

**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF Olin 103**

Plan for Lecture 24:

Rigid body rotational motion (Chap. 5)

- 1. Torque free**
- 2. Euler angles**
- 3. Motion of a symmetric top**



Course schedule

(Preliminary schedule – subject to frequent adjustment.)

Date	F&W	Reading	Topic	Assignment
1 Wed. 8/29/2012	Chap. 1	Review of basic principles	Scattering theory	#1
2 Fri. 8/31/2012	Chap. 1	Scattering theory continued		#2
3 Mon. 9/03/2012	Chap. 1	Scattering theory continued		#3
4 Wed. 9/05/2012	Chap. 1 & 2	Scattering theory/Accelerated coordinate frame		#4
5 Fri. 9/07/2012	Chap. 2	Accelerated coordinate frame		#5
6 Mon. 9/10/2012	Chap. 3	Calculus of Variation		#6
7 Wed. 9/12/2012	Chap. 3	Calculus of Variation continued		#7
8 Fri. 9/14/2012	Chap. 3	Lagrangian		#8
9 Mon. 9/17/2012	Chap. 3 & 6	Lagrangian		#9
10 Wed. 9/19/2012	Chap. 3 & 6	Lagrangian		#10
11 Fri. 9/21/2012	Chap. 3 & 6	Lagrangian		#11
12 Mon. 9/24/2012	Chap. 3 & 6	Lagrangian and Hamiltonian		#12
13 Wed. 9/26/2012	Chap. 6	Lagrangian and Hamiltonian		#13
14 Fri. 9/28/2012	Chap. 6	Lagrangian and Hamiltonian		#14
15 Mon. 10/01/2012	Chap. 4	Small oscillations		#15
16 Wed. 10/03/2012	Chap. 4	Small oscillations		
17 Fri. 10/05/2012	Chap. 4	Small oscillations		
18 Mon. 10/08/2012	Chap. 7	Wave equation		Take Home Exam
19 Wed. 10/10/2012	Chap. 7	Wave equation		Take Home Exam
20 Fri. 10/12/2012	Chap. 7	Wave equation		Take Home Exam
21 Mon. 10/15/2012	Chap. 7	Wave equation		Exam due
22 Wed. 10/17/2012	Chap. 7, 5	Moment of inertia		
23 Fri. 10/19/2012		Fall break		
24 Mon. 10/22/2012	Chap. 5	Rigid body rotation		#16
25 Wed. 10/24/2012	Chap. 5	Rigid body rotation		#17

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News

Workshop for Middle School Teachers Organized by Prof. Cho is featured in Mashable, Huffington Post, and Fox 8 News

Article in WS Journal on Tech Expo features Beet-Road Juice

Article by Lacra Neoureanu of the Salsbury Group Selected for Inaugural Contribution to Protopedia from JISO

Events

Wed Oct 24, 2012
Prof. Pablo Laguna Georgia Tech Back Homes
 4:00 PM in Olin 101
 Refreshments at 3:30 in Lobby

Wed Oct 31, 2012
Dr. Paul Kent Olin
 Lithium Ion Batteries
 4:00 PM in Olin 101
 Refreshments at 3:30 in Lobby

Oct 29-30, 2012
Sutcuart NanoDays
 NanoCarbon Technology Conference

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FOREST
UNIVERSITY

Department of Physics

WFU Physics Colloquium

TITLE: Black holes and gravitational waves: The quest to show if Einstein was right

SPEAKER: Professor Pablo Laguna,
*Center for Relativistic Astrophysics
 School of Physics
 Georgia Institute of Technology*

TIME: Wednesday October 24, 2012

PLACE: Room 101 Olin Physical Laboratory

Refreshments will be served at 3:30 PM in the Olin Lounge. All interested persons are cordially invited to attend.

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Comment about Exam problem #1:

Consider a Rutherford scattering experiment analyzed in the center of mass frame of reference in which the ratio of the (repulsive) potential to the center-of-mass energy is described by a modified interaction of the form:

$$\frac{V(r)}{E} = \begin{cases} \frac{\kappa}{r} & \text{for } r < r_{\max} \\ 0 & \text{for } r > r_{\max}. \end{cases}$$

(a) Using the approach described in Problem 1.15 of your textbook, find an expression for $\theta(b)$, the scattering angle θ as a function of the impact parameter b . (Simplify the expression, but you need not explicitly solve for b .)

(b) Show that in the limit $r_{\max} \rightarrow \infty$, your result is consistent with the Rutherford scattering result.

