

**PHY 711 Classical Mechanics and
Mathematical Methods
10-10:50 AM MWF Olin 103**

Plan for Lecture 12:

Continue reading Chapter 3 & 6

- 1. Constructing the Hamiltonian**
- 2. Hamilton's canonical equation**
- 3. Examples**

Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

| | Date | F&W Reading | Topic | Assignment |
|----|----------------|-------------|--|---------------------|
| 1 | Wed, 8/29/2012 | Chap. 1 | Review of basic principles; Scattering theory | #1 |
| 2 | Fri, 8/31/2012 | Chap. 1 | Scattering theory continued | #2 |
| 3 | Mon, 9/03/2012 | Chap. 1 | Scattering theory continued | #3 |
| 4 | Wed, 9/05/2012 | Chap. 1 & 2 | Scattering theory/Accelerated coordinate frame | #4 |
| 5 | Fri, 9/07/2012 | Chap. 2 | Accelerated coordinate frame | #5 |
| 6 | Mon, 9/10/2012 | Chap. 3 | Calculus of Variation | #6 |
| 7 | Wed, 9/12/2012 | Chap. 3 | Calculus of Variation continued | |
| 8 | Fri, 9/14/2012 | Chap. 3 | Lagrangian | #7 |
| 9 | Mon, 9/17/2012 | Chap. 3 & 6 | Lagrangian | #8 |
| 10 | Wed, 9/19/2012 | Chap. 3 & 6 | Lagrangian | #9 |
| 11 | Fri, 9/21/2012 | Chap. 3 & 6 | Lagrangian | #10 |
| 12 | Mon, 9/24/2012 | Chap. 3 & 6 | Lagrangian and Hamiltonian | #11 |



Lagrangian picture

For independent generalized coordinates $q_\sigma(t)$:

$$L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} = 0$$

\Rightarrow Second order differential equations for $q_\sigma(t)$

Switching variables – Legendre transformation

Define: $H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$

$$H = \sum_{\sigma} \dot{q}_\sigma p_\sigma - L \quad \text{where } p_\sigma = \frac{\partial L}{\partial \dot{q}_\sigma}$$

$$dH = \sum_{\sigma} \left(\dot{q}_\sigma dp_\sigma + p_\sigma d\dot{q}_\sigma - \frac{\partial L}{\partial q_\sigma} dq_\sigma - \frac{\partial L}{\partial \dot{q}_\sigma} d\dot{q}_\sigma \right) - \frac{\partial L}{\partial t} dt$$

Hamiltonian picture – continued

$$H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$$

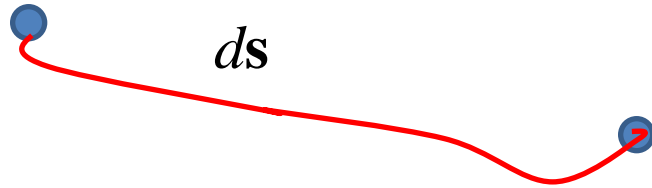
$$H = \sum_\sigma \dot{q}_\sigma p_\sigma - L \quad \text{where} \quad p_\sigma = \frac{\partial L}{\partial \dot{q}_\sigma}$$

$$dH = \sum_\sigma \left(\dot{q}_\sigma dp_\sigma + p_\sigma d\dot{q}_\sigma - \frac{\partial L}{\partial q_\sigma} dq_\sigma - \frac{\partial L}{\partial \dot{q}_\sigma} d\dot{q}_\sigma \right) - \frac{\partial L}{\partial t} dt$$

$$= \sum_\sigma \left(\frac{\partial H}{\partial q_\sigma} dq_\sigma + \frac{\partial H}{\partial p_\sigma} dp_\sigma \right) + \frac{\partial H}{\partial t} dt$$

$$\Rightarrow \dot{q}_\sigma = \frac{\partial H}{\partial p_\sigma} \quad \frac{\partial L}{\partial q_\sigma} = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} \equiv \dot{p}_\sigma = -\frac{\partial H}{\partial q_\sigma} \quad \frac{\partial L}{\partial t} = -\frac{\partial H}{\partial t}$$

Direct application of Hamiltonian's principle using the Hamiltonian function --



Generalized coordinates :
 $q_\sigma(\{x_i\})$

Define -- Lagrangian : $L \equiv T - U$

$$L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t)$$

$$\Rightarrow \text{Minimization integral : } S = \int_{t_i}^{t_f} L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) dt$$

Expressed in terms of Hamiltonian :

$$H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$$

$$H = \sum_{\sigma} \dot{q}_\sigma p_\sigma - L \quad \Rightarrow \quad L = \sum_{\sigma} \dot{q}_\sigma p_\sigma - H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$$

