

Physics 780 – General Relativity  
Solutions to Homework A

1. [5] Using the convention that  $c = 1$ , rewrite the following quantities (on personal ones, you can lie, but your lie has to be plausible)
- (a) [1.5] Your age in m  
 (b) [1.5] The length of your foot in ns  
 (c) [2] The Schwarzschild radius of the Sun, given by  $2GM$ , in km

Your answer will depend on your physical characteristic and what star your home planet orbits, but for Dr. Carlson he is about 60.8 year old, his foot is about a foot long, and the Sun's mass is about  $1.989 \times 10^{30}$  kg, so

$$t = tc = (60.8 \text{ y})(3.156 \times 10^7 \text{ s/y})(2.998 \times 10^8 \text{ m/s}) = 5.75 \times 10^{17} \text{ m},$$

$$L = \frac{L}{c} = \frac{(12.0 \text{ in})(0.0254 \text{ m/in})}{2.998 \times 10^8 \text{ m/s}} = 1.017 \times 10^{-9} \text{ s} = 1.017 \text{ ns},$$

$$2GM = \frac{2GM}{c^2} = \frac{2(6.674 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{s}^{-3})(1.989 \times 10^{30} \text{ kg})}{(2.998 \times 10^8 \text{ m/s})^2} = 2954 \text{ m} = 2.954 \text{ km}.$$

2. [5] A particle moves in a helical path given by  $(x, y, z) = (R \cos(\omega t), R \sin(\omega t), vt)$ . Find a relation between the coordinate time  $t$  and the proper time  $\tau$ , and then rewrite all four coordinates in terms of  $\tau$ .

We have

$$\begin{aligned} d\tau^2 &= dt^2 - dx^2 - dy^2 - dz^2 = dt^2 - (-R\omega \sin(\omega t) dt)^2 - (R\omega \cos(\omega t) dt)^2 - (v dt)^2 \\ &= dt^2 [1 - R^2 \omega^2 (\sin^2(\omega t) + \cos^2(\omega t)) - v^2] = dt^2 (1 - R^2 \omega^2 - v^2), \\ \tau &= \int dt \sqrt{1 - R^2 \omega^2 - v^2} = t \sqrt{1 - R^2 \omega^2 - v^2}. \end{aligned}$$

We can rewrite this as  $t = \gamma \tau$ , where  $\gamma = \frac{1}{\sqrt{1 - R^2 \omega^2 - v^2}}$ , then it is clear that our coordinates are

$$t = \gamma \tau, \quad x = R \cos(\omega \gamma \tau), \quad y = R \sin(\omega \gamma \tau), \quad z = \gamma v \tau, \quad \text{where } \gamma = \frac{1}{\sqrt{1 - \omega^2 R^2 - v^2}}.$$

3. [10] A general Lorentz transformation is a  $4 \times 4$  matrix satisfying  $\eta_{\mu\nu} = \eta_{\alpha\beta} \Lambda^\alpha_\mu \Lambda^\beta_\nu$ , which can be written in matrix form as  $\eta = \Lambda^T \eta \Lambda$

(a) [3] Using the formulas  $\det(AB) = \det(A)\det(B)$  and  $\det(A^T) = \det(A)$ , show that  $\det(\Lambda) = \pm 1$ .

Working in the matrix notation, we have

$$\det(\eta) = \det(\Lambda^T) \det(\eta) \det(\Lambda) = \det(\eta) [\det(\Lambda)]^2,$$

$$[\det(\Lambda)]^2 = 1, \quad \det(\Lambda) = \pm 1.$$

(b) [3] Using this equation in the case  $\mu = \nu = 0$ , show that  $\Lambda^0_0 \geq 1$  or  $\Lambda^0_0 \leq -1$ .

