

**PHY 711 Classical Mechanics and  
Mathematical Methods  
10-10:50 AM MWF Olin 103**

**Plan for Lecture 22:**

**Summary of mathematical methods**

- 1. Fourier transforms**
- 2. Fast Fourier transforms**

	Mon, 9/09/2013	Chap. 3	Calculus of variations — continued	
7	Wed, 9/11/2013	Chap. 3	Calculus of variations applied to Lagrangians	<a href="#">#6</a>
8	Fri, 9/13/2013	Chap. 3	Lagrangian mechanics	<a href="#">#7</a>
9	Mon, 9/16/2013	Chap. 3 & 6	Lagrangian mechanics	<a href="#">#8</a>
10	Wed, 9/18/2013	Chap. 3 & 6	Lagrangian mechanics	<a href="#">#9</a>
11	Fri, 9/20/2013	Chap. 3 & 6	Lagrangian & Hamiltonian mechanics	<a href="#">#10</a>
12	Mon, 9/23/2013	Chap. 3 & 6	Hamiltonian formalism	<a href="#">#11</a>
13	Wed, 9/25/2013	Chap. 3 & 6	Hamiltonian formalism	<a href="#">#12</a>
14	Fri, 9/27/2013	Chap. 3 & 6	Hamiltonian formalism	<a href="#">#13</a>
15	Mon, 9/30/2013	Chap. 4	Small Oscillations	<a href="#">#14</a>
16	Wed, 10/02/2013	Chap. 4	Small Oscillations	
17	Fri, 10/04/2013	Chap. 4	Small Oscillations	<a href="#">#15</a>
18	Mon, 10/07/2013	Chap. 4 & 7	Small Oscillations and waves	<a href="#">#16</a>
19	Wed, 10/09/2013	Chap. 7	Wave equation	
	Fri, 10/11/2013		No class (Fall Break)	
20	Mon, 10/14/2013	Chap. 7	Wave equation (Presentation topic due)	<a href="#">#17</a>
21	Wed, 10/16/2013	Chap. 7	Mathematical methods	<a href="#">#18</a>
22	Fri, 10/18/2013	Chap. 5	Mathematical methods	<a href="#">#19</a>
23	Mon, 10/21/2013	Chap. 5	Rigid rotations	
24	Wed, 10/23/2013	Chap. 5	Rigid rotations	
25	Fri, 10/25/2013	Chap. 5	Rigid rotations	
	Mon, 10/28/2013	No class	Take-home exam	
	Wed, 10/30/2013	No class	Take-home exam	
	Fri, 11/01/2013	No class	Take-home exam	
26	Mon, 11/04/2013		Take-home exam due	



## Use of Fourier transforms to solve wave equation

$$\text{Wave equation : } \frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0$$

Suppose  $u(x, t) = e^{-i\omega t} \tilde{F}(x, \omega)$  where  $\tilde{F}(x, \omega)$  satisfies the equation :

$$\frac{\partial^2 \tilde{F}(x, \omega)}{\partial x^2} = -\frac{\omega^2}{c^2} \tilde{F}(x, \omega) \equiv -k^2 \tilde{F}(x, \omega)$$

Further assume that fixed boundary conditions apply :  $0 \leq x \leq L$

$$\text{with } \tilde{F}(0, \omega) = 0 \quad \text{and} \quad \tilde{F}(L, \omega) = 0$$

For  $n = 1, 2, 3, \dots$

$$\tilde{F}_n(x, \omega) = \sin\left(\frac{n\pi x}{L}\right) \quad k \rightarrow k_n = \frac{n\pi}{L} \equiv \frac{\omega_n}{c}$$

$$u(x, t) = e^{-i\omega_n t} \sin(k_n x) = e^{-i\omega_n t} \frac{(e^{ik_n x} - e^{-ik_n x})}{2i} = \frac{(e^{ik_n(x-ct)} - e^{-ik_n(x+ct)})}{2i}$$

## Use of Fourier transforms to solve wave equation -- continued

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0$$

Using superposition : Suppose  $u(x, t) = \sum_n C_n e^{-i\omega_n t} \tilde{F}_n(x, \omega_n)$

$$\frac{\partial^2 \tilde{F}_n(x, \omega_n)}{\partial x^2} = -\frac{\omega_n^2}{c^2} \tilde{F}_n(x, \omega_n) \equiv -k_n^2 \tilde{F}_n(x, \omega_n)$$

$$\text{For } \tilde{F}_n(x, \omega) = \sin\left(\frac{n\pi x}{L}\right) \quad k \rightarrow k_n = \frac{n\pi}{L} \equiv \frac{\omega_n}{c}$$

$$\begin{aligned} \Rightarrow u(x, t) &= \sum_n C_n e^{-i\omega_n t} \sin(k_n x) = \sum_n \frac{C_n}{2i} e^{-i\omega_n t} (e^{ik_n x} - e^{-ik_n x}) \\ &= \sum_n \frac{C_n}{2i} (e^{ik_n(x-ct)} - e^{-ik_n(x+ct)}) \equiv f(x-ct) + g(x+ct) \end{aligned}$$