Hamilton's principle continued:

Minimization integral :

$$S = \int_{t_i}^{t_f} \left(\sum_{\sigma} \dot{q}_{\sigma} p_{\sigma} - H(\{q_{\sigma}(t)\}, \{p_{\sigma}(t)\}, t) \right) dt$$

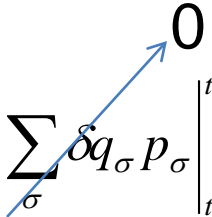
$$\delta S = \int_{t_i}^{t_f} \left(\sum_{\sigma} \left(\dot{q}_{\sigma} \delta p_{\sigma} + \delta \dot{q}_{\sigma} p_{\sigma} - \frac{\partial H}{\partial q_{\sigma}} \delta q_{\sigma} - \frac{\partial H}{\partial p_{\sigma}} \delta p_{\sigma} \right) \right) dt = 0$$

$$\Rightarrow \dot{q}_{\sigma} = \frac{\partial H}{\partial p_{\sigma}}$$

$$\Rightarrow \dot{p}_{\sigma} = -\frac{\partial H}{\partial q_{\sigma}}$$

Canonical equations

Detail :

$$\int_{t_i}^{t_f} \left(\sum_{\sigma} (\delta \dot{q}_{\sigma} p_{\sigma}) \right) dt = \int_{t_i}^{t_f} \left(\sum_{\sigma} \left(\frac{d(\delta q_{\sigma} p_{\sigma})}{dt} - \delta q_{\sigma} \dot{p}_{\sigma} \right) \right) dt = \sum_{\sigma} \delta q_{\sigma} p_{\sigma} \Big|_{t_i}^{t_f} - \int_{t_i}^{t_f} \left(\sum_{\sigma} (\delta q_{\sigma} \dot{p}_{\sigma}) \right) dt$$


Constants of the motion in Hamiltonian formalism

$$H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$$

$$\frac{dq_\sigma}{dt} = \frac{\partial H}{\partial p_\sigma} \Rightarrow \text{constant } q_\sigma \text{ if } \frac{\partial H}{\partial p_\sigma} = 0$$

$$\frac{dp_\sigma}{dt} = -\frac{\partial H}{\partial q_\sigma} \Rightarrow \text{constant } p_\sigma \text{ if } \frac{\partial H}{\partial q_\sigma} = 0$$

$$\frac{dH}{dt} = \sum_\sigma \left(\frac{\partial H}{\partial q_\sigma} \dot{q}_\sigma + \frac{\partial H}{\partial p_\sigma} \dot{p}_\sigma \right) + \frac{\partial H}{\partial t}$$

$$\Rightarrow \text{constant } H \text{ if } \frac{\partial H}{\partial t} = 0$$

Recipe for constructing the Hamiltonian

1. Construct Lagrangian function : $L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$

2. Compute generalized momenta : $p_\sigma \equiv \frac{\partial L}{\partial \dot{q}_\sigma}$

3. Construct Hamiltonian expression : $H = \sum_\sigma \dot{q}_\sigma p_\sigma - L$

4. Form Hamiltonian function : $H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$

Example 1: one - dimensional potential :

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(z)$$

$$p_x = m\dot{x} \quad p_y = m\dot{y} \quad p_z = m\dot{z}$$

$$H = m\dot{x}^2 + m\dot{y}^2 + m\dot{z}^2 - \left(\frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(z) \right)$$

$$H = \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m} + V(z)$$

Constants : p_x, p_y, H

Example 2: Motion in a central potential

$$L = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) - V(r)$$

$$p_r = m\dot{r} \quad p_\phi = mr^2\dot{\phi}$$

$$\begin{aligned} H &= m\dot{r}^2 + mr^2\dot{\phi}^2 - \left(\frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) - V(r)\right) \\ &= \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) + V(r) \end{aligned}$$

$$H = \frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} + V(r)$$

Constants: p_ϕ, H

Other examples

Lagrangian for symmetric top with Euler angles α, β, γ :

$$L = L(\alpha, \beta, \gamma, \dot{\alpha}, \dot{\beta}, \dot{\gamma}) = \frac{1}{2} I_1 (\dot{\alpha}^2 \sin^2 \beta + \dot{\beta}^2) + \frac{1}{2} I_3 (\dot{\alpha} \cos \beta + \dot{\gamma})^2 - Mgh \cos \beta$$

$$p_\alpha = I_1 \dot{\alpha} \sin^2 \beta + I_3 (\dot{\alpha} \cos \beta + \dot{\gamma}) \cos \beta$$

$$p_\beta = I_1 \dot{\beta}$$

$$p_\gamma = I_3 (\dot{\alpha} \cos \beta + \dot{\gamma})$$

$$H = \frac{1}{2} I_1 (\dot{\alpha}^2 \sin^2 \beta + \dot{\beta}^2) + \frac{1}{2} I_3 (\dot{\alpha} \cos \beta + \dot{\gamma})^2 + Mgh \cos \beta$$

$$H = \frac{(p_\alpha - p_\gamma \cos \beta)^2}{2I_1 \sin^2 \beta} + \frac{p_\beta^2}{2I_1} + \frac{p_\gamma^2}{2I_3} + Mgh \cos \beta$$

