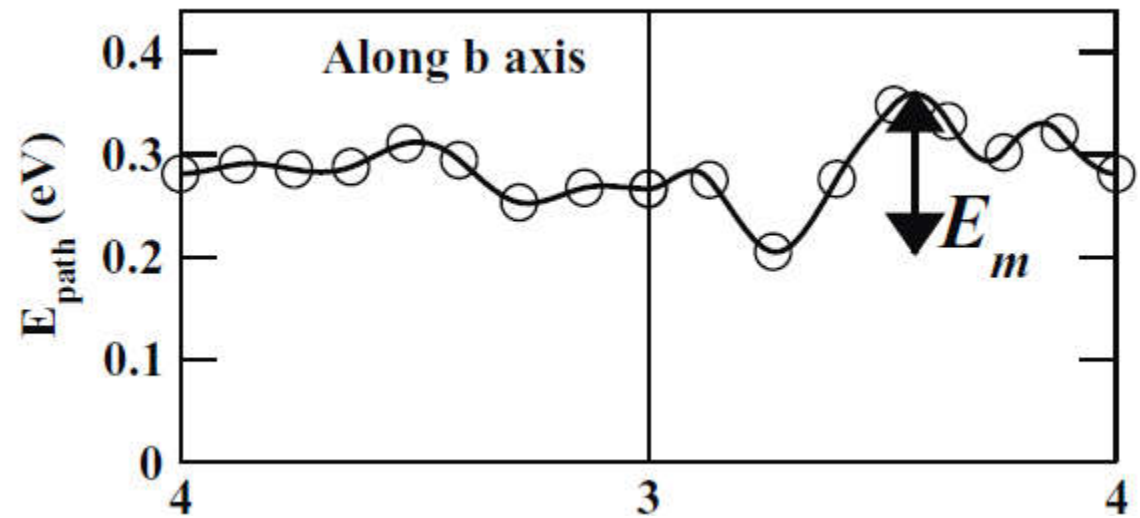
Experimentally and computationally
metastable in $P\bar{1}$ structure

Vacancy migration analysis from NEB results

for $\text{Li}_7\text{P}_3\text{S}_{11}$: Lepley & Holzwarth, *JECs* **159**, A538-A547 (2012)



● Li ● P ● S



$$E_m \approx 0.15 \text{ eV}; E_f \approx 0 \text{ eV}$$

$$\Rightarrow E_A = E_m + \frac{1}{2}E_f \approx 0.15 \text{ eV}$$

Experiment -- A Hayashi *et al.*, *J. Solid State Electrochem.* **14**, 1761 (2010):

$$\sigma \approx 2 - 3 \times 10^{-3} \text{ S/cm}$$

$$E_A \approx 0.12 - 0.18 \text{ eV}$$

Anomalous High Ionic Conductivity of Nanoporous $\beta\text{-Li}_3\text{PS}_4$

Zengcai Liu,[†] Wujun Fu,[†] E. Andrew Payzant,^{†,‡} Xiang Yu,[†] Zili Wu,^{†,§} Nancy J. Dudney,[‡] Jim Kiggans,[‡] Kunlun Hong,[†] Adam J. Rondinone,[†] and Chengdu Liang^{*,†}

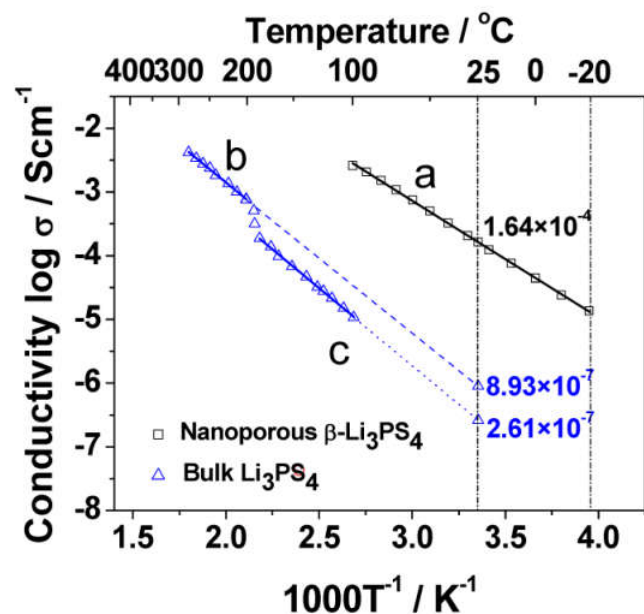


Figure 1. Arrhenius plots for nanoporous $\beta\text{-Li}_3\text{PS}_4$ (line a), bulk $\beta\text{-Li}_3\text{PS}_4$ (line b), and bulk $\gamma\text{-Li}_3\text{PS}_4$ (line c). The conductivity data for bulk Li_3PS_4 are reproduced from the work of Tachez.¹⁰

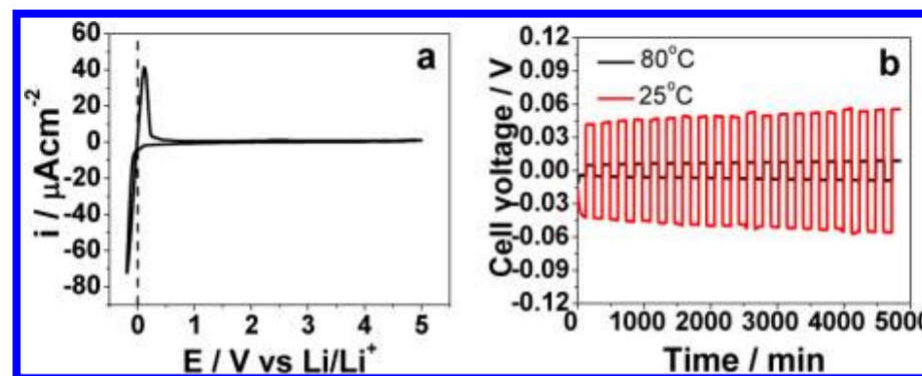
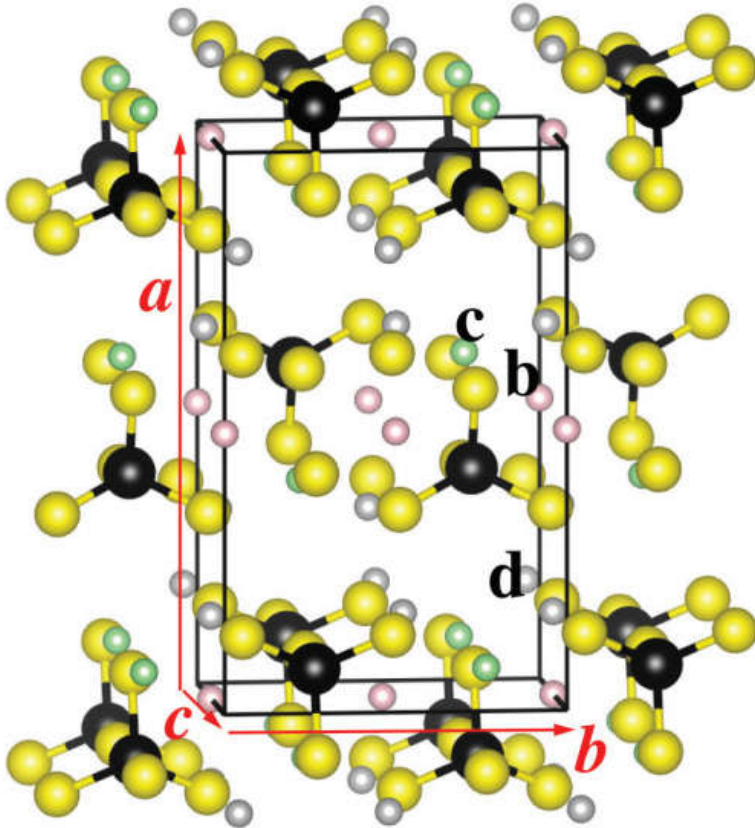


Figure 5. Electrochemical stability of $\beta\text{-Li}_3\text{PS}_4$ and cycling stability with metallic lithium electrodes. (a) CV of a $\text{Li}/\beta\text{-Li}_3\text{PS}_4/\text{Pt}$ cell, where Li and Pt serve as the reference/counter and working electrodes, respectively. (b) Lithium cyclability in a symmetric $\text{Li}/\beta\text{-Li}_3\text{PS}_4/\text{Li}$ cell. The cell was cycled at a current density of 0.1 mA cm^{-2} at room temperature and $80 \text{ }^\circ\text{C}$.

β -Li₃PS₄

● Li ● P ● S

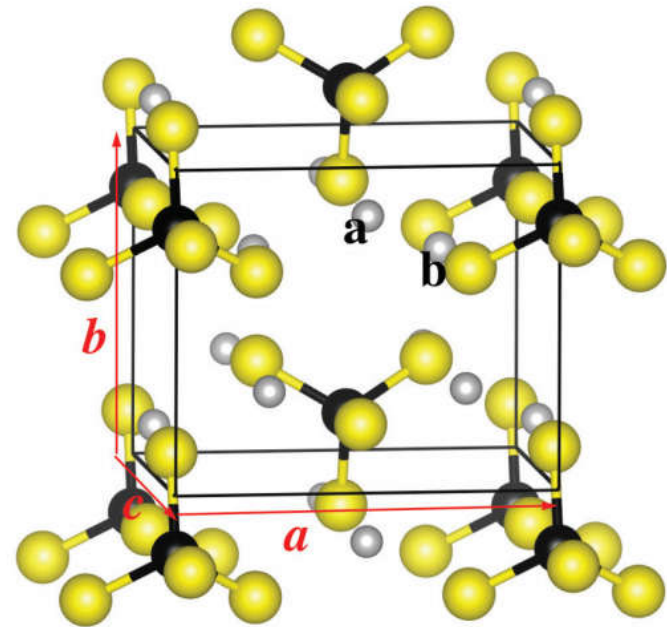


$$E_m \approx 0.3 \text{ eV}; E_f \approx 0.0 \text{ eV}$$

$$E_A = E_m + \frac{1}{2}E_f \approx 0.3 \text{ eV}$$

$$E_{exp} = 0.4 - 0.5 \text{ eV}$$

γ -Li₃PS₄



$$E_m \approx 0.3 \text{ eV}; E_f \approx 0.8 \text{ eV}$$

$$E_A = E_m + \frac{1}{2}E_f \approx 0.7 \text{ eV}$$

$$E_{exp} = 0.5 \text{ eV}$$

Lepley, Du, and Holzwarth, *PRB* **88**, 104103 (2013)

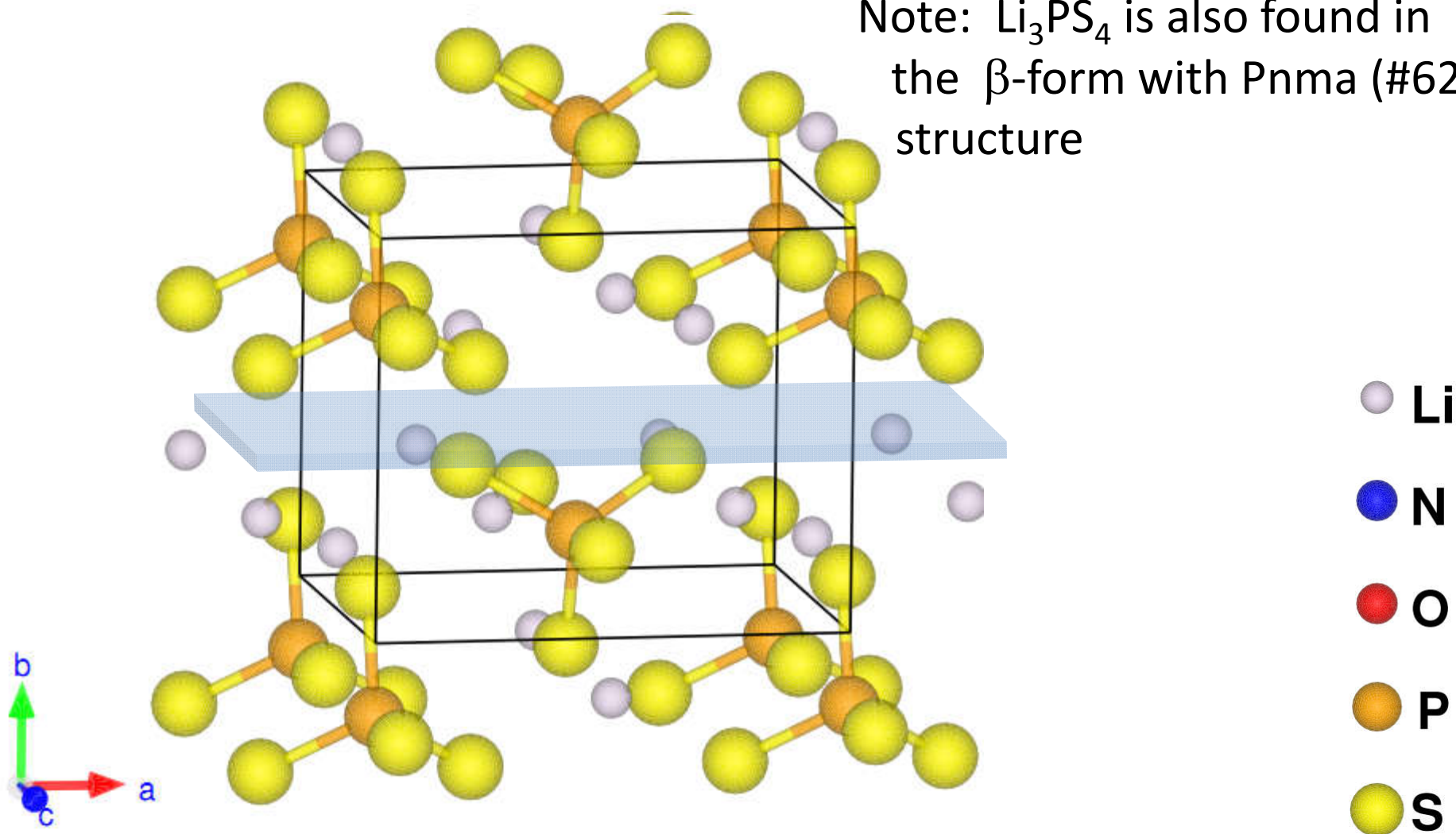
Summary of the Li_xPS_y story

- ❑ Simulations verify that thiophosphates have better ion mobility properties than their phosphate analogs
- ❑ Meta-stable crystalline $\text{Li}_7\text{P}_3\text{S}_{11}$ has been shown to have particularly favorable ion migration pathways
- ❑ γ - and β - Li_3PS_4 have very similar structures, but simulations show their ion mobilities to be different.