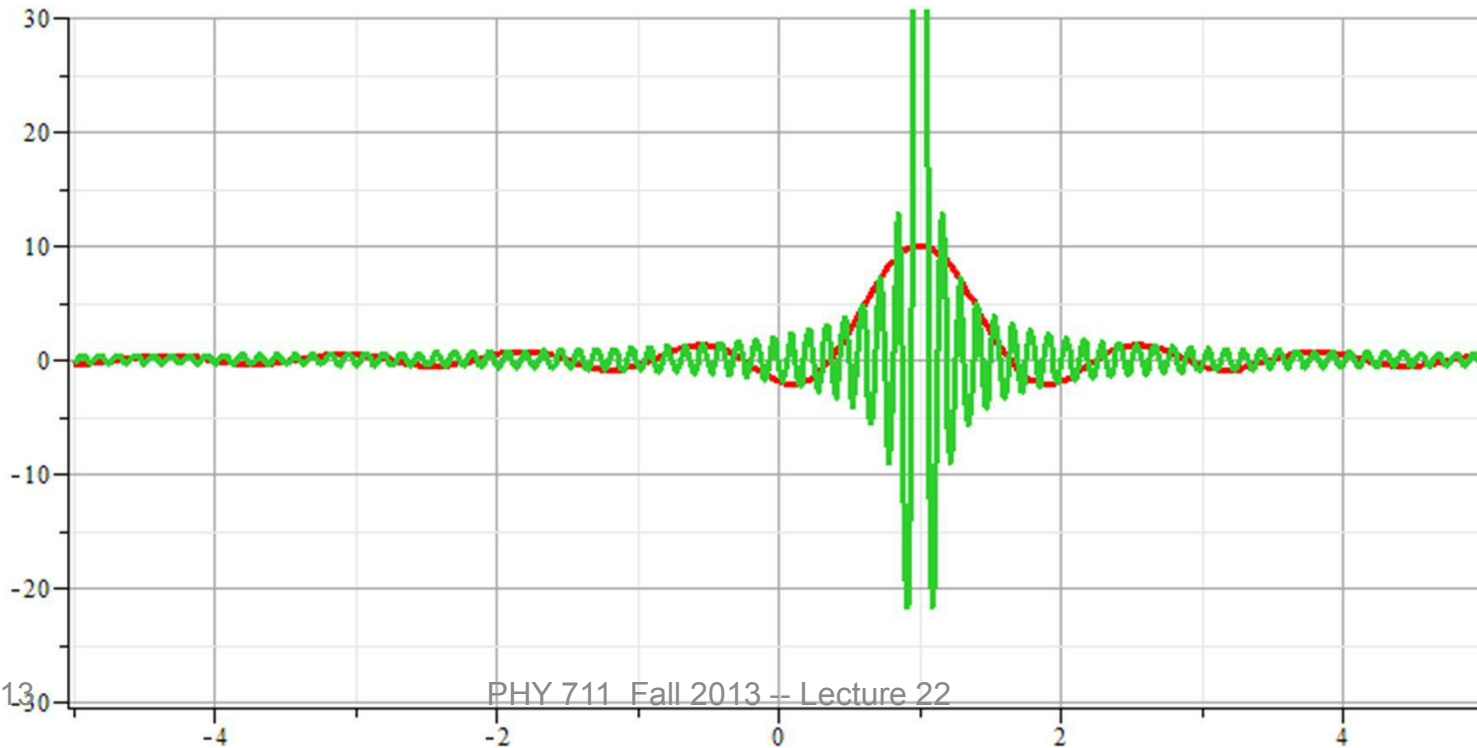
# Fourier transforms

A useful identity

$$\int_{-\infty}^{\infty} dt e^{-i(\omega - \omega_0)t} = 2\pi\delta(\omega - \omega_0)$$

Note that

$$\int_{-T}^T dt e^{-i(\omega - \omega_0)t} = \frac{2 \sin[(\omega - \omega_0)T]}{\omega - \omega_0}$$



Definition of Fourier Transform for a function  $f(t)$ :

$$f(t) = \int_{-\infty}^{\infty} d\omega F(\omega) e^{-i\omega t}$$

Backward transform:

$$F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dt f(t) e^{i\omega t}$$

Check:

$$f(t) = \int_{-\infty}^{\infty} d\omega \left( \frac{1}{2\pi} \int_{-\infty}^{\infty} dt' f(t') e^{i\omega t'} \right) e^{-i\omega t}$$

$$f(t) = \int_{-\infty}^{\infty} dt' f(t') \left( \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{i\omega(t'-t)} \right) = \int_{-\infty}^{\infty} dt' f(t') \delta(t'-t)$$