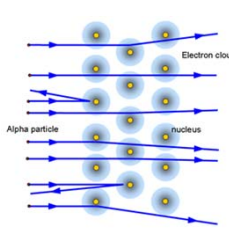
(c) For the modified interaction, assume that

$$\frac{\kappa}{r_{\max}} \gg \frac{b}{r_{\max}},$$

so that terms including (b/r_{\max}) may be neglected. Find the form of the scattering cross section for this case and compare it with the pure Rutherford scattering cross section.

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Motivation for model:



Au target particles consist of Au nuclei which are screened by electrons.

http://tap.iop.org/atoms/rutherford/img_full_47190.gif

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Calculation of the scattering angle as a function of impact parameter b :

$$\begin{aligned} \theta(b) &= \pi - 2b \int_{r_{\min}}^{r_{\max}} \frac{1}{r\sqrt{r^2 - \kappa r - b^2}} - 2b \int_{r_{\max}}^{\infty} \frac{1}{r\sqrt{r^2 - b^2}} \\ &= \pi - 2 \tan^{-1} \left(\frac{-2b^2 - \kappa r}{2b\sqrt{r^2 - \kappa r - b^2}} \right) \Bigg|_{r_{\min}}^{r_{\max}} - 2 \tan^{-1} \left(\frac{-2b^2}{2b\sqrt{r^2 - b^2}} \right) \Bigg|_{r_{\max}}^{\infty} \\ &= 2 \tan^{-1} \left(\frac{2b^2 + \kappa r_{\max}}{2b\sqrt{r_{\max}^2 - \kappa r_{\max} - b^2}} \right) + 2 \tan^{-1} \left(\frac{2b^2}{2b\sqrt{r_{\max}^2 - b^2}} \right) \\ &= 2 \tan^{-1} \left(\frac{(b/r_{\max}) + (\kappa/2b)}{\sqrt{1 - (\kappa/r_{\max}) - (b/r_{\max})^2}} \right) + 2 \tan^{-1} \left(\frac{(b/r_{\max})}{\sqrt{1 - (b/r_{\max})^2}} \right) \end{aligned}$$

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$$\begin{aligned} \tan\left(\frac{\theta}{2} - \alpha\right) &= \frac{(b/r_{\max}) + (\kappa/2b)}{\sqrt{1 - (\kappa/r_{\max}) - (b/r_{\max})^2}} \\ \tan^2\left(\frac{\theta}{2} - \alpha\right) &= \frac{(b/r_{\max})^2 + (\kappa/r_{\max}) + (\kappa/2b)^2}{1 - (\kappa/r_{\max}) - (b/r_{\max})^2} \\ \text{where } \tan \alpha &= \frac{(b/r_{\max})}{\sqrt{1 - (b/r_{\max})^2}} \end{aligned}$$

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Back to rotational motion
 Time rate of change of angular momentum

$$\frac{d\mathbf{L}}{dt} = \left(\frac{d\mathbf{L}}{dt} \right)_{\text{body}} + \boldsymbol{\omega} \times \mathbf{L}$$

For (body fixed) coordinate system that diagonalizes moment of inertia tensor :

$$\begin{aligned} \vec{\mathbf{I}} \cdot \hat{\mathbf{e}}_i &= I_i \hat{\mathbf{e}}_i & \boldsymbol{\omega} &= \tilde{\omega}_1 \hat{\mathbf{e}}_1 + \tilde{\omega}_2 \hat{\mathbf{e}}_2 + \tilde{\omega}_3 \hat{\mathbf{e}}_3 \\ \mathbf{L} &= I_1 \tilde{\omega}_1 \hat{\mathbf{e}}_1 + I_2 \tilde{\omega}_2 \hat{\mathbf{e}}_2 + I_3 \tilde{\omega}_3 \hat{\mathbf{e}}_3 \\ \frac{d\mathbf{L}}{dt} &= I_1 \dot{\tilde{\omega}}_1 \hat{\mathbf{e}}_1 + I_2 \dot{\tilde{\omega}}_2 \hat{\mathbf{e}}_2 + I_3 \dot{\tilde{\omega}}_3 \hat{\mathbf{e}}_3 + \tilde{\omega}_2 \tilde{\omega}_3 (I_3 - I_2) \hat{\mathbf{e}}_1 \\ &\quad + \tilde{\omega}_3 \tilde{\omega}_1 (I_1 - I_3) \hat{\mathbf{e}}_2 + \tilde{\omega}_1 \tilde{\omega}_2 (I_2 - I_1) \hat{\mathbf{e}}_3 \end{aligned}$$

