Based on the charge and strangeness of the octet of lightest baryons, deduce what quarks are inside them (don't worry about symmetrizing, just write the answer as | p > = |uud >). It is acceptable to repeat.

This is straightforward, and you did most of this in the reading quiz, so it is easy.

$$|p\rangle = |uud\rangle, \quad |n\rangle = |udd\rangle, \quad |\Sigma^{+}\rangle = |uus\rangle, \quad |\Sigma^{0}\rangle = |uds\rangle, \quad |\Sigma^{-}\rangle = |dds\rangle, \\ |\Lambda^{0}\rangle = |dds\rangle, \quad |\Xi^{0}\rangle = |uss\rangle, \quad |\Xi^{-}\rangle = |dss\rangle.$$

- Although the 1, 8, and 10 dimensional representations of SU(3)<sub>F</sub> are all that are used for *light* quarks, there are some objects containing heavy quarks that are different. The Λ<sub>c</sub> particle is a heavy spin ½ baryon with quark content |Λ<sub>c</sub>⟩ = ¼(|cud⟩ |cdu⟩). The charmed quark is unaffected by SU(3)<sub>F</sub>, so T<sub>a</sub> |c⟩ = 0 for all a.
  - (a) Suppose all the operators  $T_{i \to j}$  are allowed to act repeatedly on  $|\Lambda_c\rangle$ . List all of the states that you can get. Make up names for them.

The operators  $T_{i \to j}$  turn one type of quark into another, so we should probably rename them; for example,  $T_{1 \to 2}$  would be renamed  $T_{u \to d}$ . We find, for example, that

$$\begin{split} T_{u \to d} \left| \Lambda_c \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| cdd \right\rangle - \left| cdd \right\rangle \right) = 0 , \quad T_{d \to u} \left| \Lambda_c \right\rangle = \frac{1}{\sqrt{2}} \left( \left| cuu \right\rangle - \left| cuu \right\rangle \right) = 0 , \\ T_{u \to s} \left| \Lambda_c \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| csd \right\rangle - \left| cds \right\rangle \right) = \left| \Xi_c^0 \right\rangle , \quad T_{d \to s} \left| \Lambda_c \right\rangle = \frac{1}{\sqrt{2}} \left( \left| cus \right\rangle - \left| csu \right\rangle \right) = \left| \Xi_c^+ \right\rangle , \\ T_{s \to u} \left| \Lambda_c \right\rangle &= T_{s \to d} \left| \Lambda_c \right\rangle = 0 . \end{split}$$

It is not hard to see that if we let the various generators act on these states, we don't get any additional states. We find

$$\begin{split} T_{u \to d} \left| \Xi_{c}^{0} \right\rangle &= T_{u \to s} \left| \Xi_{c}^{0} \right\rangle = 0, \\ T_{d \to u} \left| \Xi_{c}^{0} \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| csu \right\rangle - \left| cus \right\rangle \right) = - \left| \Xi_{c}^{+} \right\rangle, \quad T_{d \to s} \left| \Xi_{c}^{0} \right\rangle = \frac{1}{\sqrt{2}} \left( \left| css \right\rangle - \left| css \right\rangle \right) = 0, \\ T_{s \to u} \left| \Xi_{c}^{0} \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| cud \right\rangle - \left| cdu \right\rangle \right) = \left| \Lambda_{c} \right\rangle, \quad T_{s \to d} \left| \Xi_{c}^{0} \right\rangle = \frac{1}{\sqrt{2}} \left( \left| cdd \right\rangle - \left| cdd \right\rangle \right) = 0, \\ T_{d \to u} \left| \Xi_{c}^{+} \right\rangle &= T_{d \to s} \left| \Xi_{c}^{+} \right\rangle = 0, \\ T_{u \to d} \left| \Xi_{c}^{+} \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| cds \right\rangle - \left| csd \right\rangle \right) = - \left| \Xi_{c}^{0} \right\rangle, \quad T_{u \to s} \left| \Xi_{c}^{+} \right\rangle = \frac{1}{\sqrt{2}} \left( \left| css \right\rangle - \left| css \right\rangle \right) = 0, \\ T_{s \to u} \left| \Xi_{c}^{+} \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| cuu \right\rangle - \left| cuu \right\rangle \right) = 0, \quad T_{s \to d} \left| \Xi_{c}^{+} \right\rangle = \frac{1}{\sqrt{2}} \left( \left| cud \right\rangle - \left| cdu \right\rangle \right) = \left| \Lambda_{c} \right\rangle. \end{split}$$

## (b)Identify the $T_3$ and $T_8$ values for all the states you found in part (a) (including $|\Lambda_c\rangle$ ). Make a plot of $T_3$ vs. $T_8$ for these states.

The  $T_3$  and  $T_8$  values for each of these is just the sum of the corresponding eigenvalues for the two quarks (the charm quark doesn't contribute). We therefore have

$$\begin{split} T_{3} \left| \Lambda_{c} \right\rangle &= \left( \frac{1}{2} - \frac{1}{2} \right) \left| \Lambda_{c} \right\rangle = 0, \qquad T_{8} \left| \Lambda_{c} \right\rangle = \left( \frac{1}{2\sqrt{3}} + \frac{1}{2\sqrt{3}} \right) \left| \Lambda_{c} \right\rangle = \frac{1}{\sqrt{3}} \left| \Lambda_{c} \right\rangle, \\ T_{3} \left| \Xi_{c}^{0} \right\rangle &= \left( 0 - \frac{1}{2} \right) \left| \Xi_{c}^{0} \right\rangle = -\frac{1}{2} \left| \Xi_{c}^{0} \right\rangle, \qquad T_{8} \left| \Xi_{c}^{0} \right\rangle = \left( \frac{1}{2\sqrt{3}} - \frac{1}{\sqrt{3}} \right) \left| \Xi_{c}^{0} \right\rangle = -\frac{1}{2\sqrt{3}} \left| \Xi_{c}^{0} \right\rangle, \\ T_{3} \left| \Xi^{+} \right\rangle &= \left( 0 + \frac{1}{2} \right) \left| \Xi_{c}^{+} \right\rangle = +\frac{1}{2} \left| \Xi_{c}^{0} \right\rangle, \qquad T_{8} \left| \Xi^{+} \right\rangle = \left( \frac{1}{2\sqrt{3}} - \frac{1}{\sqrt{3}} \right) \left| \Xi_{c}^{+} \right\rangle = -\frac{1}{2\sqrt{3}} \left| \Xi_{c}^{0} \right\rangle. \end{split}$$

The corresponding plot appears above right.

4. Draw at least thirteen tree-level Feynman diagrams for the process  $q\overline{q} \rightarrow ggg$ , where q is any quark, and g is a gluon. You don't have to do anything with the diagrams. For extra credit, find them all.

I have colored the three external gluon lines to help distinguish them from each other. For the seven diagrams with no an interior gluon line, I simply left that line black. There are six diagrams with three gluon/fermion vertices, six with two fermion/gluon vertex and a three gluon vertex, and with a single fermion/gluon vertex and a four gluon vertex, and three with a single fermion/gluon vertices.