Properties of Fourier transforms -- Parseval's theorem :

$$\int_{-\infty}^{\infty} dt (f(t))^* f(t) = \int_{-\infty}^{\infty} d\omega (F(\omega))^* F(\omega)$$

Fourier transform for periodic function :

Suppose  $f(t + nT) = f(t)$  for any integer  $n$

$$F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dt f(t) e^{i\omega t} = \sum_{n=-\infty}^{\infty} \left( \int_0^T dt f(t) e^{i\omega(t+nT)} \right)$$

Note that :

$$\sum_{n=-\infty}^{\infty} e^{in\omega T} = \Omega \sum_{\nu=-\infty}^{\infty} \delta(\omega - \nu\Omega), \quad \text{where } \Omega \equiv \frac{2\pi}{T}$$

Some details :

$$\sum_{n=-M}^M e^{in\omega T} = \frac{\sin\left(\left(M + \frac{1}{2}\right)\omega T\right)}{\sin\left(\frac{1}{2}\omega T\right)}$$

$$\lim_{M \rightarrow \infty} \left( \frac{\sin\left(\left(M + \frac{1}{2}\right)\omega T\right)}{\sin\left(\frac{1}{2}\omega T\right)} \right) = 2\pi \sum_{\nu} \delta(\omega T - \nu\Omega T) = \frac{2\pi}{T} \sum_{\nu} \delta(\omega - \nu\Omega)$$

$$\Rightarrow \sum_{n=-\infty}^{\infty} e^{in\omega T} = \Omega \sum_{\nu=-\infty}^{\infty} \delta(\omega - \nu\Omega), \quad \text{where } \Omega \equiv \frac{2\pi}{T}$$

$$\begin{aligned} \Rightarrow F(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dt f(t) e^{i\omega t} = \sum_{\nu=-\infty}^{\infty} \Omega \delta(\omega - \nu\Omega) \left( \int_0^T dt f(t) e^{i\omega t} \right) \\ &\equiv \sum_{\nu=-\infty}^{\infty} \Omega \delta(\omega - \nu\Omega) \bar{F}(\nu\Omega) \end{aligned}$$



## Fourier transform for periodic function – continued:

$$f(t) = \int_{-\infty}^{\infty} d\omega F(\omega) e^{-i\omega t} \quad \text{where } f(t) = f(t + nT)$$

$$\begin{aligned} \Rightarrow F(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dt f(t) e^{i\omega t} = \sum_{\nu=-\infty}^{\infty} \Omega \delta(\omega - \nu\Omega) \left( \int_0^T dt f(t) e^{i\omega t} \right) \\ &\equiv \sum_{\nu=-\infty}^{\infty} \Omega \delta(\omega - \nu\Omega) \bar{F}(\nu\Omega) \end{aligned}$$

$$\Omega \equiv \frac{2\pi}{T}$$

Thus, for a periodic function

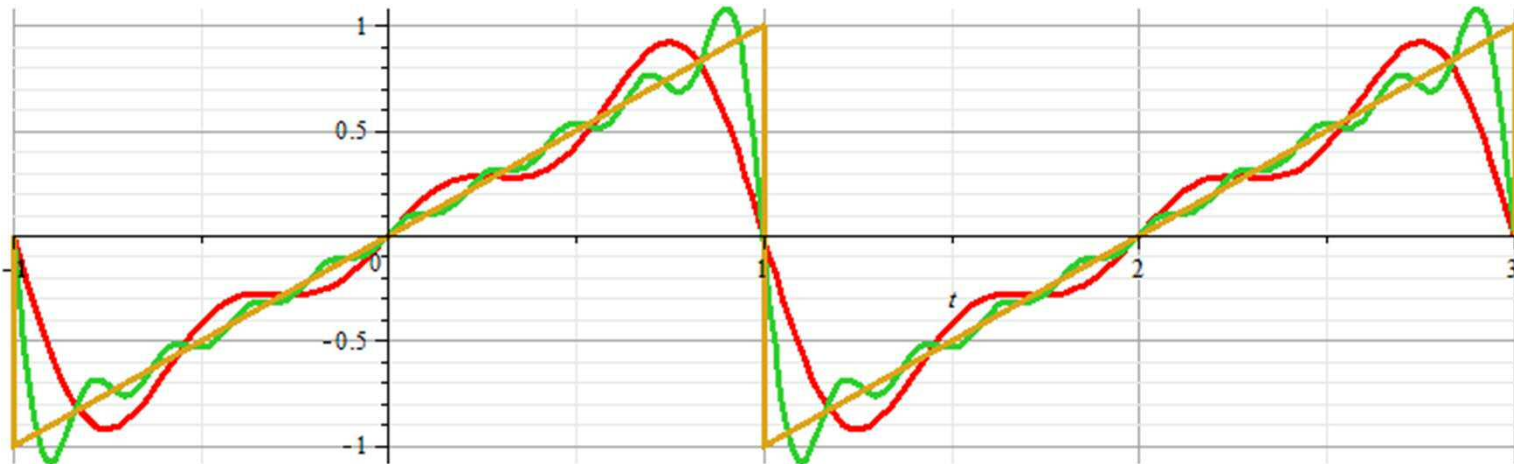
$$f(t) = \Omega \sum_{\nu=-\infty}^{\infty} \bar{F}(\nu\Omega) e^{-i\nu\Omega t}$$