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Descriptions of rotation about a given origin -- continued

Note that the torque equation

$$\frac{d\mathbf{L}}{dt} = \left(\frac{d\mathbf{L}}{dt} \right)_{body} + \boldsymbol{\omega} \times \mathbf{L} = \boldsymbol{\tau}$$

is very difficult to solve directly in the body fixed frame.
 For $\boldsymbol{\tau} = 0$ we can solve the Euler equations :

$$\begin{aligned} \frac{d\mathbf{L}}{dt} = I_1 \dot{\tilde{\omega}}_1 \hat{\mathbf{e}}_1 + I_2 \dot{\tilde{\omega}}_2 \hat{\mathbf{e}}_2 + I_3 \dot{\tilde{\omega}}_3 \hat{\mathbf{e}}_3 + \tilde{\omega}_2 \tilde{\omega}_3 (I_3 - I_2) \hat{\mathbf{e}}_1 \\ + \tilde{\omega}_3 \tilde{\omega}_1 (I_1 - I_3) \hat{\mathbf{e}}_2 + \tilde{\omega}_1 \tilde{\omega}_2 (I_2 - I_1) \hat{\mathbf{e}}_3 = 0 \end{aligned}$$

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Euler equations for rotation in body fixed frame :

$$\begin{aligned} I_1 \dot{\tilde{\omega}}_1 + \tilde{\omega}_2 \tilde{\omega}_3 (I_3 - I_2) &= 0 \\ I_2 \dot{\tilde{\omega}}_2 + \tilde{\omega}_3 \tilde{\omega}_1 (I_1 - I_3) &= 0 \\ I_3 \dot{\tilde{\omega}}_3 + \tilde{\omega}_1 \tilde{\omega}_2 (I_2 - I_1) &= 0 \end{aligned}$$

Solution for asymmetric top -- $I_3 \neq I_2 \neq I_1$:

Suppose : $\dot{\tilde{\omega}}_3 \approx 0$ Define : $\Omega_1 \equiv \tilde{\omega}_3 \frac{I_3 - I_2}{I_1}$

Define : $\Omega_2 \equiv \tilde{\omega}_3 \frac{I_3 - I_1}{I_2}$

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$$\begin{aligned} I_1 \dot{\tilde{\omega}}_1 + \tilde{\omega}_2 \tilde{\omega}_3 (I_3 - I_2) &= 0 \\ I_2 \dot{\tilde{\omega}}_2 + \tilde{\omega}_3 \tilde{\omega}_1 (I_1 - I_3) &= 0 \\ I_3 \dot{\tilde{\omega}}_3 + \tilde{\omega}_1 \tilde{\omega}_2 (I_2 - I_1) &= 0 \end{aligned}$$

$\dot{\tilde{\omega}}_3 \approx 0$ Define : $\Omega_1 \equiv \tilde{\omega}_3 \frac{I_3 - I_2}{I_1}$ $\Omega_2 \equiv \tilde{\omega}_3 \frac{I_3 - I_1}{I_2}$

$\dot{\tilde{\omega}}_1 = -\Omega_1 \tilde{\omega}_2$ $\dot{\tilde{\omega}}_2 = \Omega_2 \tilde{\omega}_1$

If Ω_1 and Ω_2 are both positive or both negative :

$$\tilde{\omega}_1(t) \approx A \cos(\sqrt{\Omega_1 \Omega_2} t + \varphi)$$

$$\tilde{\omega}_2(t) \approx A \sqrt{\frac{\Omega_2}{\Omega_1}} \sin(\sqrt{\Omega_1 \Omega_2} t + \varphi)$$

\Rightarrow If Ω_1 and Ω_2 have opposite signs, solution is unstable.

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Transformation between body-fixed and inertial coordinate systems – Euler angles

http://en.wikipedia.org/wiki/Euler_angles

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$\tilde{\omega} = \dot{\alpha} \hat{e}_3^0 + \dot{\beta} \hat{e}_2' + \dot{\gamma} \hat{e}_3$

Need to express all components in body-fixed frame:

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$\tilde{\omega} = \dot{\alpha} \hat{e}_3^0 + \dot{\beta} \hat{e}_2' + \dot{\gamma} \hat{e}_3$

$\hat{e}_2' = \sin \gamma \hat{e}_1 + \cos \gamma \hat{e}_2$

Matrix representation :

$$\hat{e}_2' = \begin{pmatrix} \cos \gamma & \sin \gamma & 0 \\ -\sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \sin \gamma \\ \cos \gamma \\ 0 \end{pmatrix}$$