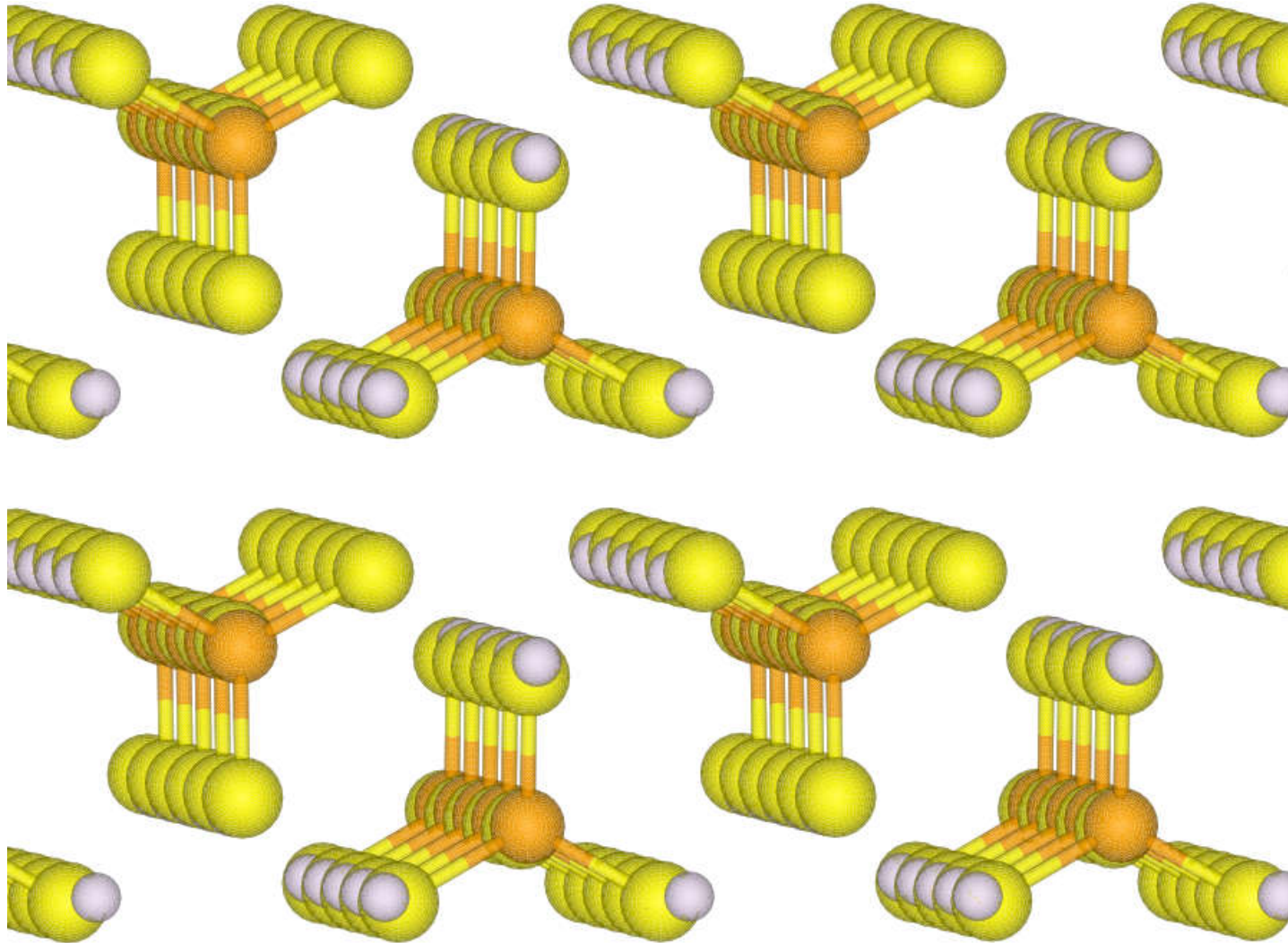
Models of Idealized Interfaces

Crystal structure of bulk Li_3PS_4 – γ -form $\text{Pmn}2_1$ (#31)

Note: Li_3PS_4 is also found in the β -form with Pnma (#62) structure

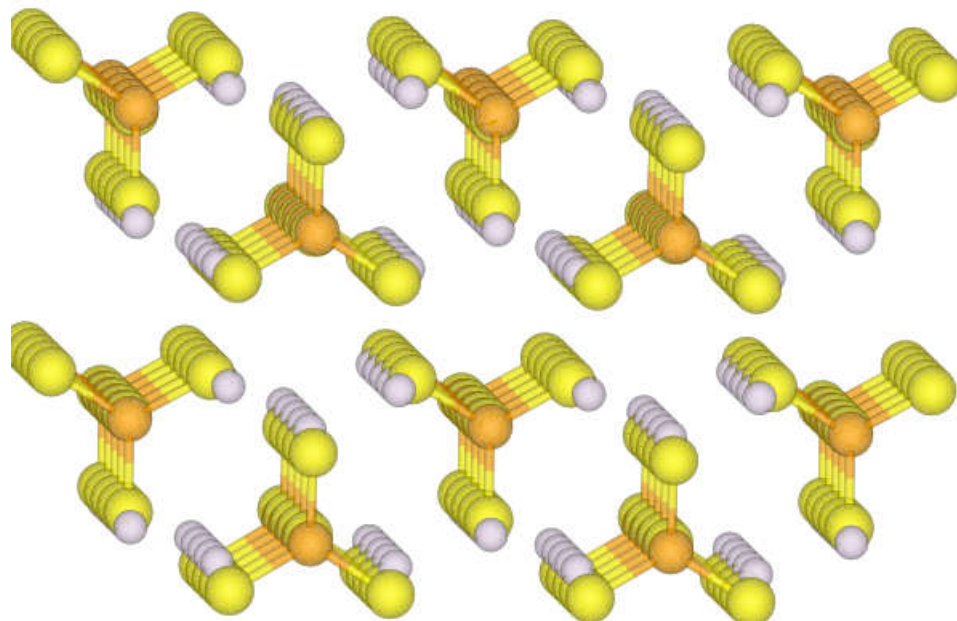


γ -Li₃PS₄ [0 1 0] surface

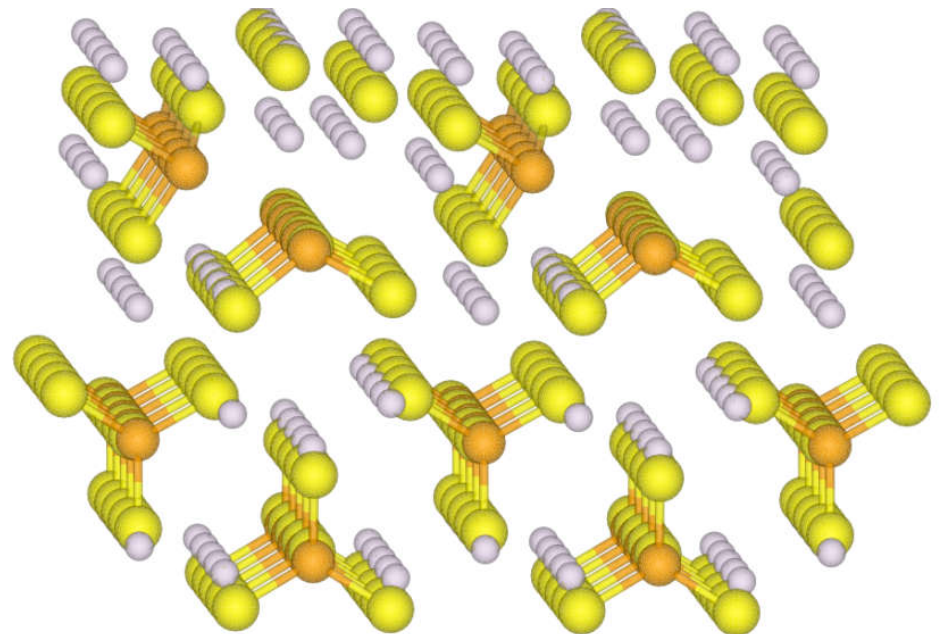


Simulations of ideal γ -Li₃PS₄ [0 1 0] surface in the presence of Li

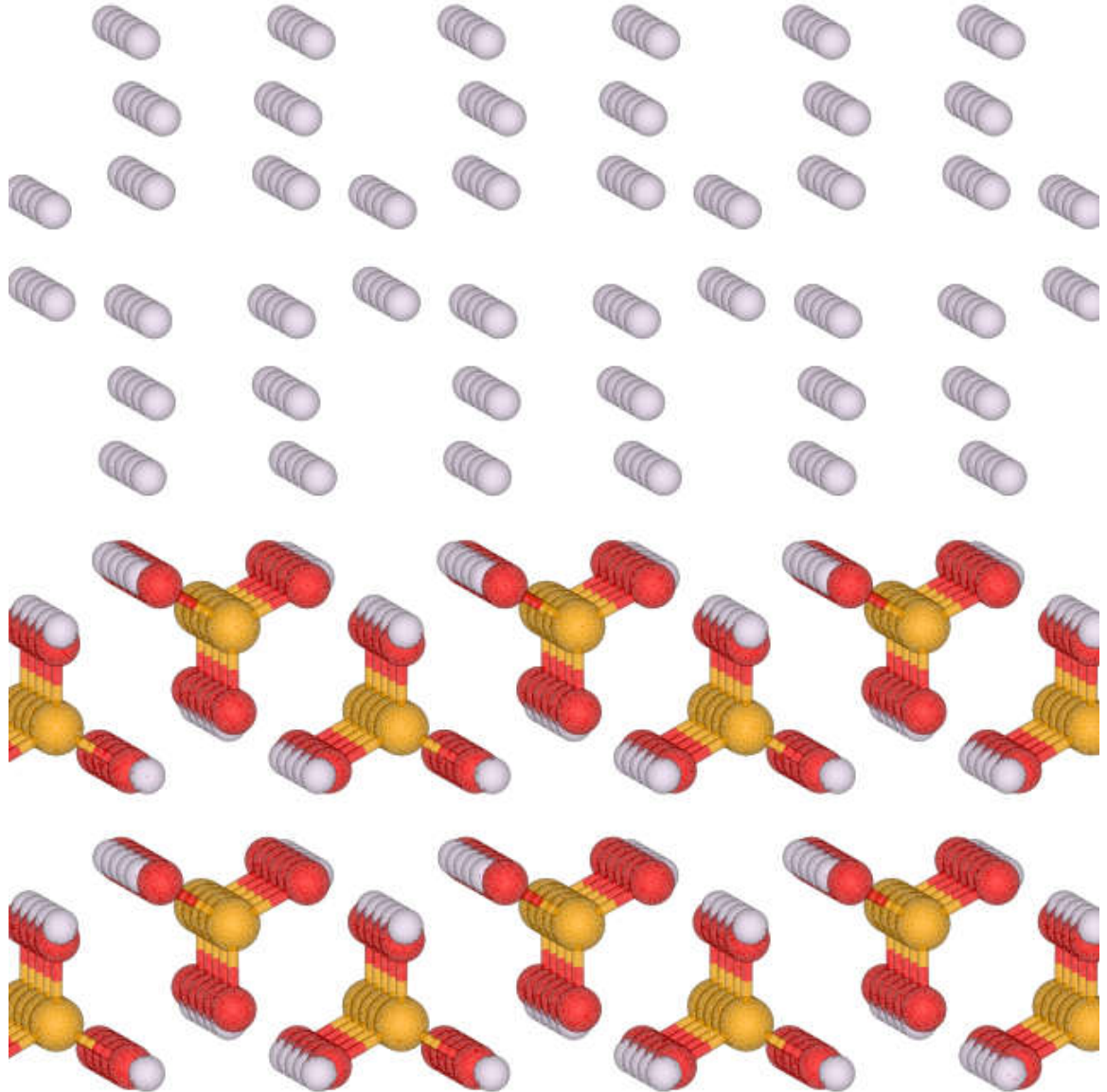
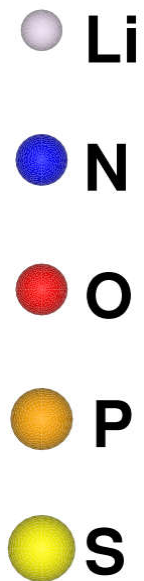
Initial configuration:



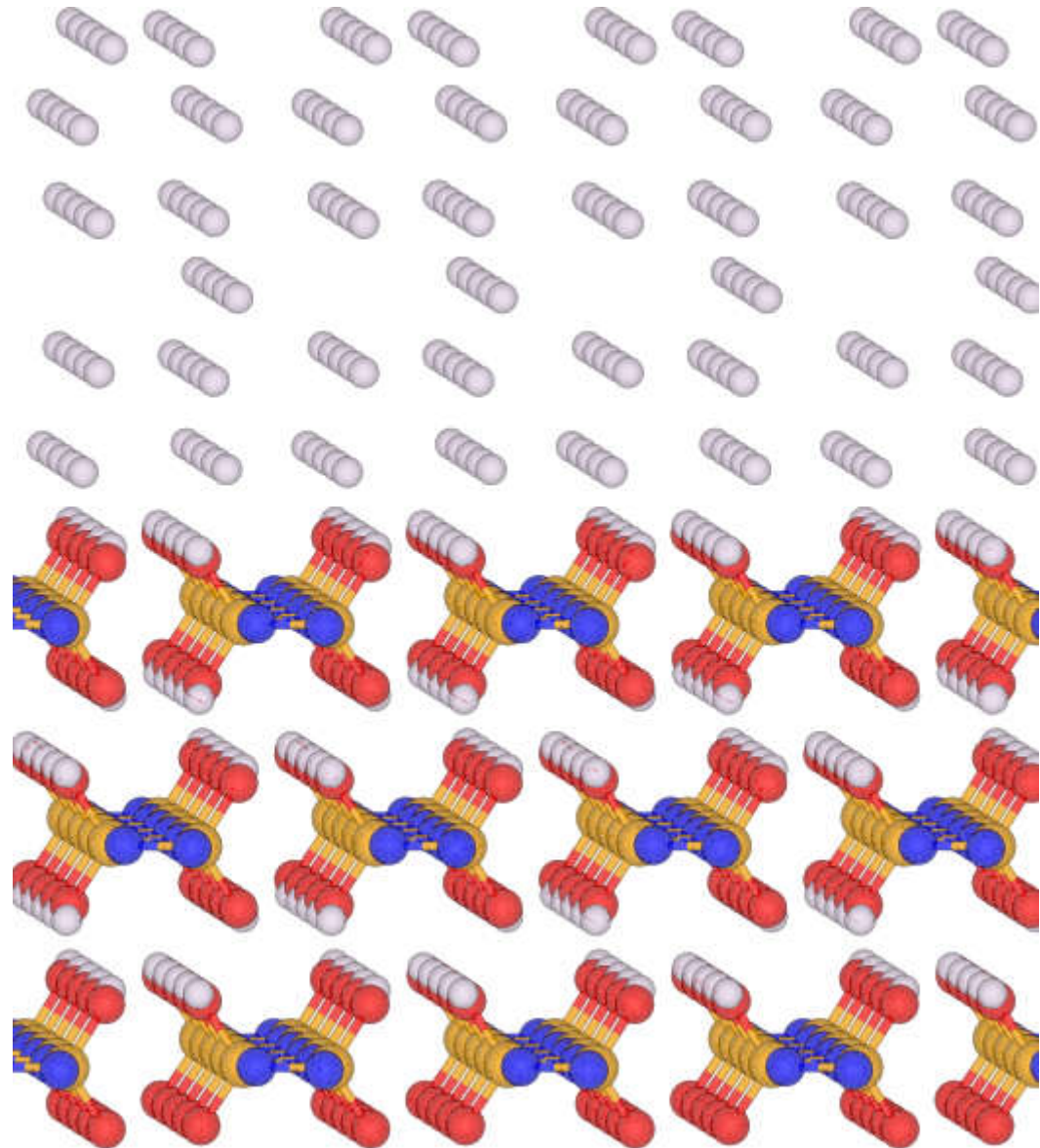
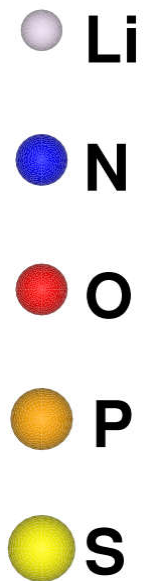
Computed optimized
structure:



Computational counter example – stable interface: $\text{Li}/\beta\text{-Li}_3\text{PO}_4$



Computational counter example – stable interface: $\text{Li}/\text{SD-Li}_2\text{PO}_2\text{N}$



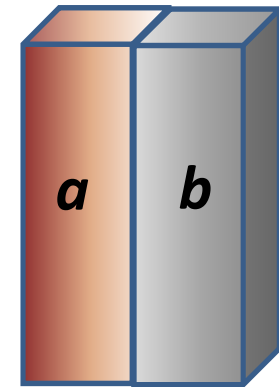
Quantitative study of interfaces –
 (Lepley & Holzwarth, *PRB* **92** 214201 (2015))

Within any given periodic simulation cell with n_a units of material a and with n_b units of material b , we can define an interface energy:

$$\tilde{\gamma}_{ab}(\tilde{\Omega}, n_a, n_b) = \frac{\tilde{E}_{ab}(\tilde{\Omega}, A, n_a, n_b) - n_a E_a - n_b E_b}{A}$$

area of interface
within supercell

bulk energies

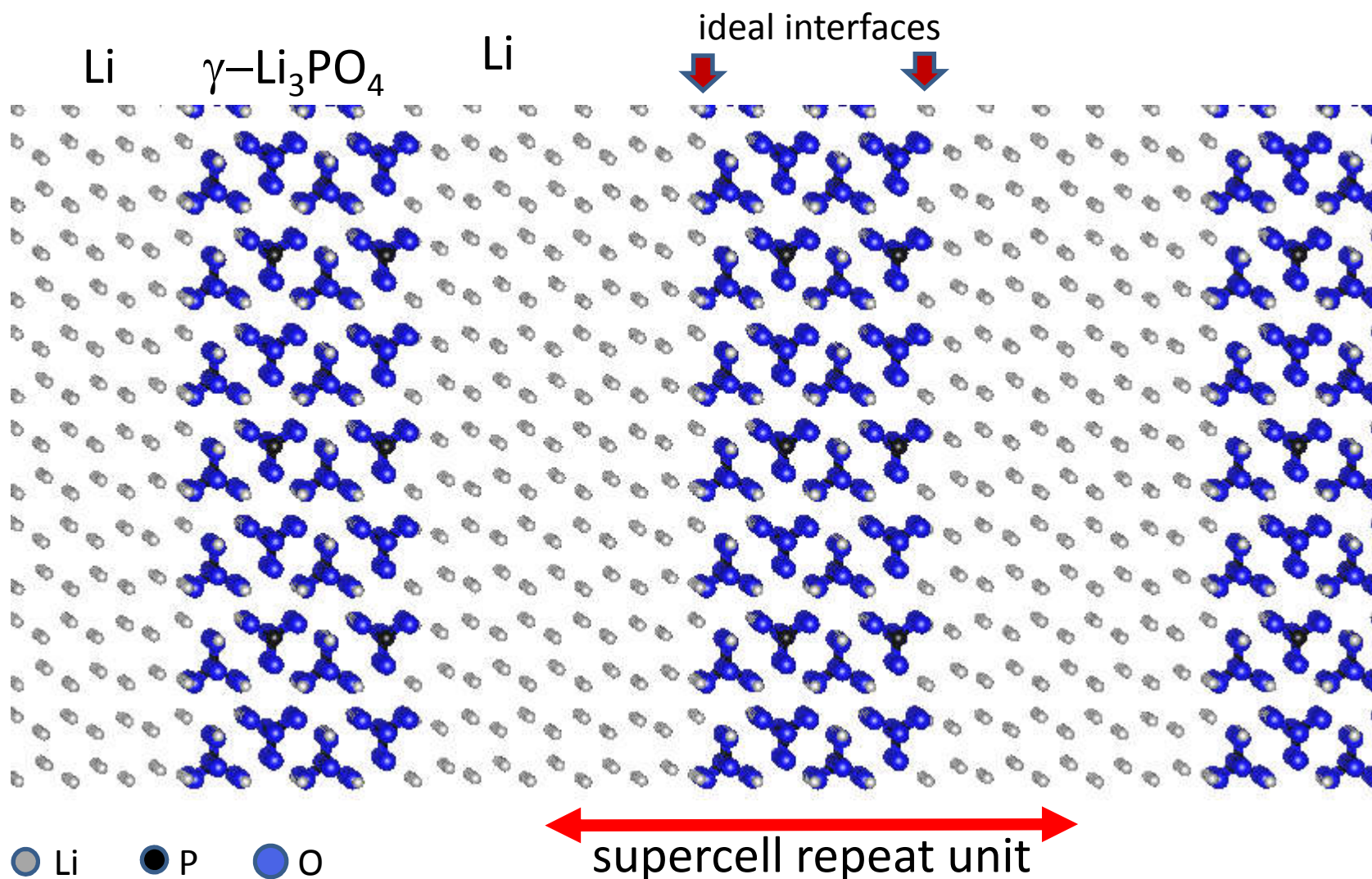


In order approximately remove the effects of lattice strain:

- Design the supercell to be commensurate with lattice a
- Now the strain will scale with the amount of material b

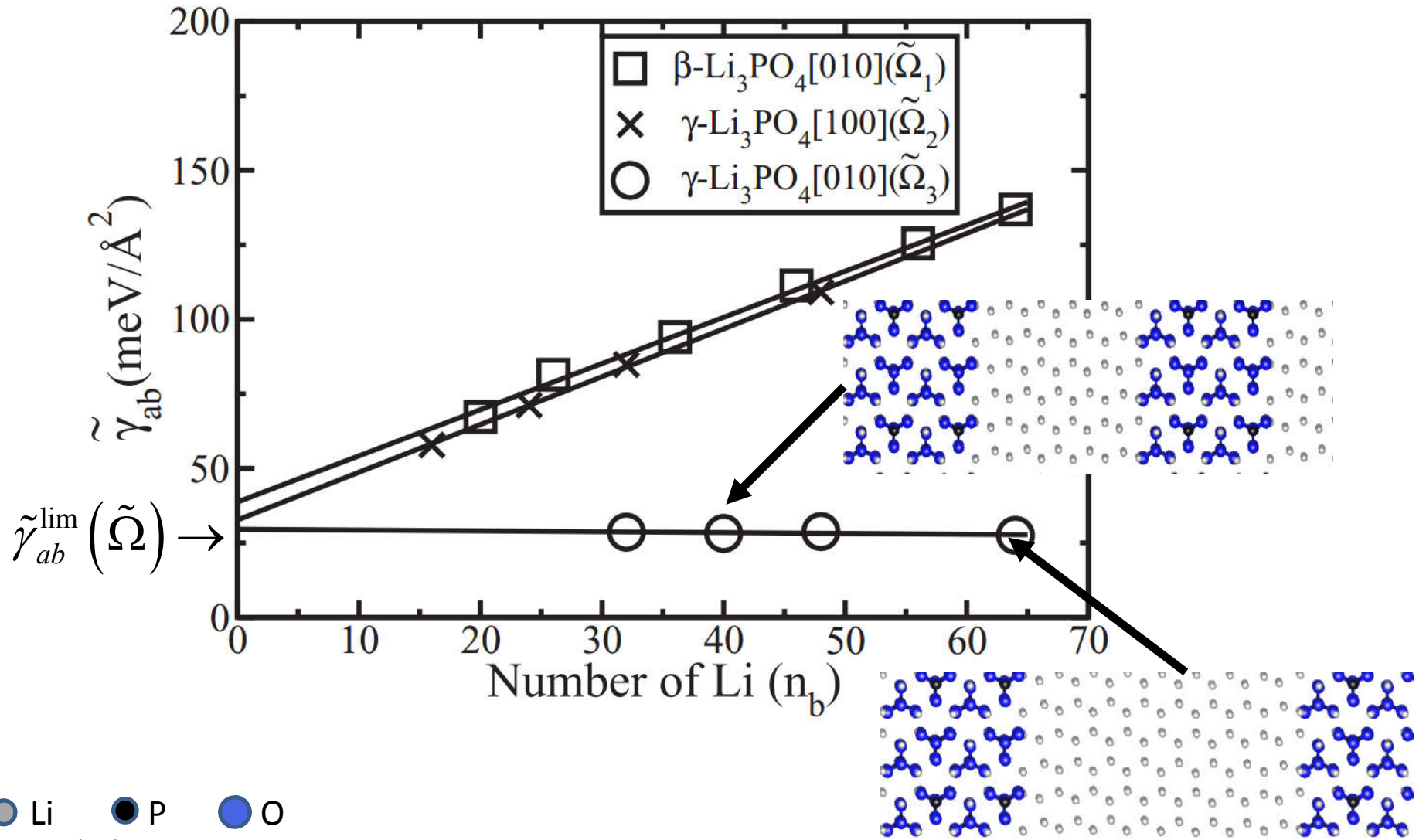
$$\Rightarrow \tilde{\gamma}_{ab}(\tilde{\Omega}, n_a, n_b) = \tilde{\gamma}_{ab}^{\text{lim}}(\tilde{\Omega}) + n_b \sigma$$