Using the fact that  $\eta_{\alpha\beta}$  has only diagonal elements, the equation for  $\mu = \nu = 0$  yields

$$\begin{aligned} -1 &= \eta_{00} = \eta_{\alpha\beta} \Lambda^\alpha_0 \Lambda^\beta_0 = \eta_{00} \Lambda^0_0 \Lambda^0_0 + \eta_{11} \Lambda^1_0 \Lambda^1_0 + \eta_{22} \Lambda^2_0 \Lambda^2_0 + \eta_{33} \Lambda^3_0 \Lambda^3_0 \\ &= -(\Lambda^0_0)^2 + (\Lambda^1_0)^2 + (\Lambda^2_0)^2 + (\Lambda^3_0)^2, \\ (\Lambda^0_0)^2 &= 1 + (\Lambda^1_0)^2 + (\Lambda^2_0)^2 + (\Lambda^3_0)^2 \end{aligned}$$

It follows that  $(\Lambda^0_0)^2 \geq 1$ , so  $|\Lambda^0_0| \geq 1$  which implies  $\Lambda^0_0 \geq 1$  or  $\Lambda^0_0 \leq -1$ .

(c) [2] Argue that if you start with the identity Lorentz transformation ( $\Lambda = 1$ ), and then continuously change it, by making small rotations or boosts, the sign of  $\det(\Lambda)$  and  $\Lambda^0_0$  will never change. Call these Lorentz transformations *proper Lorentz transformation*.

As we perform small changes to  $\Lambda$ , the determinant and the value of  $\Lambda^0_0$  can only change in a continuous manner. Therefore the determinant, constrained to the values  $\pm 1$ , cannot jump from positive to negative, so it must always have the value  $+1$ . Similarly,  $\Lambda^0_0$  starts at  $+1$ , and since it can't slowly change to a value slightly less than one, it will always have  $\Lambda^0_0 \geq 1$ , also positive.

(d) [2] Show that time reversal  $\Lambda = \mathcal{T} = \text{diag}(-1, 1, 1, 1)$  parity,  $\Lambda = \mathcal{P} = \text{diag}(1, -1, -1, -1)$ , and the combination  $\Lambda = \mathcal{PT}$  are improper Lorentz transformations.

It is trivial to see that  $\mathcal{T}$  and  $\mathcal{P}$  both have determinant minus one, and therefore are improper. The combination  $\mathcal{PT}$  has determinant  $+1$ , but it has  $\Lambda^0_0 = -1$  (as also does  $\mathcal{T}$ ), and hence it is also improper. Hence you can't by small changes turn around in time, turn around like a mirror image, nor can you do both.

4. [5] The concept of future and past do not work exactly the same in special relativity, but some things are the same. We will say that a point  $x^\mu$  is in the absolute future of another point  $y^\mu$  if  $x^0 - y^0 > |\mathbf{x} - \mathbf{y}|$ . In standard physics, the future of the future is the future. Show that in relativity, if  $x$  is in the future of  $y$  and  $y$  is in the future of  $z$ , then  $x$  is in the future of  $z$ . You may want to look up the triangle inequality.

We know that  $x^0 - y^0 > |\mathbf{x} - \mathbf{y}|$  and  $y^0 - z^0 > |\mathbf{y} - \mathbf{z}|$ . Adding these two inequalities, we have  $x^0 - z^0 > |\mathbf{x} - \mathbf{y}| + |\mathbf{y} - \mathbf{z}|$ . We want to prove that  $x^0 - z^0 > |\mathbf{x} - \mathbf{z}|$ . All the we need to finish the proof is to show that

$$|\mathbf{x} - \mathbf{y}| + |\mathbf{y} - \mathbf{z}| \geq |\mathbf{x} - \mathbf{z}|.$$

This theorem is the triangle inequality, since it says that the sum of two sides of a triangle is always at least as long as the third side. Mathematically, you can prove this by first defining  $\mathbf{a} = \mathbf{x} - \mathbf{y}$  and  $\mathbf{b} = \mathbf{y} - \mathbf{z}$ . We then substitute this into the expression, and then square it, so we are trying to prove

$$\begin{aligned} |\mathbf{a}| + |\mathbf{b}| &\geq |\mathbf{a} + \mathbf{b}|, \\ a^2 + 2ab + b^2 &\geq (\mathbf{a} + \mathbf{b})^2, \\ a^2 + 2ab + b^2 &\geq a^2 + 2\mathbf{a} \cdot \mathbf{b} + b^2, \\ ab &\geq \mathbf{a} \cdot \mathbf{b} = ab \cos \theta. \end{aligned}$$

Since we always have  $\cos \theta \leq 1$ , this is obvious.