Example:

Suppose:  $f(t) = \frac{t - nT}{T}$  for  $(n-1)T \leq t \leq (n+1)T$ ;  $n = 0, 2, 4, 6 \dots$

Note, in this case the repeat period is  $2T$  and the convenient sample time interval is  $-T \leq t \leq T$ .

$$\bar{F}(v\Omega) = \frac{2\pi}{2T} i \int_{-T}^T \frac{t}{T} \sin\left(\frac{v\pi t}{2T}\right) dt \quad f(t) = \sum_{v=1}^{\infty} 2|\bar{F}(v\Omega)| \sin\left(\frac{v\pi t}{2T}\right)$$



Thus, for a periodic function

$$f(t) = \sum_{\nu=-\infty}^{\infty} F(\nu\Omega) e^{-i\nu\Omega t}$$

Now suppose that the transformed function is bounded;

$$|F(\nu\Omega)| \leq \varepsilon \quad \text{for } |\nu| \geq N$$

Define a periodic transform function

$$\tilde{F}(\nu\Omega) \equiv \tilde{F}(\nu\Omega + \nu'((2N+1)\Omega))$$

Effect on time domain :

$$f(t) = \sum_{\nu=-\infty}^{\infty} \tilde{F}(\nu\Omega) e^{-i\nu\Omega t} = \frac{2\pi}{(2N+1)\Omega} \sum_{\nu=-N}^N \tilde{F}(\nu\Omega) e^{-i\nu\Omega t} \sum_{\mu} \delta\left(t - \frac{\mu T}{2N+1}\right)$$

# Doubly periodic functions

$$t \rightarrow \frac{\mu T}{2N+1}$$

$$\tilde{f}_\mu = \frac{1}{2N+1} \sum_{\nu=-N}^N \tilde{F}_\nu e^{-i2\pi\nu\mu/(2N+1)}$$

$$\tilde{F}_\nu = \sum_{\mu=-N}^N \tilde{f}_\mu e^{i2\pi\nu\mu/(2N+1)}$$

# More convenient notation

$$2N + 1 \rightarrow M$$

$$\tilde{f}_\mu = \frac{1}{M} \sum_{\nu=0}^{M-1} \tilde{F}_\nu e^{-i2\pi\nu\mu/M}$$

$$\tilde{F}_\nu = \sum_{\mu=0}^{M-1} \tilde{f}_\mu e^{i2\pi\nu\mu/M}$$

Note that for  $W = e^{i2\pi/M}$

$$\tilde{F}_0 = \tilde{f}_0 W^0 + \tilde{f}_1 W^0 + \tilde{f}_2 W^0 + \tilde{f}_3 W^0 + \dots$$

$$\tilde{F}_1 = \tilde{f}_0 W^0 + \tilde{f}_1 W^1 + \tilde{f}_2 W^2 + \tilde{f}_3 W^3 + \dots$$

$$\tilde{F}_2 = \tilde{f}_0 W^0 + \tilde{f}_1 W^2 + \tilde{f}_2 W^4 + \tilde{f}_3 W^6 + \dots$$

Note that for  $W = e^{i2\pi/M}$

$$\tilde{F}_0 = \tilde{f}_0 W^0 + \tilde{f}_1 W^0 + \tilde{f}_2 W^0 + \tilde{f}_3 W^0 + \dots$$

$$\tilde{F}_1 = \tilde{f}_0 W^0 + \tilde{f}_1 W^1 + \tilde{f}_2 W^2 + \tilde{f}_3 W^3 + \dots$$

$$\tilde{F}_2 = \tilde{f}_0 W^0 + \tilde{f}_1 W^2 + \tilde{f}_2 W^4 + \tilde{f}_3 W^6 + \dots$$

$$\text{However, } W^M = \left( e^{i2\pi/M} \right)^M = 1$$

$$\text{and } W^{M/2} = \left( e^{i2\pi/M} \right)^{M/2} = -1$$

Cooley-Tukey algorithm: J. W. Cooley and J. W. Tukey, "An algorithm for machine calculation of complex Fourier series" Math. Computation 19, 297-301 (1965)