It is convenient to model the interface between a solid electrolyte and solid electrode in the slab geometry using a periodic simulation cell:



Lepley's linear equation for the interface

energy: $\tilde{\gamma}_{ab}(\tilde{\Omega}, n_a, n_b) = \tilde{\gamma}_{ab}^{\text{lim}}(\tilde{\Omega}) + n_b \sigma$



● Li ● P ● O

8/17/2016

IMRC 2016

55

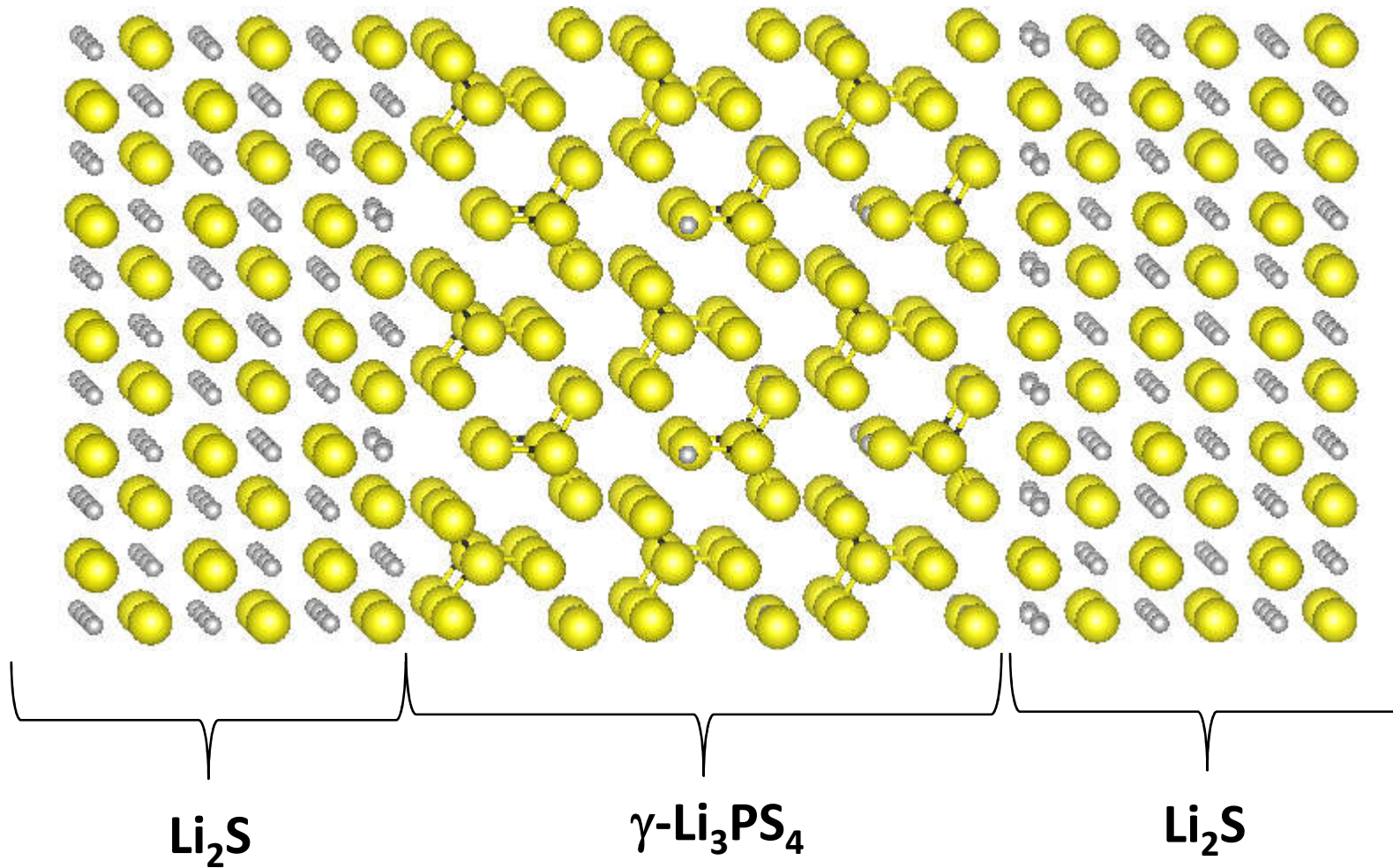
Some interface energy results

System	$\tilde{\gamma}_{ab}^{\text{lim}}$ (meV/Å ²)	σ (meV/Å ²)
Li ₂ O[110]/Li(Ω ₁)	30	6.1
Li ₂ O[110]/Li(Ω ₂)	26	0.2
Li ₂ S[110]/Li(Ω ₃)	19	0.2
Li ₂ S[100]/Li(Ω ₄)	19	0.0
γ-Li ₃ PO ₄ [010]/Li(Ω ₃)	31	0.0
γ-Li ₃ PS ₄ [010]/Li ₂ S [110]	16	1.0
γ-Li ₃ PS ₄ [010]/Li	-216	-0.1



$\gamma\text{-Li}_3\text{PS}_4$ [010]/ Li_2S [110]

● Li ● P ● S



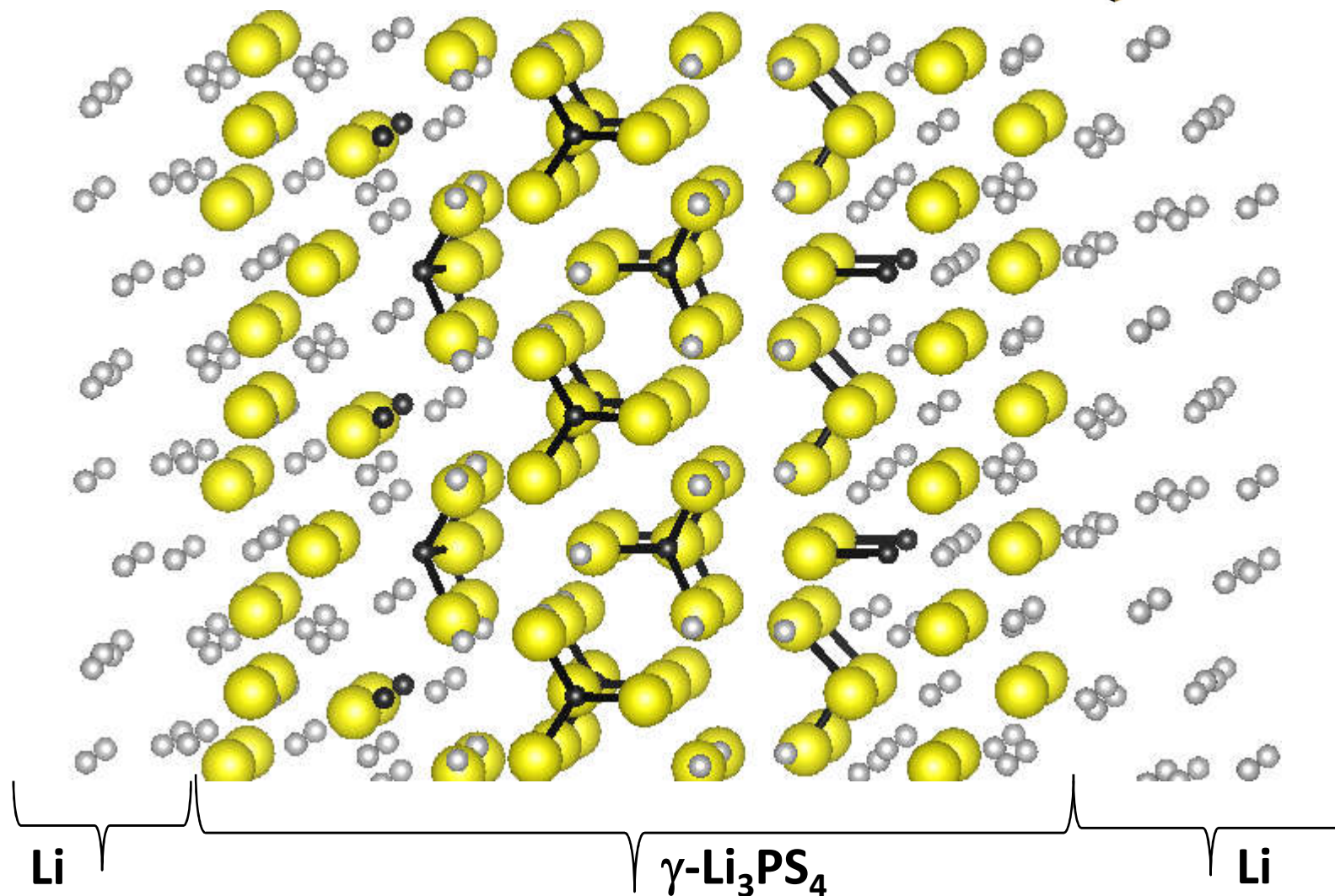
Stable interface; composite electrolyte system

$\gamma\text{-Li}_3\text{PS}_4$ [010]/Li

● Li ● P ● S

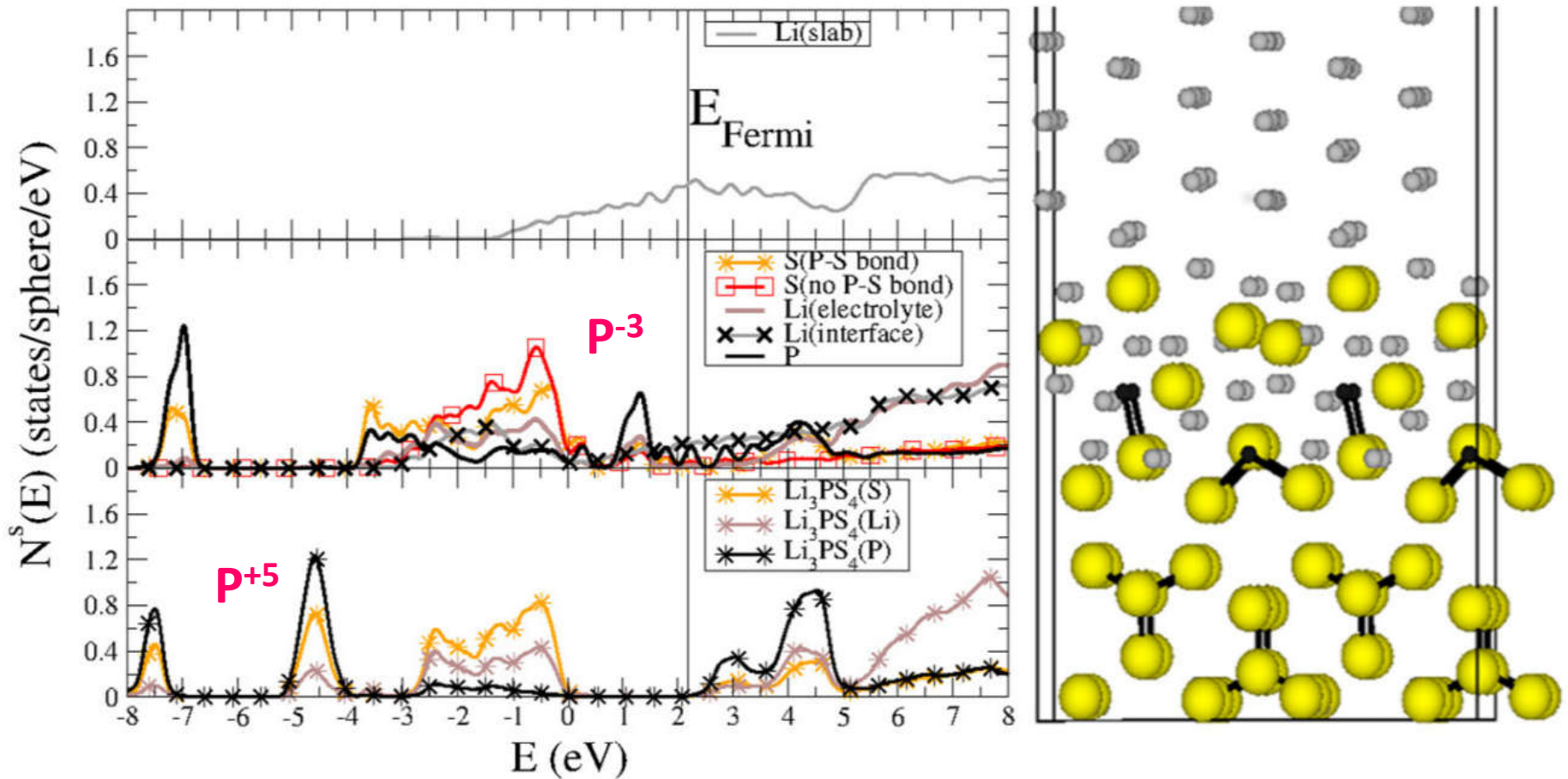


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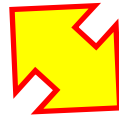