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$$\tilde{\omega} = \dot{\alpha} \hat{e}_3^0 + \dot{\beta} \hat{e}_2' + \dot{\gamma} \hat{e}_3'$$

$$\hat{e}_3^0 = -\sin \beta \hat{e}_1' + \cos \beta \hat{e}_3'$$

Matrix representation :

$$\hat{e}_3^0 = \begin{pmatrix} \cos \gamma & \sin \gamma & 0 \\ -\sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cos \beta & 0 & -\sin \beta \\ 0 & 1 & 0 \\ \sin \beta & 0 & \cos \beta \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$= \begin{pmatrix} -\sin \beta \cos \gamma \\ \sin \beta \sin \gamma \\ \cos \beta \end{pmatrix}$$

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$$\tilde{\omega} = \dot{\alpha} \hat{e}_3^0 + \dot{\beta} \hat{e}_2' + \dot{\gamma} \hat{e}_3'$$

$$\tilde{\omega} = \dot{\alpha} \begin{pmatrix} -\sin \beta \cos \gamma \\ \sin \beta \sin \gamma \\ \cos \beta \end{pmatrix} + \dot{\beta} \begin{pmatrix} \sin \gamma \\ \cos \gamma \\ 0 \end{pmatrix} + \dot{\gamma} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\tilde{\omega} = \tilde{\omega}_1 \hat{e}_1' + \tilde{\omega}_2 \hat{e}_2' + \tilde{\omega}_3 \hat{e}_3'$$

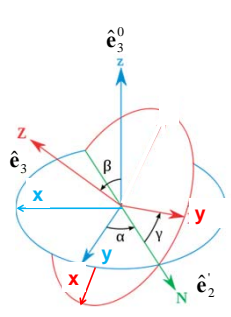
$$\tilde{\omega} = \dot{\alpha} \begin{pmatrix} -\sin \beta \cos \gamma \\ \sin \beta \sin \gamma \\ \cos \beta \end{pmatrix} + \dot{\beta} \begin{pmatrix} \sin \gamma \\ \cos \gamma \\ 0 \end{pmatrix} + \dot{\gamma} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\tilde{\omega}_1 = \dot{\alpha}(-\sin \beta \cos \gamma) + \dot{\beta} \sin \gamma$$

$$\tilde{\omega}_2 = \dot{\alpha}(\sin \beta \sin \gamma) + \dot{\beta} \cos \gamma$$

$$\tilde{\omega}_3 = \dot{\alpha} \cos \beta + \dot{\gamma}$$

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$$\tilde{\omega} = \dot{\alpha} \hat{e}_3^0 + \dot{\beta} \hat{e}_2' + \dot{\gamma} \hat{e}_3'$$


$$\tilde{\omega} = [\dot{\alpha}(-\sin \beta \cos \gamma) + \dot{\beta} \sin \gamma] \hat{e}_1' + [\dot{\alpha}(\sin \beta \sin \gamma) + \dot{\beta} \cos \gamma] \hat{e}_2' + [\dot{\alpha} \cos \beta + \dot{\gamma}] \hat{e}_3'$$

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Rotational kinetic energy

$$\begin{aligned}
 T(\alpha, \beta, \gamma, \dot{\alpha}, \dot{\beta}, \dot{\gamma}) &= \frac{1}{2} I_1 \tilde{\omega}_1^2 + \frac{1}{2} I_2 \tilde{\omega}_2^2 + \frac{1}{2} I_3 \tilde{\omega}_3^2 \\
 &= \frac{1}{2} I_1 [\dot{\alpha}(-\sin \beta \cos \gamma) + \dot{\beta} \sin \gamma]^2 \\
 &\quad + \frac{1}{2} I_2 [\dot{\alpha}(\sin \beta \sin \gamma) + \dot{\beta} \cos \gamma]^2 \\
 &\quad + \frac{1}{2} I_3 [\dot{\alpha} \cos \beta + \dot{\gamma}]^2
 \end{aligned}$$

If $I_1 = I_2$:

$$T(\alpha, \beta, \gamma, \dot{\alpha}, \dot{\beta}, \dot{\gamma}) = \frac{1}{2} I_1 (\dot{\alpha}^2 \sin^2 \beta + \dot{\beta}^2) + \frac{1}{2} I_3 [\dot{\alpha} \cos \beta + \dot{\gamma}]^2$$

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