Constants : p_α, p_γ, H

Other examples

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{q}{2c} B_0 (\dot{x}y - \dot{y}x)$$

$$p_x = m\dot{x} + \frac{q}{2c} B_0 y$$

$$p_y = m\dot{y} - \frac{q}{2c} B_0 x$$

$$p_z = m\dot{z}$$

$$H = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

$$H = \left(\frac{p_x - \frac{q}{2c} B_0 y}{2m} \right)^2 + \left(\frac{p_y + \frac{q}{2c} B_0 x}{2m} \right)^2 + \frac{p_z^2}{2m}$$

Constants : p_z, H

Poisson brackets:

Recall:

$$H = H(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$$

$$\frac{dq_\sigma}{dt} = \frac{\partial H}{\partial p_\sigma} \Rightarrow \text{constant } q_\sigma \text{ if } \frac{\partial H}{\partial p_\sigma} = 0$$

$$\frac{dp_\sigma}{dt} = -\frac{\partial H}{\partial q_\sigma} \Rightarrow \text{constant } p_\sigma \text{ if } \frac{\partial H}{\partial q_\sigma} = 0$$

$$\frac{dH}{dt} = \sum_\sigma \left(\frac{\partial H}{\partial q_\sigma} \dot{q}_\sigma + \frac{\partial H}{\partial p_\sigma} \dot{p}_\sigma \right) + \frac{\partial H}{\partial t}$$

Similarly for an arbitrary function : $F = F(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$

$$\frac{dF}{dt} = \sum_\sigma \left(\frac{\partial F}{\partial q_\sigma} \dot{q}_\sigma + \frac{\partial F}{\partial p_\sigma} \dot{p}_\sigma \right) + \frac{\partial F}{\partial t} = \sum_\sigma \left(\frac{\partial F}{\partial q_\sigma} \frac{\partial H}{\partial p_\sigma} - \frac{\partial F}{\partial p_\sigma} \frac{\partial H}{\partial q_\sigma} \right) + \frac{\partial F}{\partial t}$$

Poisson brackets -- continued:

For an arbitrary function : $F = F(\{q_\sigma(t)\}, \{p_\sigma(t)\}, t)$

$$\frac{dF}{dt} = \sum_{\sigma} \left(\frac{\partial F}{\partial q_{\sigma}} \dot{q}_{\sigma} + \frac{\partial F}{\partial p_{\sigma}} \dot{p}_{\sigma} \right) + \frac{\partial F}{\partial t} = \sum_{\sigma} \left(\frac{\partial F}{\partial q_{\sigma}} \frac{\partial H}{\partial p_{\sigma}} - \frac{\partial F}{\partial p_{\sigma}} \frac{\partial H}{\partial q_{\sigma}} \right) + \frac{\partial F}{\partial t}$$

Define :

$$[F, G]_{PB} \equiv \sum_{\sigma} \left(\frac{\partial F}{\partial q_{\sigma}} \frac{\partial G}{\partial p_{\sigma}} - \frac{\partial F}{\partial p_{\sigma}} \frac{\partial G}{\partial q_{\sigma}} \right) = -[G, F]_{PB}$$

So that :
$$\frac{dF}{dt} = [F, H]_{PB} + \frac{\partial F}{\partial t}$$

Poisson brackets -- continued:

$$[F, G]_{PB} \equiv \sum_{\sigma} \left(\frac{\partial F}{\partial q_{\sigma}} \frac{\partial G}{\partial p_{\sigma}} - \frac{\partial F}{\partial p_{\sigma}} \frac{\partial G}{\partial q_{\sigma}} \right) = -[G, F]_{PB}$$

Examples :

$$[x, x]_{PB} = 0 \quad [x, p_x]_{PB} = 1 \quad [x, p_y]_{PB} = 0$$

Liouville theorem

Let $D \equiv$ density of particles in phase space :

$$\frac{dD}{dt} = [D, H]_{PB} + \frac{\partial D}{\partial t} = 0$$