Initially unstable interface; (meta)-stable buffer layer formed

Partial density of states analysis of unstable $\text{Li}_3\text{PS}_4/\text{Li}$ interface:



Bulk reactions from estimated heats of formation

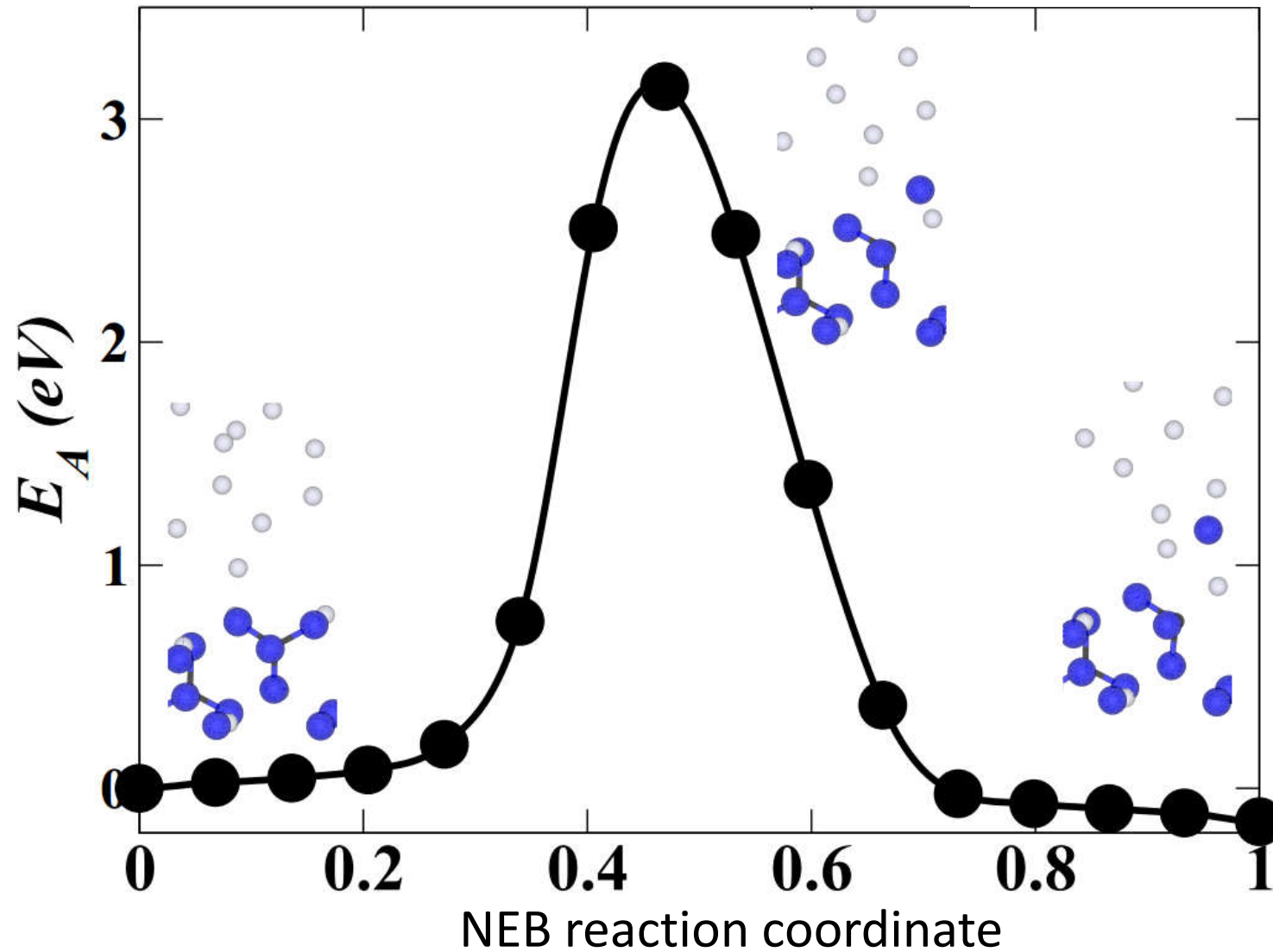


Decomposition at interface



(Meta-)stable interface

Evidence of kinetic barrier at $\text{Li}_3\text{PO}_4/\text{Li}$ interface



Summary of ideal interface story

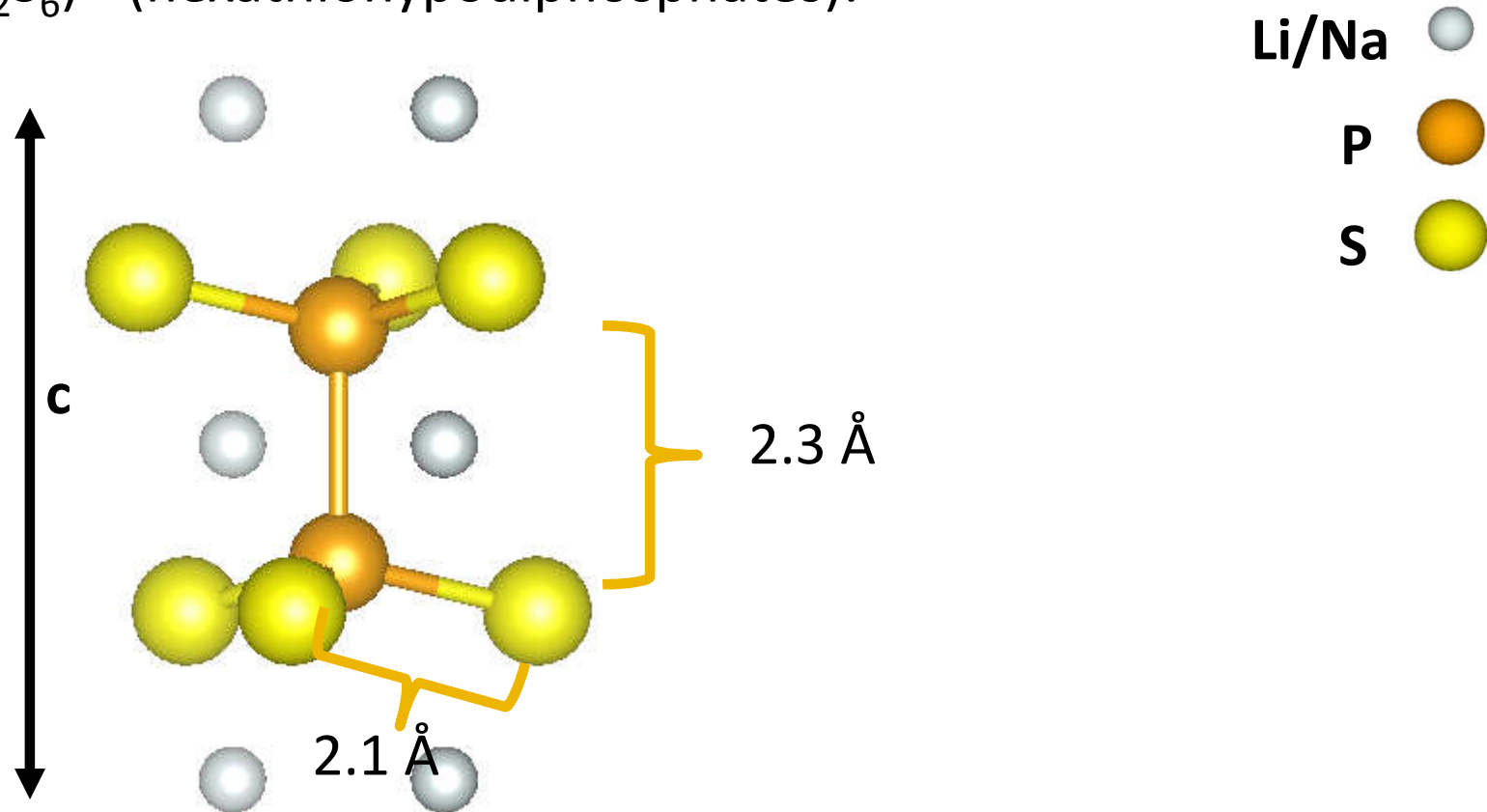
- ❑ A practical scheme was developed to compute an intensive measure of the interface interaction $\tilde{\gamma}_{ab}^{\text{int}}$, explicitly accounting for the effects of lattice strain.
- ❑ Discussed bulk reactivity as related to the interface stability of the interfaces of
 - ❑ $\text{Li}_3\text{PO}_4/\text{Li}$ (having a significant kinetic barrier to decomposition)
 - ❑ $\text{Li}_3\text{PS}_4/\text{Li}$ (having localized decomposition).

Other thiophosphate compounds; comparison of $\text{Li}_4\text{P}_2\text{S}_6$ and $\text{Na}_4\text{P}_2\text{S}_6$

Zachary D. Hood, Cameron Kates, Melanie Kirkham, Shiba Adhikari, Chengdu Liang, and
N. A. W. Holzwarth, *Solid State Ionics* **284**, 61-70 (2015)

Larry E. Rush Jr. and N.A.W. Holzwarth, *Solid State Ionics* **286**, 45-50 (2016)

High temperature processing of Li and Na thiophosphates have shown to produce dimer units $(P_2S_6)^{-4}$ (hexathiohypodiphosphates):



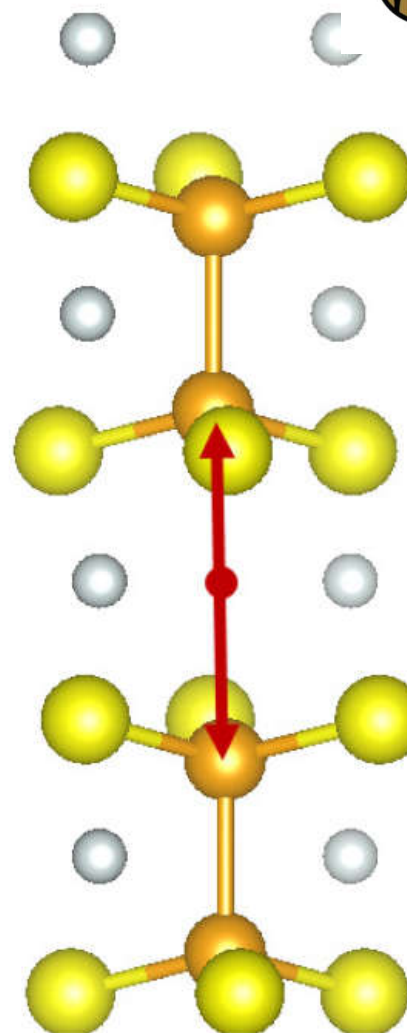
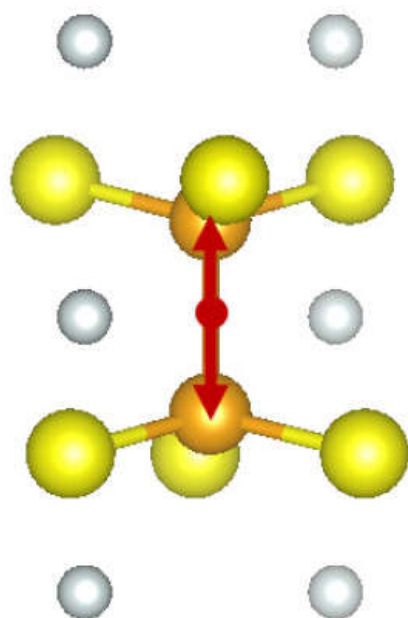
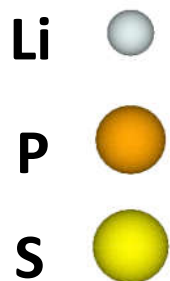
R. Mercier, J.P. Malugani, B. Fahys, J. Douglade, G. Robert, *J. Solid State Chem.* **43**, 151–162 (1982).

A. Kuhn, R. Eger, J. Nuss, B.V. Lotsch, *Z. Anorg. Allg. Chem.* **640**, 689–692 (2014).

Mercier's analysis of $\text{Li}_4\text{P}_2\text{S}_6$

-- disorder in P-P placements

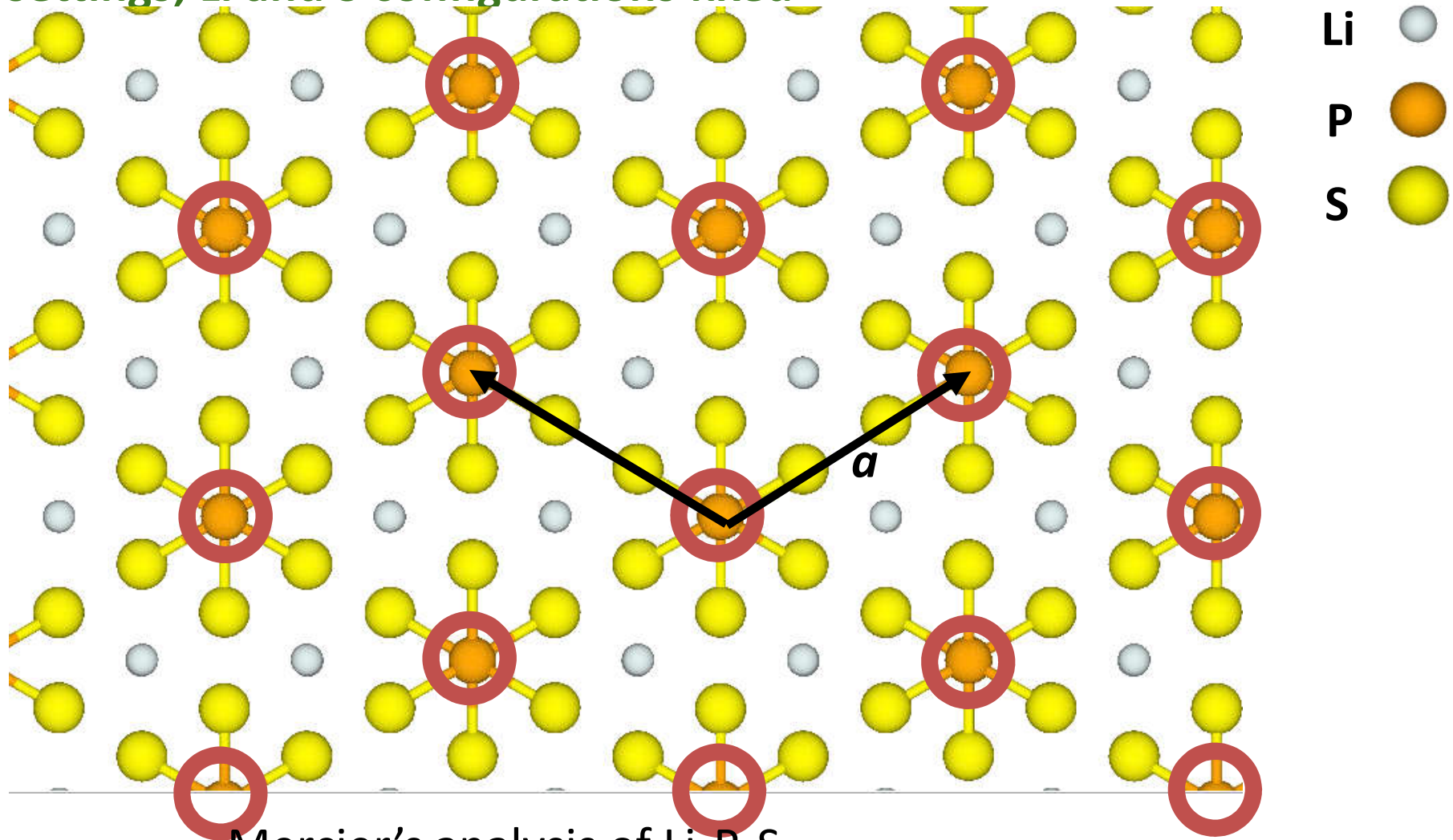
Crystal Space Group $P6_3/mcm$ (#193)



$$P_{\uparrow} \equiv \pm z_P c$$

$$P_{\downarrow} \equiv \pm \left(\frac{1}{2} - z_P \right) c$$

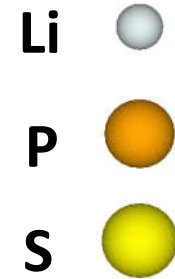
Structural variation can be mapped on to a two-dimensional hexagonal lattice with each P configuration taking $P\uparrow$ or $P\downarrow$ settings; Li and S configurations fixed



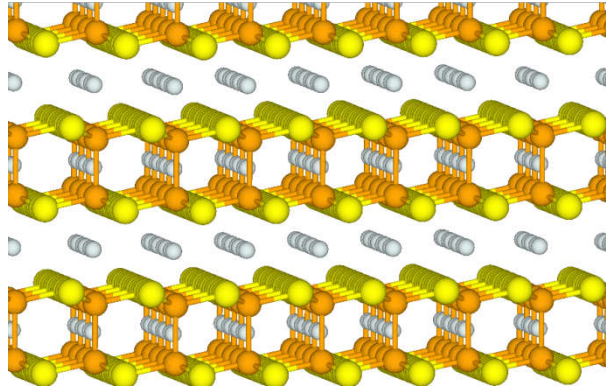
Mercier's analysis of $\text{Li}_4\text{P}_2\text{S}_6$

Examples:

Mercier's analysis of $\text{Li}_4\text{P}_2\text{S}_6$



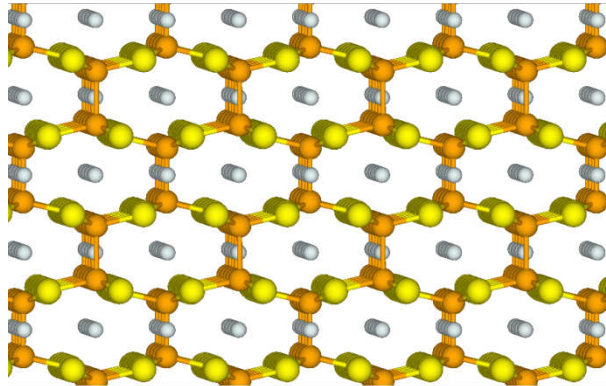
Structure "b"
 $P\bar{3}1m$



$\Delta E = 0.03 \text{ eV}$

100% P↑

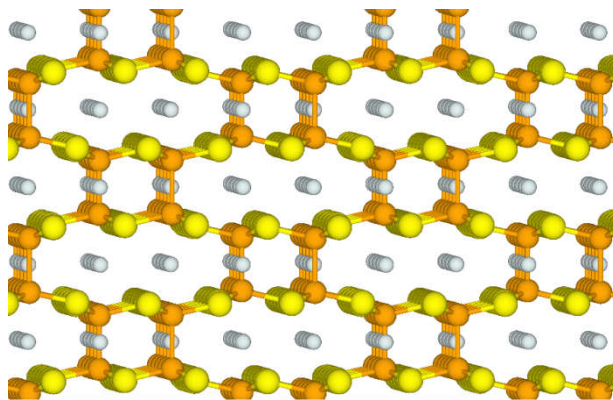
Structure "c"
 $Pn\bar{1}m$



$\Delta E = 0$

50% P↑
50% P↓

Structure "d"
 $Pnma$



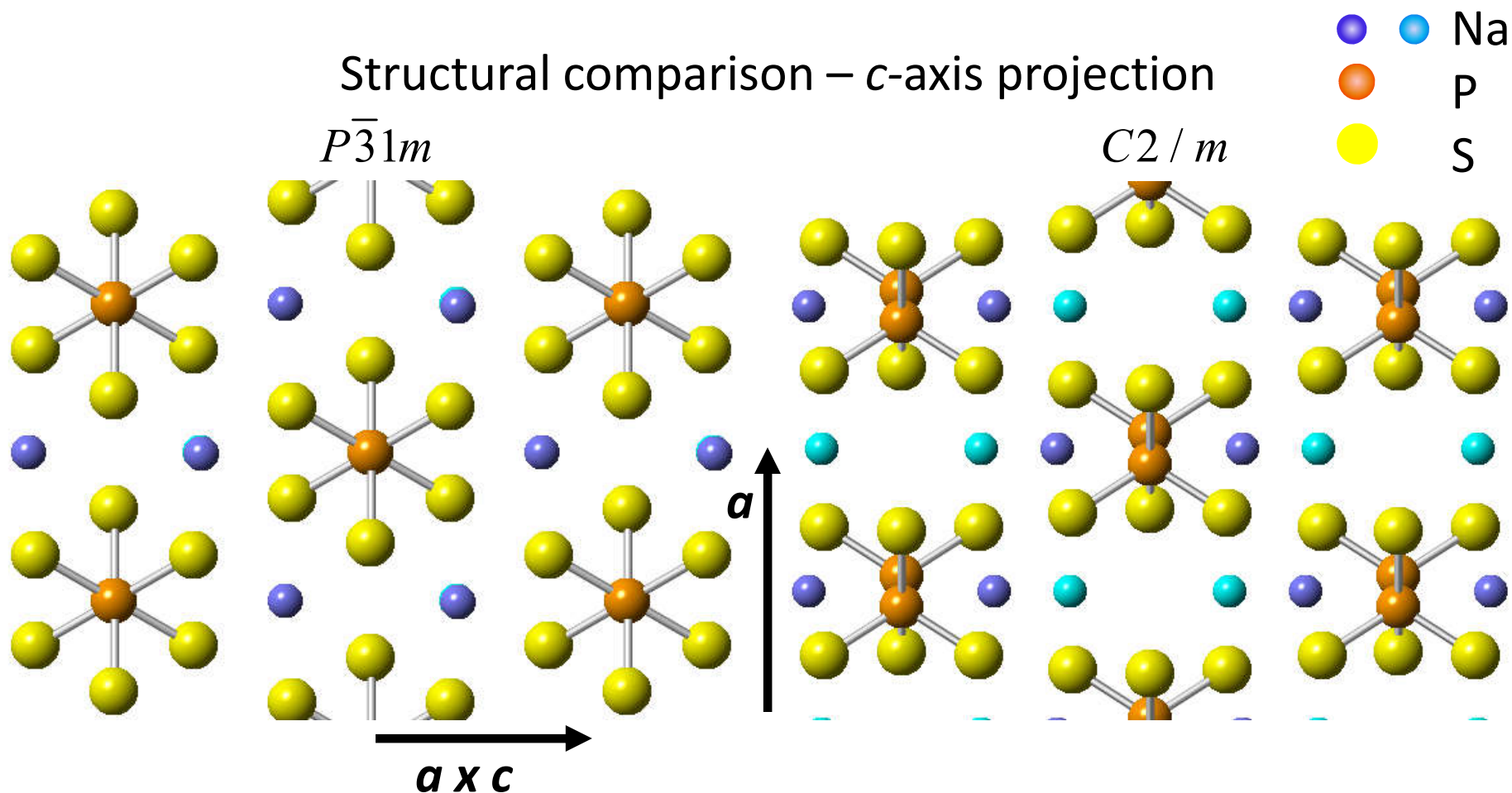
$\Delta E = 0$

50% P↑
50% P↓

Two model configurations of disordered ground state structure

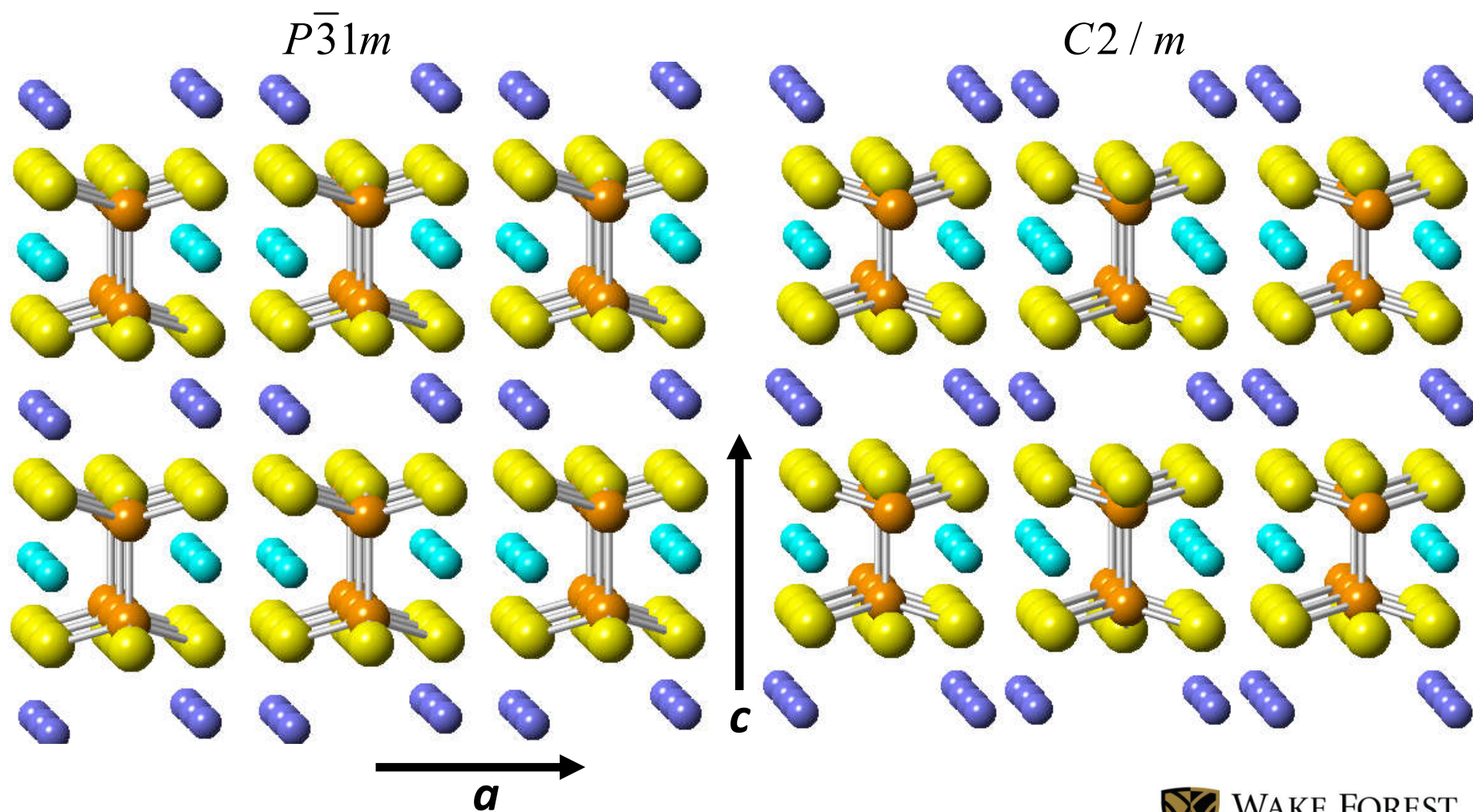
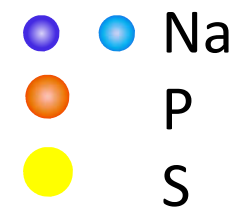
Structure of $\text{Na}_4\text{P}_2\text{S}_6$:

Kuhn et al., ZAAC **640**, 689-692 (2014) synthesized single crystals with a monoclinic structure having space group $C2/m$ with similarities to the trigonal structure with $P\bar{3}1m$ space group





Structural comparison – view including c-axis



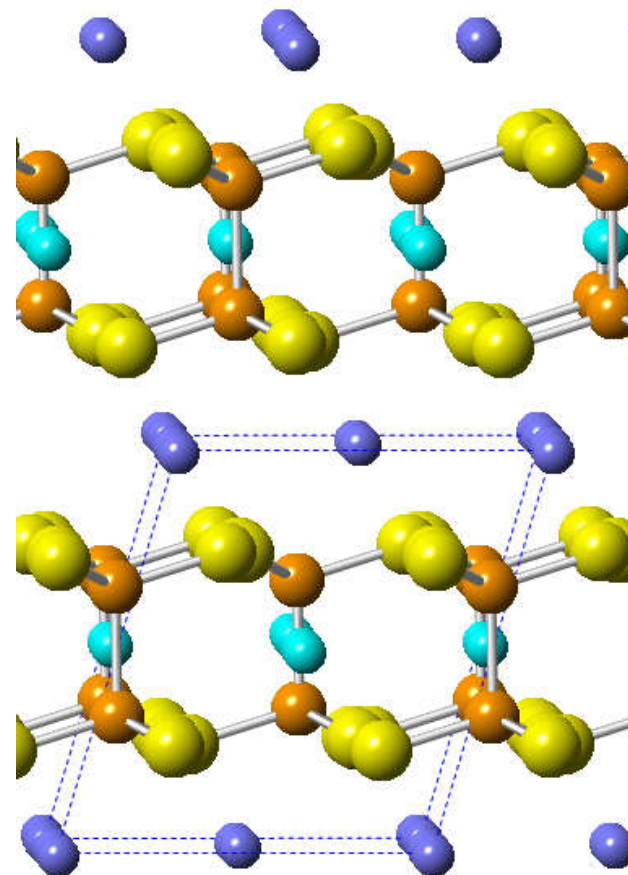
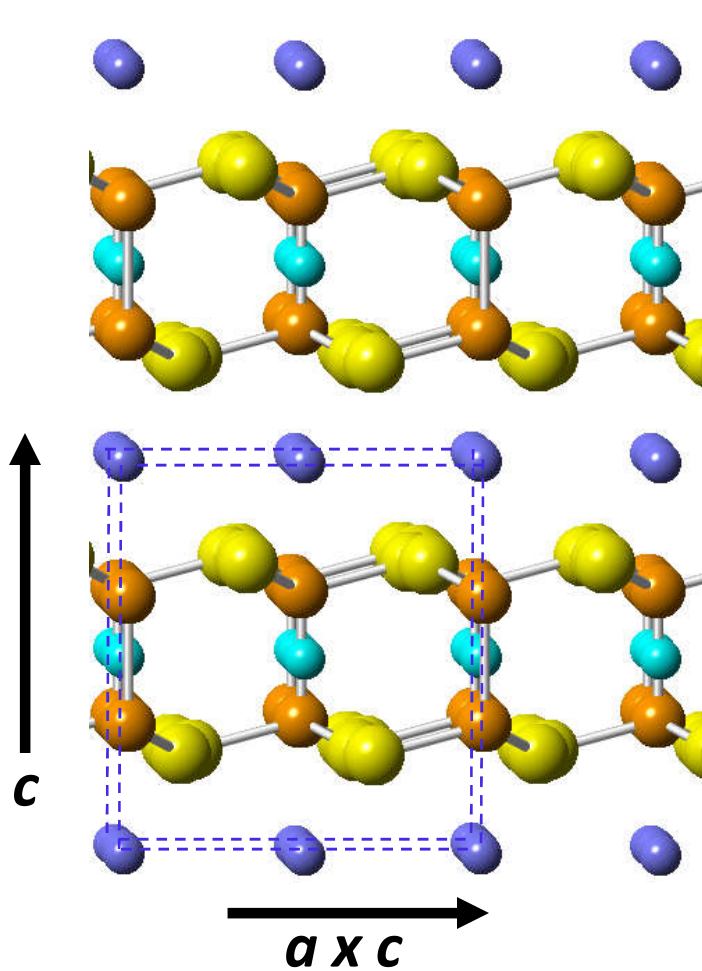


Structural comparison – view including c-axis

- Na
- P
- S

$P\bar{3}1m$

$C2/m$



Results for $\text{Na}_4\text{P}_2\text{S}_6$:

Calculated heats of formation (eV per formula unit) for $\text{Na}_4\text{P}_2\text{S}_6$ and $\text{Li}_4\text{P}_2\text{S}_6$ in 4 structural models

	$\text{Na}_4\text{P}_2\text{S}_6$	$\text{Li}_4\text{P}_2\text{S}_6$
Kuhn structure	-11.47 eV	-12.07 eV
Structure "b"	-11.47 eV	-12.42 eV
Structure "c"	-11.56 eV	-12.46 eV
Structure "d"	-11.56 eV	-12.46 eV

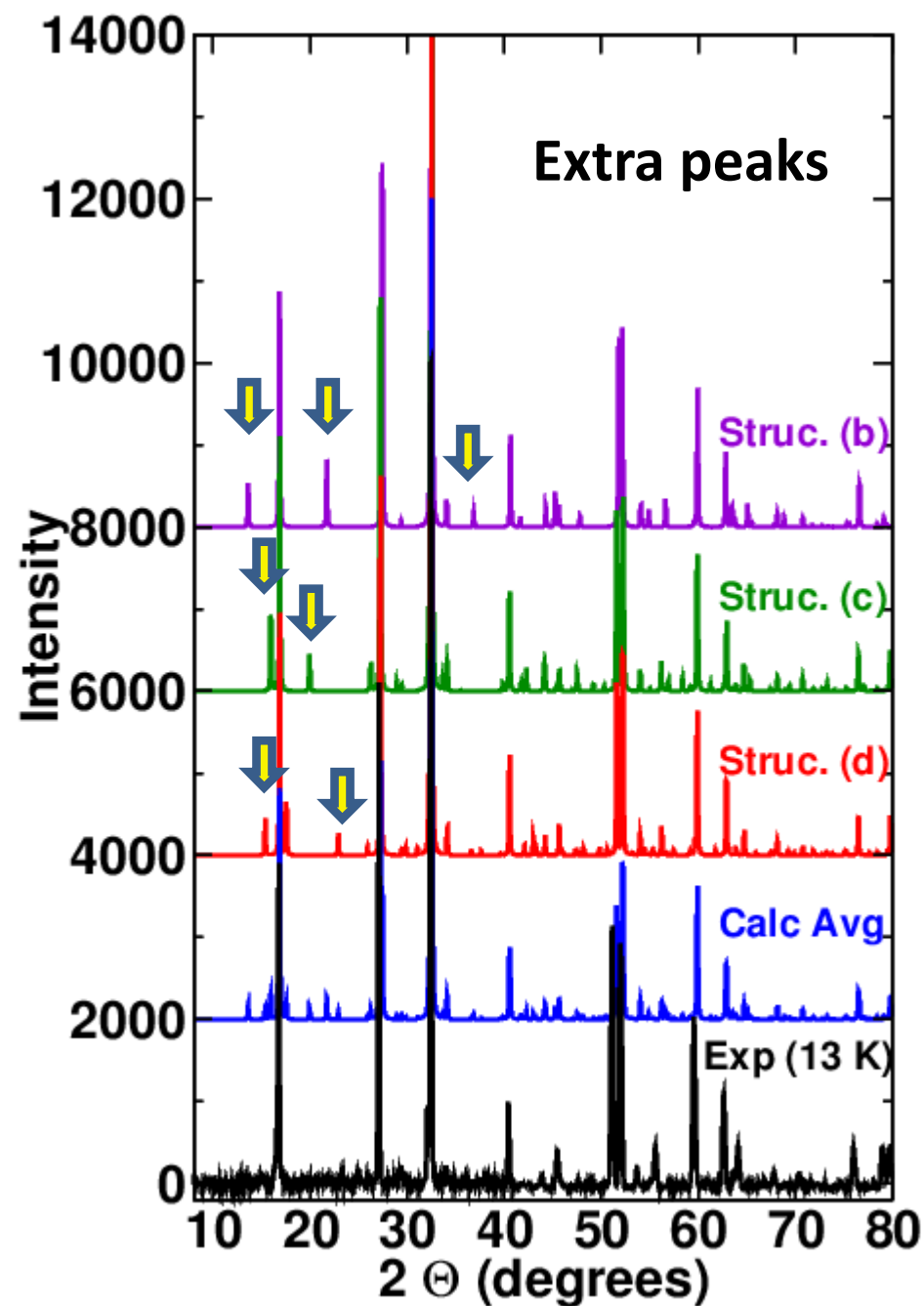
Models of disordered Mercier structure

→ Calculations find the most stable structure for both $\text{Na}_4\text{P}_2\text{S}_6$ and $\text{Li}_4\text{P}_2\text{S}_6$ to be the disordered Mercier structure, suggesting that the Kuhn structure is meta-stable.

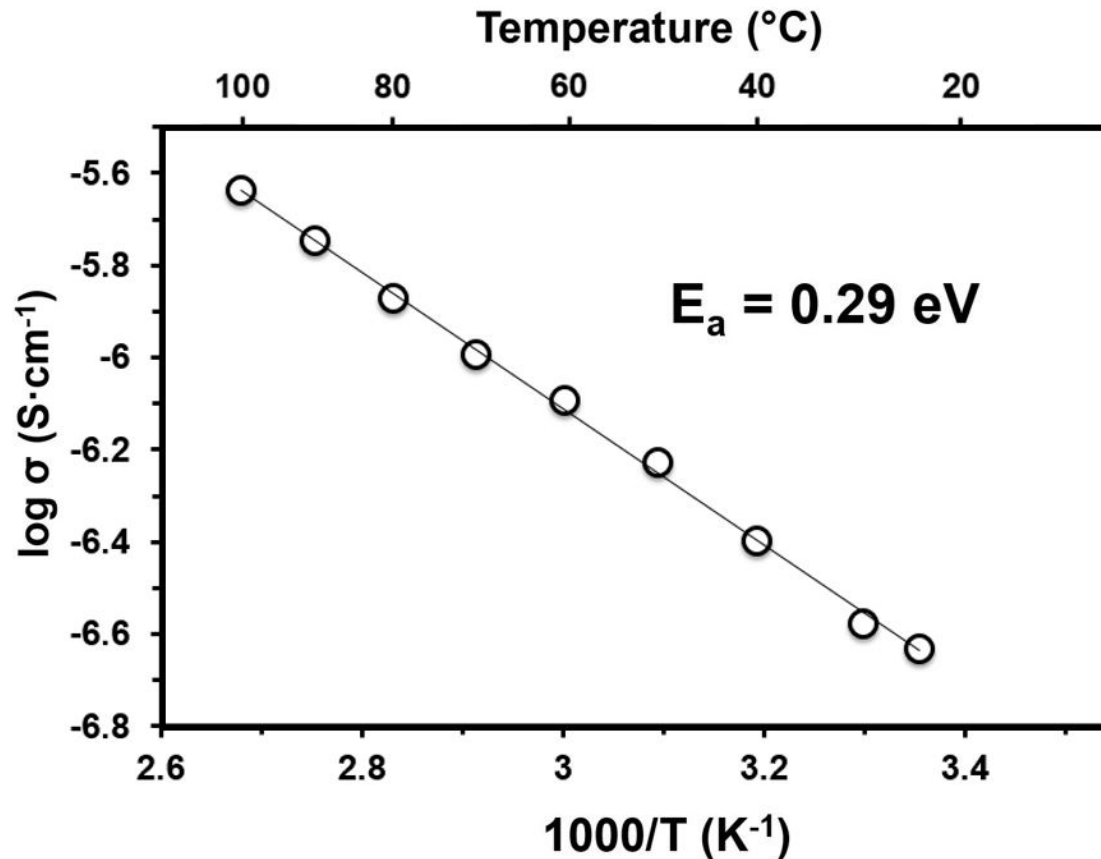
Comparison of X-ray data for $\text{Li}_4\text{P}_2\text{S}_6$ with simulations

Note: simulations scaled by 102% to compensate for systematic LDA error.

Simulations consistent with incoherent average over all $\text{P}\uparrow$ and $\text{P}\downarrow$ configurations



Ionic conductivity and Activation Energy for $\text{Li}_4\text{P}_2\text{S}_6$

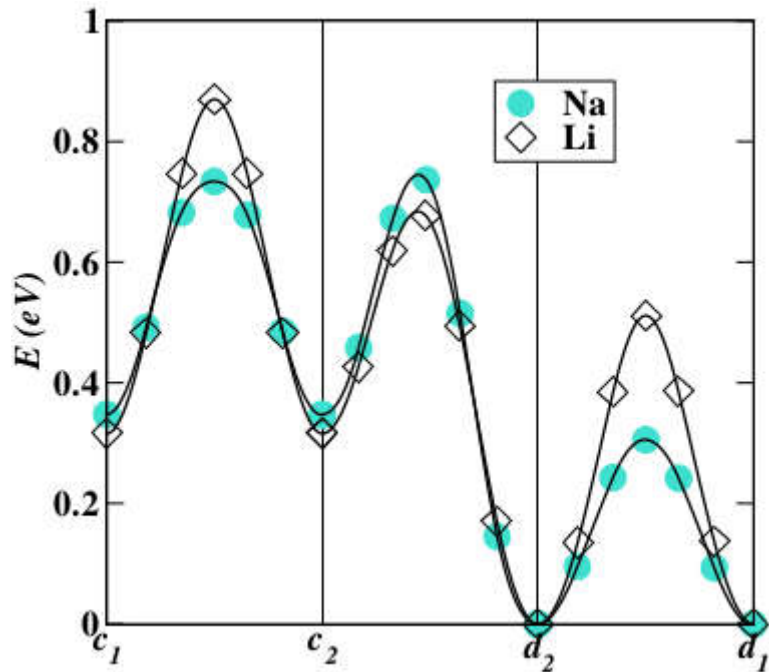


$2.38 \times 10^{-7} \text{ S/cm}$ at 25°C and $2.33 \times 10^{-6} \text{ S/cm}$ at 100°C
 $\text{Li}_4\text{P}_2\text{S}_6$ pressed pellets with blocking (Al/C) electrodes

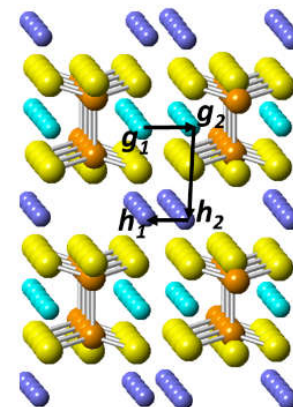
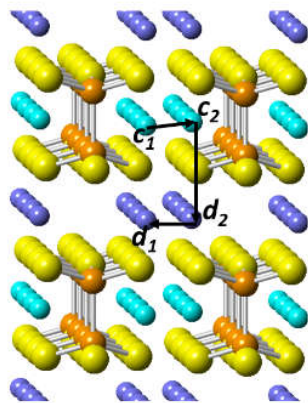
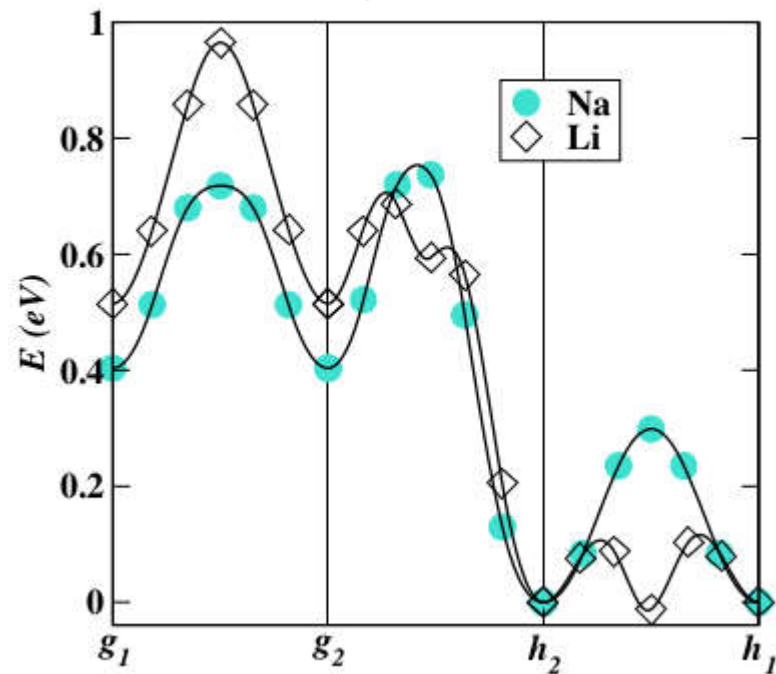
Li/ $\text{Li}_4\text{P}_2\text{S}_6$ /Li cells could not be cycled

Comparison of vacancy migration of $\text{Na}_4\text{P}_2\text{S}_6$ and $\text{Li}_4\text{P}_2\text{S}_6$

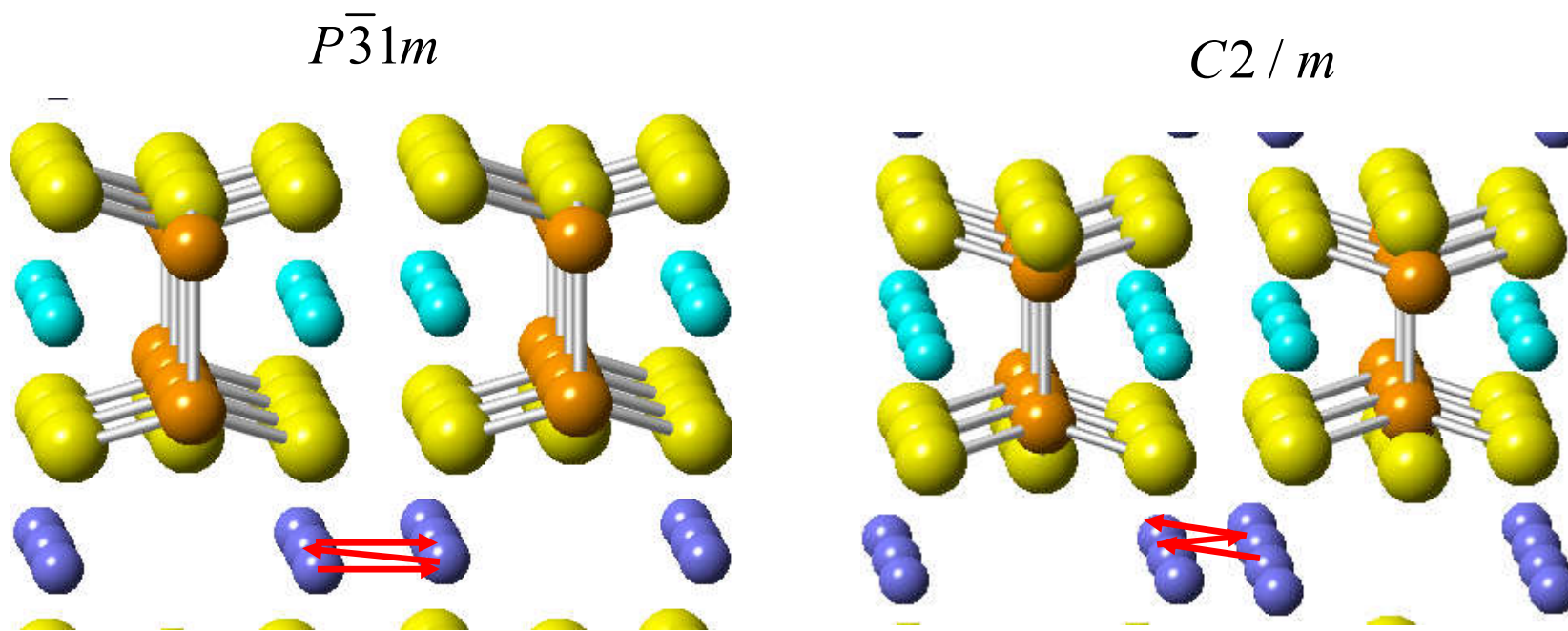
$P\bar{3}1m$



$C2/m$

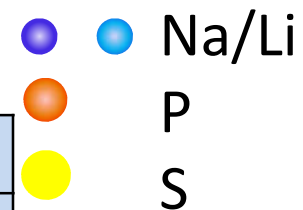


- Na/Li
- P
- S



Minimum ion vacancy migration energies

	$P\bar{3}1m$	$C2/m$
$\text{Na}_4\text{P}_2\text{S}_6$	0.3 eV	0.3 eV
$\text{Li}_4\text{P}_2\text{S}_6$	0.5 eV	0.1 eV



Conclusions on studies of $\text{Li}_4\text{P}_2\text{S}_6$ and $\text{Na}_4\text{P}_2\text{S}_6$:

- $\text{Li}_4\text{P}_2\text{S}_6$ and $\text{Na}_4\text{P}_2\text{S}_6$ have interesting structural properties; simulations find the most stable structure for both to be the disordered Mercier structure, suggesting that the Kuhn structure is meta-stable.
- Experimental structural studies for $\text{Li}_4\text{P}_2\text{S}_6$ agree with the simulations; material is found to be remarkably temperature independent and thermally stable relative to other thio-phosphates.
- Measurements find $\text{Li}_4\text{P}_2\text{S}_6$ to have low ionic conductivity; simulations suggest that $\text{Na}_4\text{P}_2\text{S}_6$ may have more favorable ionic conductivity.
- Models of ideal $\text{Li}_4\text{P}_2\text{S}_6/\text{Li}$ interfaces find broken P—S bonds; $\text{Na}_4\text{P}_2\text{S}_6/\text{Na}$ interfaces in the Kuhn structure may be slightly more stable

Additional thoughts

- **Limitations of first principles modeling**
 - Small simulation cells**
 - Zero temperature**
- **Possible extensions**
 - Develop approximation schemes for treatment of larger supercells**
 - Use molecular dynamics and/or Monte Carlo techniques**
- **Ideal research effort in materials includes close collaboration of both simulations and experimental measurements.**
- **For battery technology, there remain many opportunities for new materials development.**